# LOCAL LIMIT THEOREM FOR SYMMETRIC RANDOM WALKS IN GROMOV-HYPERBOLIC GROUPS 

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## 1. Introduction

Consider a countable group $\Gamma$ (with identity denoted by $e$ ), together with a probability measure $\mu$ whose support generates $\Gamma$ as a semigroup (we say that $\mu$ is admissible). Multiplying random elements of $\Gamma$ distributed independently according to $\mu$, one obtains a random walk on $\Gamma$. The local limit problem consists in determining good asymptotics for the transition probabilities $p_{n}(x, y)$ of this random walk. Let us assume for simplicity that $\mu$ is finitely supported. For $\Gamma=\mathbb{Z}^{d}$, simple Fourier computations show that $p_{n}(e, e) \sim C n^{-d / 2}$ if the walk is centered, and $p_{n}(e, e) \sim C R^{-n} n^{-d / 2}$ for some $R>1$ if the walk is not centered. Similar asymptotics hold in nilpotent groups by the deep results of Varopoulos and Alexopoulos Ale02.

When the group is not amenable, $p_{n}(e, e)$ decays exponentially fast. The situation is well understood for semisimple Lie groups and absolutely continuous measures since the work of Bougerol [Bou81]: the probability to return to a fixed neighborhood of the identity behaves like $C R^{-n} n^{-a}$ for some $R>1$ (depending on the measure one considers) and some $a>1$ only depending on the geometry of the group. In the simplest case of rank one groups, $a=3 / 2$. It is reasonable to conjecture that similar asymptotics (with the same $a$ ) hold for random walks on cocompact lattices of such semisimple Lie groups, but the proofs of Bougerol (based on representation theory) do not adapt well, and this question is essentially open.

A notable exception is the case of free groups: in this situation, the generating function of the transition probabilities (also called the Green function) $G_{r}(x, y)=$ $\sum r^{n} p_{n}(x, y)$ is an algebraic function of $r$. A careful study of its first singularity then yields the asymptotics of $p_{n}(x, y)$. For free groups, this is due to Lalley [Lal93], and the asymptotics is of the form $p_{n}(x, y) \sim C(x, y) R^{-n} n^{-3 / 2}$, in accordance with the results of Bougerol in rank one Lie groups. Most free products can also be treated similarly; see [Woe00, Chapter III] and references therein.

Recently, together with Lalley, we were able to treat in GL13 some nonamenable groups where the Green function is not expected to be algebraic. We proved that, for a cocompact lattice of $\operatorname{PSL}(2, \mathbb{R})$, and for a finitely supported symmetric measure $\mu$, the above asymptotics $p_{n}(x, y) \sim C(x, y) R^{-n} n^{-3 / 2}$ still holds. Henceforth, we will refer to this situation as the $\operatorname{PSL}(2, \mathbb{R})$-case. The overall strategy can in fact be

[^0]formulated in any Gromov-hyperbolic group (including in particular all cocompact lattices in rank one semisimple Lie groups), but a crucial point in the proof really relies on two-dimensional geometry. In this article, we provide a completely different argument for this crucial point, making it possible to extend the results of GL13] to any Gromov-hyperbolic group.

We say that the walk is aperiodic if there exists an odd integer $n$ such that $p_{n}(e, e)>0$. In this case, $p_{n}(e, e)>0$ for all large enough $n$.

Theorem 1.1. Let $\Gamma$ be a finitely generated nonelementary Gromov-hyperbolic group. Let $\mu$ be an admissible finitely supported symmetric probability measure on $\Gamma$. Denote by $R>1$ the inverse of the spectral radius of the corresponding random walk. For any $x, y \in \Gamma$, there exists $C(x, y)>0$ such that

$$
p_{n}(x, y) \sim C(x, y) R^{-n} n^{-3 / 2}
$$

if the walk is aperiodic. If the walk is periodic, this asymptotics holds for even (resp. odd) $n$ if the distance from $x$ to $y$ is even (resp. odd).

The proof of the analogous theorem in the $\operatorname{PSL}(2, \mathbb{R})$-case in GL13 is divided into three steps, as follows:
(1) One shows that Ancona's results Anc87 on the Martin boundary extend up to $r=R$. In particular, the Martin kernel $K_{r, \xi}(x)=G_{r}(x, \xi) / G_{r}(e, \xi)$ converges when $\xi$ tends to a point in the geometric boundary of $\Gamma$, uniformly in $r \in[1, R]$.
(2) Using the Cannon automaton coding geodesics in the group, and thermodynamic formalism in the resulting subshift of finite type, one gets estimates for the sums $\sum_{x \in \Gamma} G_{r}(e, x) G_{r}(x, e)$ when $r \rightarrow R$ in terms of a pressure function. This implies that $r \mapsto G_{r}(e, e)$ almost satisfies a differential equation. Asymptotics of this function follow.
(3) From the asymptotics of $G_{r}(e, e)$, one deduces the asymptotics of $p_{n}(e, e)$ using tauberian theorems (and a little bit of spectral theory). The asymptotics of $p_{n}(x, y)$ are proved in the same way.

From this point on, this article is subdivided into three sections, each devoted to one of those three steps. We will give further comments, explain quickly the arguments in [GL13], and insist on the differences between the $\operatorname{PSL}(2, \mathbb{R})$-case and the general case of Gromov-hyperbolic groups. The main difference is in the first step: the proof of GL13] is deeply two-dimensional, and the general argument is completely different. For the second step, a significant technical complication appears: in the $\operatorname{PSL}(2, \mathbb{R})$-case, the Cannon automaton (a combinatorial object coding the geodesics in the group) is transitive, while this is not the case in general. To overcome this difficulty, we use additional information from the first step and a technique of Calegari and Fujiwara [F10. Finally, the third step is exactly the same in the $\operatorname{PSL}(2, \mathbb{R})$-case or in the general case, and we will give some details only for the convenience of the reader.

While we have tried to make this article as self-contained as possible, GL13] provides a good introduction to some concepts and techniques that we use. The letter $C$ denotes a constant that may vary from line to line. Since most arguments work exactly in the same way for symmetric or nonsymmetric measures, we have written most proofs without using the assumption of symmetry. It only plays a role in the proof of Lemma [2.6] (the central lemma to obtain Ancona inequalities)
and in Section 4 We expect that the first step (Ancona inequalities) should be true without any symmetry assumption on the measure. While we are not able to prove it in general, we are able to obtain it for cocompact discrete subgroups of $\operatorname{PSL}(2, \mathbb{R})$. The identification of the Martin boundary at the spectral radius follows. The argument is given in Appendix A

## 2. Ancona inequalities up to the spectral radius

2.1. The Green function. Consider an admissible finitely supported probability measure $\mu$ on a countable group $\Gamma$. It defines a random walk on $\Gamma$. Let $R=$ $R(\mu)=\lim \sup p_{n}(e, e)^{-1 / n}$ (when $\mu$ is symmetric, this is the inverse of the spectral radius of the Markov operator associated to the random walk on $\ell^{2}$ ). The Green function is defined for $1 \leqslant r<R$ and $x, y \in \Gamma$ by $G_{r}(x, y)=\sum r^{n} p_{n}(x, y)$. By a result of Guivarc'h, it is convergent even for $r=R$ if the group carries no recurrent random walk (this is in particular true for nonamenable groups). One should think of $G_{r}(x, y)$ as the average number of passages in $y$ if the random walk starts from $x$, but for the measure $r \mu$ instead of $\mu$. In particular, for larger $r, G_{r}$ gives more weight to longer paths.

If $\gamma=\left(x, x_{1}, \ldots, x_{n-1}, y\right)$ is a path of length $n$ from $x$ to $y$, its $r$-weight $w_{r}(\gamma)$ is $r^{n} \prod_{i=0}^{n-1} p\left(x_{i}, x_{i+1}\right)$ [where $x_{0}=x$ and $x_{n}=y$ by convention, and we write $p(a, b)=$ $\mu\left(a^{-1} b\right)$ for the probability to jump from $a$ to $\left.b\right]$. By definition, $G_{r}(x, y)=\sum w_{r}(\gamma)$, where the sum is over all paths from $x$ to $y$.

If $\Omega$ is a subset of $\Gamma$, one defines the restricted Green function $G_{r}(x, y ; \Omega)$ as $\sum w_{r}(\gamma)$ where the sum is over all paths $\gamma=\left(x, x_{1}, \ldots, x_{n-1}, y\right)$ such that $x_{i} \in \Omega$ for $1 \leqslant i \leqslant n-1$. If $A$ is a subset of $\Gamma$ not containing $x, y$ such that any trajectory of the random walk from $x$ to $y$ has to go through $A$, one has

$$
\begin{equation*}
G_{r}(x, y)=\sum_{a \in A} G_{r}\left(x, a ; A^{c}\right) G_{r}(a, y)=\sum_{a \in A} G_{r}(x, a) G_{r}\left(a, y ; A^{c}\right), \tag{2.1}
\end{equation*}
$$

where $A^{c}$ denotes the complement of $A$. Indeed, the first (resp. second) formula is proved by splitting a path from $x$ to $y$ according to its first (resp. last) visit to $A$. More generally, if $\Omega$ is a subset of $\Gamma$ containing $x$ and $y$, the above formula holds restricted to $\Omega$, i.e.,
$G_{r}(x, y ; \Omega)=\sum_{a \in A \cap \Omega} G_{r}\left(x, a ; A^{c} \cap \Omega\right) G_{r}(a, y ; \Omega)=\sum_{a \in A \cap \Omega} G_{r}(x, a ; \Omega) G_{r}\left(a, y ; A^{c} \cap \Omega\right)$.
Assuming that $\Gamma$ is finitely generated, we can consider a word distance $d$ on $\Gamma$ coming from a finite symmetric generating set. If $x$ and $y$ are at distance $d$, there is a path from $x$ to $y$ whose probability is bounded from below by $C^{-d}$, and staying close to a geodesic segment from $x$ to $y$. We deduce that, for any $z$,

$$
\begin{equation*}
C^{-d(x, y)} \leqslant G_{r}(x, z) / G_{r}(y, z) \leqslant C^{d(x, y)}, \tag{2.2}
\end{equation*}
$$

and similar inequalities hold for the Green functions restricted to any set containing a fixed size neighborhood of a geodesic segment from $x$ to $y$. These inequalities are called Harnack inequalities.

The first visit Green function is $F_{r}(x, y)=G_{r}\left(x, y ;\{y\}^{c}\right)$. It takes into account only the first visits to $y$. For $r=1$, this is the probability to reach $y$ starting from $x$. One has $G_{r}(x, y)=F_{r}(x, y) G_{r}(y, y)=F_{r}(x, y) G_{r}(e, e)$. Moreover,
$F_{r}(x, y) G_{r}(y, z) \leqslant G_{r}(x, z)$ (since the concatenation of a path from $x$ to $y$ with a path from $y$ to $z$ gives a path from $x$ to $z$ ). Dividing by $G_{r}(e, e)$, one gets

$$
\begin{equation*}
F_{r}(x, y) F_{r}(y, z) \leqslant F_{r}(x, z) . \tag{2.3}
\end{equation*}
$$

We also obtain

$$
\begin{equation*}
G_{r}(x, y) G_{r}(y, z) \leqslant G_{R}(e, e) G_{r}(x, z) \tag{2.4}
\end{equation*}
$$

The Martin boundary is the set of pointwise limits of sequences of functions $K_{r, y_{n}}(x)=G_{r}\left(x, y_{n}\right) / G_{r}\left(e, y_{n}\right)$ when $y_{n}$ tends to infinity. Since these functions are normalized by $K_{r, y}(e)=1$, and $r$-harmonic except at $y$, limits exist, are nonzero, and are $r$-harmonic everywhere (since the measure $\mu$ has finite support). Understanding the Martin boundary amounts to understanding for which sequences $y_{n}$ the functions $K_{r, y_{n}}$ converge.

The derivative of $G_{r}$ with respect to $r$ can be computed. Indeed,

$$
\begin{equation*}
\left(r G_{r}(x, y)\right)^{\prime}=\sum_{z \in \Gamma} G_{r}(x, z) G_{r}(z, y) \tag{2.5}
\end{equation*}
$$

where the prime indicates the derivative with respect to $r$. This equation for $x=$ $y=e$ shows that

$$
\begin{equation*}
\eta(r):=\sum_{z} G_{r}(e, z) G_{r}(z, e)<+\infty \quad \text { for all } r<R \tag{2.6}
\end{equation*}
$$

Let us introduce a convenient notation: we shall write

$$
\begin{equation*}
H_{r}(x, y)=G_{r}(x, y) G_{r}(y, x) \tag{2.7}
\end{equation*}
$$

In the symmetric case, this is simply the square of the Green function. With this notation, $\eta(r)=\sum_{z \in \Gamma} H_{r}(e, z)$. Since $G_{r}$ satisfies (2.4), $H_{r}$ also satisfies this inequality.
2.2. Ancona inequalities. Consider a finitely generated group $\Gamma$. Its Cayley graph is endowed with the word metric coming from any finite set of generators. One says that $\Gamma$ is Gromov-hyperbolic (or simply hyperbolic) if there exists $\delta$ such that any geodesic triangle in this Cayley graph is $\delta$-thin; i.e., each side of the triangle is contained in the $\delta$-neighborhood of the union of the two other sides. This notion is invariant under quasi-isometry, and therefore independent of the choice of the generators (see GdlH90 for more details on hyperbolic groups). The geometric intuition to have is that any finite set of points in a hyperbolic group is isometric to a finite set of points in a tree, up to some constant depending only on the number of points. In particular, statements regarding the relative positions of points can be reduced to statements in trees that are easy to check combinatorially. This intuition is made precise by the following theorem [GdlH90, Theorem 2.12].

Theorem 2.1. For any $n \in \mathbb{N}$ and $\delta>0$, there exists a constant $C=C(n, \delta)$ with the following property. Consider a subset $A$ of a $\delta$-hyperbolic space of cardinality at most $n$. There exists a map $\Phi$ from $A$ to a metric tree such that, for any $x, y \in A$,

$$
d(x, y)-C \leqslant d(\Phi(x), \Phi(y)) \leqslant d(x, y)
$$

A hyperbolic group $\Gamma$ (or more generally any geodesic Gromov-hyperbolic space) has a well-defined geometric boundary $\partial \Gamma$ : this is the set of semi-infinite geodesics, where two such geodesics are identified if they stay a bounded distance away. This boundary is a compact space, and $\Gamma \cup \partial \Gamma$ is also compact.

Consider now an admissible finitely supported probability measure $\mu$ on a nonelementary hyperbolic group $\Gamma$ (i.e., not quasi-isometric to $\{0\}$ or $\mathbb{Z}$ ). Let $R=R(\mu)$, which is strictly larger than 1 since $\Gamma$ is not amenable. Ancona proved in Anc87, that, for any $r<R$, the Martin boundary for $r$-harmonic functions coincides with the geometric boundary: $K_{r, y_{n}}$ converges pointwise if and only if $y_{n}$ converges to a point $\xi \in \partial \Gamma$, and the limits are different for different points of the boundary.

A crucial inequality in Ancona's proof is the fact that the converse inequality to (2.4) holds for any $r<R$ whenever $y$ is close to a geodesic from $x$ to $z$ (with a constant a priori depending on $r$ ). In other words, typical trajectories from $x$ to $z$ follow the geodesic sufficiently well so that they are likely to pass close to $y$. When $r$ increases, $G_{r}$ gives more and more weight to long trajectories, which are more likely to go further from the geodesic. Hence, this Ancona estimate is more and more subtle when $r$ increases. Proving such an estimate for $r=R$ is a crucial step in the proof of Theorem 1.1.

Definition 2.2. A probability measure $\mu$ on a Gromov-hyperbolic group $\Gamma$ satisfies uniform Ancona inequalities if there exists a constant $C>0$ such that, for any $x, z \in \Gamma$ and for any $y$ close to a geodesic segment from $x$ to $z$, for any $r \in[1, R(\mu)]$,

$$
G_{r}(x, z) \leqslant C G_{r}(x, y) G_{r}(y, z)
$$

Theorem 2.3. If $\mu$ is admissible, finitely supported, and symmetric on a nonelementary Gromov-hyperbolic group, it satisfies uniform Ancona inequalities.

This result has been proved in GL13 for cocompact lattices of $\operatorname{PSL}(2, \mathbb{R})$, using very specific two-dimensional arguments. The main idea in the new argument to follow is to combine a supermultiplicativity estimate (originating in DPPS11 for counting problems) with a geometric construction of random barriers in hyperbolic space. We will write $|x|$ for the distance of $x$ to the identity $e$ (for some fixed word distance), and $\mathbb{S}_{k}$ for the sphere of radius $k$ around $e$. The rest of this section is devoted to the proof of Theorem 2.3. We fix a nonelementary Gromov-hyperbolic group $\Gamma$ and an admissible probability measure $\mu$. We do not assume yet that $\mu$ is symmetric, since it will only be important in Lemma 2.6 below.

Lemma 2.4. There exists $C>0$ such that, for any $x, y \in \Gamma$, there exists $a \in \Gamma$ of length at most $C$ such that $|x a y| \geqslant|x|+|y|$.

Proof. Fix $C_{0}=C(4, \delta)>0$ such that any configuration of at most four points can be approximated by a tree with error at most $C_{0}$, as in Theorem 2.1.

We will rely on the classical construction of free groups with two generators in $\Gamma$ as follows. A hyperbolic element of $\Gamma$ is an element $u$ of $\Gamma$ such that the left-multiplication by $u$ has two fixed points at infinity, an attracting one and a repelling one, denoted by $u_{+}$and $u_{-}$. Consider two hyperbolic elements $u, v$ in $\Gamma$ such that the four points $u_{+}, u_{-}, v_{+}, v_{-}$are distinct (this is possible since $\Gamma$ is nonelementary; see the proof of GdlH90, Theorem 8.37]), and fix small disjoint neighborhoods $V\left(u_{+}\right), V\left(u_{-}\right), V\left(v_{+}\right), V\left(v_{-}\right)$of those points in $\Gamma \cup \partial \Gamma$. If $N$ is large enough, any $x$ in the complement of $V\left(a_{+}\right)$(for $a \in F=\left\{u, u^{-1}, v, v^{-1}\right\}$ ) shares only a short beginning with $a^{N}$. More precisely, there exists $K>0$ independent of $N$ such that, in a tree approximation $\Phi$ of $e, x, a^{N}$ with error at most $C_{0}$, the branches from $\Phi(e)$ leading to $\Phi(x)$ and $\Phi\left(a^{N}\right)$ split before time $K$. Increasing $N$, we can also assume that $\left|a^{N}\right| \geqslant 4 K+3 C_{0}$ for all $a \in F$.


Figure 1. Approximating tree for $e, x, x a^{N}, x a^{N} y$.
Consider now two points $x, y \in \Gamma$. One can choose $a \in F$ such that $x^{-1} \notin$ $V\left(a_{+}\right)$and $y \notin V\left(a_{-}\right)$. Consider a tree approximation $\Phi:\left\{e, x, x a^{N}, x a^{N} y\right\} \rightarrow$ $T$. The geodesic paths from $\Phi(x)$ to, respectively, $\Phi(e)$ and $\Phi\left(x a^{N}\right)$, split before time $K$ by construction since $x^{-1} \notin V\left(a_{+}\right)$. In the same way, the paths from $\Phi\left(x a^{N}\right)$ to, respectively, $\Phi(x)$ and $\Phi\left(x a^{N} y\right)$ also split before time $K$ since $y \notin V\left(a_{-}\right)$ (see Figure [1). Hence,

$$
\begin{aligned}
\left|x a^{N} y\right| & \geqslant d\left(\Phi(e), \Phi\left(x a^{N} y\right)\right) \\
& \geqslant d(\Phi(e), \Phi(x))+d\left(\Phi(x), \Phi\left(x a^{N}\right)\right)+d\left(\Phi\left(x a^{N}\right), \Phi\left(x a^{N} y\right)\right)-2 \cdot 2 K \\
& \geqslant|x|-C_{0}+\left|a^{N}\right|-C_{0}+|y|-C_{0}-4 K \\
& \geqslant|x|+|y| .
\end{aligned}
$$

We recall the notation $H_{r}(x, y)=G_{r}(x, y) G_{r}(y, x)$ from (2.7).
Lemma 2.5. There exists $C>0$ such that, for any $k \in \mathbb{N}, \sum_{x \in \mathbb{S}_{k}} H_{R}(e, x) \leqslant C$.
Proof. Fix $r<R$. Write $u_{k}(r)=\sum_{x \in \mathbb{S}_{k}} H_{r}(e, x)$. To $x \in \mathbb{S}_{k}$ and $y \in \mathbb{S}_{\ell}$ one can associate thanks to the previous lemma a point $\Psi(x, y)=x a y \in \bigcup_{k+\ell \leqslant i \leqslant k+\ell+C} \mathbb{S}_{i}$. By (2.4), we have

$$
\begin{aligned}
H_{r}(e, x) H_{r}(e, y) & \leqslant C H_{r}(e, x) H_{r}(e, a) H_{r}(e, y)=C H_{r}(e, x) H_{r}(x, x a) H_{r}(x a, x a y) \\
& \leqslant C H_{r}(e, x a y) .
\end{aligned}
$$

Let us estimate the number of preimages under $\Psi$ of some point $z$. Let $\gamma_{z}$ be a geodesic segment from $e$ to $z$. If $z=x a y$ and $x$ is far away from $\gamma_{z}$, a tree approximation shows that $|z|$ is significantly smaller than $|x|+|a|+|y|$. This is impossible by construction. Therefore, $x$ is contained in a ball of fixed radius $B\left(\gamma_{z}(k), C\right)$. In particular, the number of possibilities for $x$ is uniformly bounded. Arguing in the same way for $y$, we deduce that, for some $C>0$, each point has at most $C$ preimages under $\Psi$.

Finally,

$$
\begin{aligned}
u_{k}(r) u_{\ell}(r) & =\sum_{x \in \mathbb{S}_{k}, y \in \mathbb{S}_{\ell}} H_{r}(e, x) H_{r}(e, y) \leqslant C \sum_{x \in \mathbb{S}_{k}, y \in \mathbb{S}_{\ell}} H_{r}(e, \Psi(x, y)) \\
& \leqslant C \sum_{i=k+\ell}^{k+\ell+C} \sum_{z \in \mathbb{S}_{i}} H_{r}(e, z) \leqslant C \sum_{i=k+\ell}^{k+\ell+C} u_{i}(r) .
\end{aligned}
$$

As $r<R$, the sum $\sum_{x \in \Gamma} H_{r}(e, x)$ is finite by (2.6). In particular, the sequence $u_{k}(r)$ is summable, and reaches its maximum $M(r)$ at some index $k_{0}(r)$. Using the previous equation with $k=\ell=k_{0}(r)$, we get $M(r)^{2} \leqslant C(C+1) M(r)$; hence $M(r) \leqslant C(C+1)=D$.

Finally, for every $r<R$, for every $k \in \mathbb{N}$, one has $\sum_{x \in \mathbb{S}_{k}} H_{r}(e, x) \leqslant D$. The lemma follows by letting $r$ tend to $R$.

The following lemma is the main estimate in the proof of Theorem 2.3. It gives superexponentially small estimates for the $R$-probabilities of paths staying too far away from geodesics, implying that such paths are very unlikely and will not contribute a lot to $G_{R}$.
Lemma 2.6. Assume that $\mu$ is finitely supported and symmetric. There exist $n_{0}>0$ and $\varepsilon>0$ such that, for any $n \geqslant n_{0}$, for any $x, y, z \in \Gamma$ on a geodesic segment (in this order) with $d(x, y) \geqslant n$ and $d(y, z) \geqslant n$,

$$
G_{R}\left(x, z ; B(y, n)^{c}\right) \leqslant 2^{-e^{\varepsilon n}}
$$

Proof. Without loss of generality, one can assume $y=e$.
Fix some $\varepsilon>0$ very small, and let $N=\left\lfloor e^{\varepsilon n}\right\rfloor$. In this proof, we will write $C$ for a generic constant independent of $\varepsilon$. The idea of the proof is to construct $N$ barriers $A_{1}, \ldots, A_{N}$ such that any trajectory of the random walk going from $x$ to $z$ outside of $B(e, n)$ has to go through $A_{1}$, then $A_{2}$, and so on. Decomposing a trajectory according to its first visit to $A_{1}$, then $A_{2}$, and so on, we obtain as in (2.1)

$$
\begin{align*}
& G_{R}\left(x, z ; B(e, n)^{c}\right)  \tag{2.8}\\
& \quad \leqslant \sum_{a_{1} \in A_{1}} \cdots \sum_{a_{N} \in A_{N}} G_{R}\left(x, a_{1}\right) G_{R}\left(a_{1}, a_{2}\right) \cdots G_{R}\left(a_{N-1}, a_{N}\right) G_{R}\left(a_{N}, z\right)
\end{align*}
$$

We will construct the barriers so that, writing $A_{0}=\{x\}$ and $A_{N+1}=\{z\}$, one has for any $0 \leqslant i \leqslant N$

$$
\begin{equation*}
\sum_{a \in A_{i}} \sum_{b \in A_{i+1}} G_{R}(a, b)^{2} \leqslant 1 / 4 \tag{2.9}
\end{equation*}
$$

This implies the desired estimate on $G_{R}\left(x, z ; B(y, n)^{c}\right)$ by Cauchy-Schwarz, as follows. To write it formally, it is more convenient to express things in terms of operators, as in Ser83 Led13. Define an operator $L_{i}: \ell^{2}\left(A_{i+1}\right) \rightarrow \ell^{2}\left(A_{i}\right)$ by $L_{i} f(a)=\sum_{b \in A_{i+1}} G_{R}(a, b) f(b)$. The sum to estimate in (2.8) is $\left(L_{0} \cdots L_{N} \delta_{z}\right)(x) ;$ it is therefore bounded by $\Pi\left\|L_{i}\right\|$. Moreover,

$$
\begin{aligned}
\left\|L_{i} f\right\|_{\ell^{2}}^{2} & =\sum_{a \in A_{i}}\left|\sum_{b \in A_{i+1}} G_{R}(a, b) f(b)\right|^{2} \leqslant \sum_{a \in A_{i}}\left(\sum_{b \in A_{i+1}} G_{R}(a, b)^{2}\right)\left(\sum_{b \in A_{i+1}}|f(b)|^{2}\right) \\
& =\left(\sum_{a \in A_{i}} \sum_{b \in A_{i+1}} G_{R}(a, b)^{2}\right)\|f\|_{\ell^{2}}^{2} .
\end{aligned}
$$

With (2.9), we obtain $\left\|L_{i}\right\| \leqslant\left(\sum_{a \in A_{i}} \sum_{b \in A_{i+1}} G_{R}(a, b)^{2}\right)^{1 / 2} \leqslant 1 / 2$, and the result of the lemma follows.

It remains to construct barriers satisfying (2.9). The construction is geometric and is done in the hyperbolic space $\mathbb{H}^{m}$ for some $m \geqslant 2$ (or rather its model as the euclidean unit ball in $\mathbb{R}^{m}$, with the boundary at infinity identified with the unit
sphere $S^{m-1}$ in $\mathbb{R}^{m}$ ). By $\overline{\mathrm{BSO0}}$, the group $\Gamma$ with its word metric is roughly similar to a subset of such a space: if $m$ is large enough, there exist a mapping $\Psi: \Gamma \rightarrow \mathbb{H}^{m}$ and $\lambda>0, C>0$ such that $\left|\lambda d_{\mathbb{H}}(\Psi(u), \Psi(v))-d(u, v)\right| \leqslant C$ for all $u, v \in \Gamma$. The image under $\Psi$ of a geodesic in $\Gamma$ is a quasigeodesic in $\mathbb{H}^{m}$; therefore it remains uniformly close to a true hyperbolic geodesic (see for instance GdlH90, Theorem 5.11]). It follows that it does not make a serious difference to use geodesics in $\Gamma$ or in $\mathbb{H}^{m}$.

The hyperbolic geodesic from $\Psi(x)$ to $\Psi(z)$ can be extended bi-infinitely. Composing with a hyperbolic isometry, we can assume that this geodesic goes through the center $O$ of the ball model of $\mathbb{H}^{m}$, and that $\Psi(e)$ is a bounded distance away from $O$. Let $\xi$ be the endpoint of this hyperbolic geodesic in negative time. To an angle $\theta \in[0, \pi]$, we associate the union of all the semi-infinite geodesics $[O \zeta)$ (with $\zeta \in S^{m-1}$ ) making an angle $\theta$ with $[O \xi$ ) (its boundary at infinity is the circle of points at distance $\theta$ of $\xi$ in $S^{m-1}$ ). Let then $A(\theta)$ be the set of points $a$ in $B(e, n)^{c} \subset \Gamma$ such that $\Psi(a)$ is at a distance at most $C_{0}$ of such a geodesic. If $C_{0}$ is chosen large enough, a path of the random walk going from $x$ to $z$ in $B(e, n)^{c}$ can not jump over $A(\theta)$ since $\mu$ has finite support, so that $A(\theta)$ is a barrier.

In $X=[0, \pi]$, consider $X_{i}=[(2 i-1) / N, 2 i / N]$ (for $1 \leqslant i \leqslant N$ ). Those intervals are separated by $1 / N \sim e^{-\varepsilon n}$. In each of them, we will choose an angle $\theta_{i}$ and let $A_{i}=A\left(\theta_{i}\right)$. One should then ensure that (2.9) is satisfied. To do so, we will choose each $\theta_{i}$ at random as follows. Let $\Omega=\prod_{i=1}^{N} X_{i}$, endowed with the product of the probability measures $N \mathrm{dLeb}$ on $X_{i}$. Define a function $f_{i}$ on $\Omega$ by

$$
f_{i}\left(\theta_{1}, \ldots, \theta_{N}\right)=\sum_{a \in A\left(\theta_{i}\right), b \in A\left(\theta_{i+1}\right)} G_{R}(a, b)^{2}
$$

where by convention $A\left(\theta_{0}\right)=\{x\}$ and $A\left(\theta_{N+1}\right)=\{z\}$. One should find a value of $\bar{\theta}=\left(\theta_{1}, \ldots, \theta_{N}\right)$ such that $f_{i}(\bar{\theta}) \leqslant 1 / 4$ for all $i$. We will show that

$$
\begin{equation*}
\int f_{i} \leqslant C e^{-\rho n} \tag{2.10}
\end{equation*}
$$

for some $\rho>0$ independent of $\varepsilon$. It follows that $\int\left(\sum f_{i}\right) \leqslant C(1+N) e^{-\rho n} \leqslant C(1+$ $\left.e^{\varepsilon n}\right) e^{-\rho n}$. Choosing $\varepsilon$ small enough, this is exponentially small, and is in particular bounded by $1 / 4$ for large enough $n$. This yields a point $\bar{\theta}$ with $\sum f_{i}(\bar{\theta}) \leqslant 1 / 4$, for which the corresponding barriers satisfy (2.9).

Let us now prove (2.10). We will only give the argument for $1 \leqslant i \leqslant N-1$ : the case of $f_{0}$ and $f_{N}$ is slightly different (since $A_{0}=\{x\}$ and $A_{N+1}=\{z\}$ are fixed), and it turns out to be analogous to the general case, but simpler. Fix some $i \in[1, N-1]$. To each $a \in \Gamma$ and $j \in\{i, i+1\}$, we associate the set $X_{j}(a)$ of angles $\theta$ in $X_{j}$ such that $a \in A(\theta)$. By definition,

$$
\int f_{i}=\sum_{a, b \in \Gamma} G_{R}(a, b)^{2} \cdot N \operatorname{Leb}\left(X_{i}(a)\right) \cdot N \operatorname{Leb}\left(X_{i+1}(b)\right) .
$$

For $a \in \mathbb{S}_{k}$, its image under $\Psi$ is at a distance at least $\alpha k$ of $O$ in $\mathbb{H}^{m}$, for some $\alpha>0$. If one moves away from this point by at most $C_{0}$, the visual angle from $O$ varies by at most $C e^{-\alpha k}$. It follows that $\operatorname{Leb}\left(X_{j}(a)\right) \leqslant C e^{-\alpha k}$. Since $N \leqslant e^{\varepsilon n}$, we obtain

$$
\int f_{i} \leqslant C e^{2 \varepsilon n} \sum_{a, b} G_{R}(a, b)^{2} e^{-\alpha|a|} e^{-\alpha|b|}
$$

where the sum is restricted to those $a$ and $b$ outside of $B(e, n)$ and whose images under $\Psi$ belong to the $C_{0}$-neighborhoods of the sectors delimited, respectively, by $X_{i}$ and $X_{i+1}$. Writing $u=a^{-1} b$ (with $|u| \leqslant|a|+|b|$ ), we get

$$
\int f_{i} \leqslant C e^{2 \varepsilon n} \sum_{u} G_{R}(e, u)^{2} e^{-\alpha|u|} M(u)
$$

where $M(u)$ is the number of ways to decompose $u$ as $a^{-1} b$. Fix a point $u$, and such a decomposition $u=a^{-1} b$.

The hyperbolic geodesics from $O$ to, respectively, $\Psi(a)$ and $\Psi(b)$, make an angle at least $e^{-\varepsilon n} / 2$. Therefore, they are far away from each other outside of the ball $B(O, 2 \varepsilon n)$. It follows from a tree approximation that

$$
d_{\mathbb{H}}(\Psi(a), \Psi(b)) \geqslant|\Psi(a)|+|\Psi(b)|-4 \varepsilon n-C .
$$

Since $\Psi$ is a quasisimilarity, we deduce that $|u|=d(a, b) \geqslant|a|+|b|-C \varepsilon n-C$. In particular, if $\varepsilon$ is small enough, since $|a| \geqslant n$ and $|b| \geqslant n$, we obtain $|u| \geqslant n$. It also follows from this argument that a geodesic in the group from $a$ to $b$ has to pass through the ball $B(e, C \varepsilon n)$, since geodesics in the group and in hyperbolic space remain a bounded distance away. Let $\gamma$ be a geodesic segment from $e$ to $u$ in $\Gamma$. Then $a \gamma$ is a geodesic segment from $a$ to $b$. There exists a time $j$ such that $a \gamma(j) \in$ $B(e, C \varepsilon n)$. Finally, $a \in \bigcup_{j=0}^{|u|} \gamma(j)^{-1} B(e, C \varepsilon n)$, which gives at most $(|u|+1) C^{C \varepsilon n}$ possibilities for $a$. Arguing similarly for $b$, we obtain $M(u) \leqslant(|u|+1)^{2} e^{C \varepsilon n}$ for some $C>0$.

Finally, we have

$$
\begin{aligned}
\int f_{i} & \leqslant C e^{2 \varepsilon n} \sum_{|u| \geqslant n} G_{R}(e, u)^{2} e^{-\alpha|u|}(|u|+1)^{2} e^{C \varepsilon n} \\
& \leqslant C e^{(C+2) \varepsilon n} \sum_{|u| \geqslant n} G_{R}(e, u)^{2} e^{-\alpha|u| / 2}
\end{aligned}
$$

Since $\sum_{|u|=k} G_{R}(e, u)^{2}$ equals $\sum_{|u|=k} H_{R}(e, u)$ by symmetry of $\mu$, it is uniformly bounded by Lemma 2.5. Therefore, $\int f_{i}$ is bounded by $C e^{C \varepsilon n} e^{-\alpha n / 2}$. If $\varepsilon$ is small enough, this is at most $C e^{-\alpha n / 4}$. This proves (2.10) and concludes the proof of the lemma.

The following lemma is proved in GL13 and is elementary (see the proof of Theorems 4.1 and 4.3 there).
Lemma 2.7. Let $\mu$ be an admissible measure on a Gromov-hyperbolic group. Assume that, for all $K>0$, there exists $n_{0}$ such that, for all $n \geqslant n_{0}$, for all points $x, y, z$ on a geodesic segment (in this order) with $d(x, y) \in[n, 100 n]$ and $d(y, z) \in[n, 100 n]$, one has $G_{R}\left(x, z ; B(y, n)^{c}\right) \leqslant K^{-n}$. Then $\mu$ satisfies uniform Ancona inequalities. If it has finite support, it even satisfies strong uniform Ancona inequalities (as defined below in Definition [2.8]).

To prove this lemma, one uses recursively its assumptions to show that most $r$-weight is concentrated on paths staying close enough to the geodesic from $x$ to $z$, and in particular passing in a ball of fixed radius around $y$. This lemma, together with Lemma 2.6 proves uniform Ancona inequalities for symmetric measures, i.e., Theorem [2.3. For finitely supported measures, strong uniform Ancona inequalities (see below) are then deduced as in [GL13. Theorem 4.6].

We need the following strengthening of Ancona inequalities.
Definition 2.8. A measure $\mu$ on a Gromov-hyperbolic group satisfies strong uniform Ancona inequalities if it satisfies uniform Ancona inequalities and, additionally, there exist constants $C>0$ and $\rho>0$ such that, for all points $x, x^{\prime}, y, y^{\prime}$ whose configuration is approximated by a tree as follows:

for any $r \in[1, R]$,

$$
\begin{equation*}
\left|\frac{G_{r}(x, y) / G_{r}\left(x^{\prime}, y\right)}{G_{r}\left(x, y^{\prime}\right) / G_{r}\left(x^{\prime}, y^{\prime}\right)}-1\right| \leqslant C e^{-\rho n} \tag{2.11}
\end{equation*}
$$

Ancona inequalities ensure that the quantity $\frac{G_{r}(x, y) / G_{r}\left(x^{\prime}, y\right)}{G_{r}\left(x, y^{\prime}\right) / G_{r}\left(x^{\prime}, y^{\prime}\right)}$ in the definition is bounded from above and from below. Strong Ancona inequalities ensure that this quantity is exponentially close to 1 in terms of the distance between the sets of points $\left\{x, x^{\prime}\right\}$ and $\left\{y, y^{\prime}\right\}$. These bounds are not formal consequences of Ancona inequalities, but they are consequences of Ancona inequalities in suitable domains (that follow from Lemma 2.6). Applying Lemmas 2.6 and 2.7 we obtain the following result, strengthening Theorem 2.3
Theorem 2.9. If $\mu$ is admissible, finitely supported, and symmetric on a nonelementary Gromov-hyperbolic group, it satisfies strong uniform Ancona inequalities.

When one takes $x^{\prime}=e$, then the quantity appearing in (2.11) is the ratio $K_{r, y}(x) / K_{r, y^{\prime}}(x)$ of the Martin kernels. The theorem implies that, when $y_{i}$ tends to a point $\xi \in \partial \Gamma$, the sequence $K_{r, y_{i}}(x)$ is a Cauchy sequence (since the points $x, e$ and $y_{i}, y_{j}$ satisfy the assumptions of the definition with a large $n$ for large enough $i, j)$. Hence, it converges to a function $K_{r, \xi}(x)$. This is the main step in the proof that the Martin boundary for $r$-harmonic functions coincides with the geometric boundary (one should also check that $K_{r, \xi} \neq K_{r, \eta}$ for $\xi \neq \eta$, which is easy). We omit the (classical) details; see, for instance, INO08.

## 3. Asymptotics of the Green function

Let $\Gamma$ be a nonelementary Gromov-hyperbolic group. Our goal in this section is to prove the following theorem.
Theorem 3.1. Let $\mu$ be an admissible probability measure on $\Gamma$ satisfying strong uniform Ancona inequalities. For any $x, y \in \Gamma$, there exists $C(x, y)>0$ such that, when $r$ tends to $R=R(\mu)$,

$$
\partial G_{r}(x, y) / \partial r \sim C(x, y)(R-r)^{-1 / 2}
$$

Throughout this section, we fix a measure $\mu$ satisfying the assumptions of this theorem. Theorem 2.9 shows that it is the case for finitely supported symmetric measures, but symmetry or finite support will play no additional role in this section. We will concentrate mainly on the proof of Theorem 3.1 for $x=y=e$, since the general case will follow easily.

In a sense, the proof of Theorem 3.1 is essentially done in [GL13, but there is an important technical difference: a (well-chosen) Markov automaton for a surface group is transitive, while there can be several components in a general hyperbolic
group. This means that, in the thermodynamic formalism, we will have to deal with several dominating components. This problem is solved thanks to a technique of Calegari and Fujiwara CF10 and to Lemma 2.5. This sketch of the argument might be sufficient for experts, but since there are several technical subtleties, we will give most details below. There is a significant overlap with some arguments in GL13, but this seems necessary to keep the argument understandable. Two significant differences with [GL13] (in addition to the existence of several dominating components in the automaton, and directly related to this issue) are that we need some a priori estimates (proved in Subsection 3.1), and that for $r<R$ we will associate to $G_{r}$ a measure living on the group, not on the boundary.
3.1. A priori estimates. The main idea behind the proof of Theorem 3.1] as in GL13, is that the function $G_{r}(e, e)$ almost satisfies a differential equation. By (2.5), its derivative with respect to $r$ is essentially $\sum_{x} G_{r}(e, x) G_{r}(x, e)=$ $\sum_{x} H_{r}(e, x)$ [recall the notation (2.7)], and its second derivative is essentially $\sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)$. To prove that $G_{r}(e, e)$ almost satisfies a differential equation, we should relate those quantities. The next proposition gives such a (crude) relation.

Proposition 3.2. There exists $C>0$ such that, for all $r \in[1, R)$,

$$
C^{-1} \leqslant \frac{\sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)}{\left(\sum_{x} G_{r}(e, x) G_{r}(x, e)\right)^{3}} \leqslant C
$$

Proof. Consider two points $x, y$. The triangle with vertices $e, x, y$ is thin, so there exists a point $w$ (defined uniquely up to a finite set) that is close to each of its sides. By the Ancona inequality, we have $G_{r}(e, y)=G_{r}\left(w^{-1}, w^{-1} y\right) \leqslant$ $C G_{r}\left(w^{-1}, e\right) G_{r}\left(e, w^{-1} y\right)$, since a geodesic segment from $w^{-1}$ to $w^{-1} y$ passes close to $e$ by construction. Similar estimates hold along $[y, x]$ and $[x, e]$, and we obtain

$$
\begin{aligned}
& G_{r}(e, y) G_{r}(y, x) G_{r}(x, e) \\
& \quad \leqslant C G_{r}\left(w^{-1}, e\right) G_{r}\left(e, w^{-1}\right) \cdot G_{r}\left(w^{-1} x, e\right) G_{r}\left(e, w^{-1} x\right) \cdot G_{r}\left(w^{-1} y, e\right) G_{r}\left(e, w^{-1} y\right)
\end{aligned}
$$

The points $w^{-1}, w^{-1} x$, and $w^{-1} y$ determine $x$ and $y$. Using the notation $H_{r}$ and summing over $x$ and $y$, we get

$$
\sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e) \leqslant C \sum_{a, b, c} H_{r}(e, a) H_{r}(e, b) H_{r}(e, c)
$$

This is one of the inequalities of the proposition.
For the reverse inequality, for any $u \in \Gamma$, write $B(u)$ for the set of points $x$ such that a geodesic from $e$ to $x$ passes close to $u$. Lemma 2.4 ensures that, for any $z \in \Gamma$, there exists a uniformly bounded $a$ such that $u a z \in B(u)$. Harnack inequalities (2.2) give $G_{r}(e, z) \leqslant C(u) G_{r}(e, u a z)$ (and $H_{r}$ satisfies the same inequality). Hence, $\sum_{z \in \Gamma} H_{r}(e, z) \leqslant C(u) \sum_{v \in B(u)} H_{r}(e, v)$. Choose now three geodesic segments $\gamma_{1}$, $\gamma_{2}$, and $\gamma_{3}$ (with endpoints denoted by $u_{1}, u_{2}, u_{3}$ ), long enough and going in three different directions (this is possible since the group is nonelementary) so that the sets $B\left(u_{i}\right)$ are pairwise disjoint, and so that any geodesic from $B\left(u_{i}\right)$ to $B\left(u_{j}\right)$ $(i \neq j)$ has to pass close to $e$. We get

$$
\left(\sum_{z} H_{r}(e, z)\right)^{3} \leqslant C\left(u_{1}, u_{2}, u_{3}\right) \sum_{v_{i} \in B\left(u_{i}\right)} H_{r}\left(e, v_{1}\right) H_{r}\left(e, v_{2}\right) H_{r}\left(e, v_{3}\right) .
$$

By (2.4), we have $G_{r}\left(v_{1}, e\right) G_{r}\left(e, v_{2}\right) \leqslant C G_{r}\left(v_{1}, v_{2}\right)$, and similarly for circular permutations. This sum is therefore bounded by $\sum G_{r}\left(v_{1}, v_{2}\right) G_{r}\left(v_{2}, v_{3}\right) G_{r}\left(v_{3}, v_{1}\right)$. Let $y=v_{1}^{-1} v_{2}$ and $x=v_{1}^{-1} v_{3}$. The point close to the three sides of a geodesic triangle with vertices $e, x$, and $y$ is close to $v_{1}^{-1}$ by construction; hence $x$ and $y$ determine $v_{1}$ (and then $v_{2}$ and $v_{3}$ ) up to a finite number of possibilities. We finally get

$$
\left(\sum_{z} H_{r}(e, z)\right)^{3} \leqslant C \sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)
$$

proving the other inequality of the lemma.
Corollary 3.3. There exist $A \geqslant 0$ and $C>0$ such that, for all $r \in[1, R)$,

$$
\frac{C^{-1}}{\sqrt{A+(R-r)}} \leqslant \sum_{x \in \Gamma} H_{r}(e, x) \leqslant \frac{C}{\sqrt{A+(R-r)}} .
$$

Moreover, $A=0$ if and only if $\sum_{x \in \Gamma} H_{R}(e, x)=+\infty$.
Proof. Let $\eta(r)=\sum_{x} H_{r}(e, x)$, and $F(r)=r^{2} \eta(r)$. By (2.5),

$$
F^{\prime}(r)=2 r \sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e) .
$$

Therefore, Proposition 3.2 shows that $2 F^{\prime}(r) / F(r)^{3}=\left(-1 / F(r)^{2}\right)^{\prime}$ is bounded from above and below. Integrating this estimate on an interval $[r, s$ ], we get

$$
C^{-1}(s-r) \leqslant 1 / F(r)^{2}-1 / F(s)^{2} \leqslant C(s-r) .
$$

When $s$ increases to $R, F(s)$ converges either to a positive constant or to infinity. Hence, $1 / F(s)^{2}$ converges to $A \in[0, \infty)$. We get $A+C^{-1}(R-r) \leqslant 1 / F(r)^{2} \leqslant$ $A+C(R-r)$, from which the corollary follows.

We shall see later that $A$ is in fact equal to 0 . Therefore, this corollary gives the right order of magnitude $1 / \sqrt{R-r}$ for the function $\eta(r)=\sum_{x \in \Gamma} H_{r}(e, x)$. However, to obtain Theorem 3.1. we need to get asymptotics, of the form $\eta(r) \sim C / \sqrt{R-r}$. The strategy will be the same, relying on the differential equation, but we will need to improve Proposition 3.2, to get convergence instead of mere bounds. This is most conveniently done using the transfer operator on a Markov automaton, as we will explain in the next subsection. Before doing this, let us state a final technical lemma that relies on Corollary 3.3 and will be important later on.

Lemma 3.4. Fix $a \in \Gamma$. There exists $C>0$ such that, for any $x \in \Gamma$, for any $r \in[1, R]$,

$$
\left|\log \left(\frac{G_{r}(e, x)}{G_{r}(a, x)}\right)-\log \left(\frac{G_{R}(e, x)}{G_{R}(a, x)}\right)\right| \leqslant C \sqrt{R-r}
$$

and

$$
\left|\log \left(\frac{G_{r}(x, e)}{G_{r}(x, a)}\right)-\log \left(\frac{G_{R}(x, e)}{G_{R}(x, a)}\right)\right| \leqslant C \sqrt{R-r}
$$

Proof. The second estimate of the lemma can be deduced from the first one applied to the measure $\check{\mu}(g)=\mu\left(g^{-1}\right)$. We will therefore concentrate on the first one.

Fix some $x \in \Gamma$. Let $f(r)=\log \left(G_{r}(e, x) / G_{r}(a, x)\right)$. We will show that its derivative is bounded in absolute value by $C / \sqrt{A+R-r}$, where $C$ is a constant that does not depend on $x$ (of course, it may depend on $a$ ). By integration, this
gives $|f(r)-f(R)| \leqslant C(\sqrt{A+R-r}-\sqrt{A})$, which is bounded by $C \sqrt{R-r}$ as desired.

We write $f(r)=\log \left(r G_{r}(e, x)\right)-\log \left(r G_{r}(a, x)\right)$. With the formula (2.5) for the derivative of $r G_{r}$, we get

$$
\begin{aligned}
f^{\prime}(r) & =\frac{\sum_{y} G_{r}(e, y) G_{r}(y, x)}{r G_{r}(e, x)}-\frac{\sum_{y} G_{r}(a, y) G_{r}(y, x)}{r G_{r}(a, x)} \\
& =\frac{1}{r} \sum_{y}\left(1-\frac{G_{r}(a, y) / G_{r}(e, y)}{G_{r}(a, x) / G_{r}(e, x)}\right) \frac{G_{r}(e, y) G_{r}(y, x)}{G_{r}(e, x)} .
\end{aligned}
$$

Consider a geodesic segment $\gamma$ from $e$ to $x$ and write $\gamma(n)(0 \leqslant n \leqslant|x|)$ for the point on $\gamma$ at distance $n$ of $e$. Let $\Gamma_{n}$ denote the set of points $y \in \Gamma$ whose projection on $\gamma$ is $\gamma(n)$, i.e., $d(y, \gamma(n)) \leqslant d(y, \gamma(i))$ for $i \neq n$. Note that there can be several such projections - in the following argument, the multiplicity is not important; otherwise one can avoid it by using only the first projection. For $y \in \Gamma_{n}$, the points $e, a$ and $x, y$ are in the configuration of Definition 2.8, with a separating distance at least $n-C$ (for some $C$ depending only on $a$ ). Since $\mu$ satisfies strong uniform Ancona inequalities by assumption, we obtain

$$
\left|f^{\prime}(r)\right| \leqslant C \sum_{n=0}^{|x|} \sum_{y \in \Gamma_{n}} e^{-\rho n} \frac{G_{r}(e, y) G_{r}(y, x)}{G_{r}(e, x)}
$$

For $y \in \Gamma_{n}$, geodesics from $e$ to $y$ and from $y$ to $x$ pass close to $\gamma(n)$. Hence, Ancona inequalities give

$$
\begin{aligned}
G_{r}(e, y) G_{r}(y, x) & \leqslant C G_{r}(e, \gamma(n)) G_{r}(\gamma(n), y) G_{r}(y, \gamma(n)) G_{r}(\gamma(n), x) \\
& \leqslant C G_{r}(e, x) H_{r}(\gamma(n), y)
\end{aligned}
$$

Finally,

$$
\left|f^{\prime}(r)\right| \leqslant C \sum_{n=0}^{|x|} \sum_{y \in \Gamma_{n}} e^{-\rho n} H_{r}(\gamma(n), y) \leqslant C \sum_{n=0}^{|x|} e^{-\rho n} \sum_{y \in \Gamma} H_{r}\left(e, \gamma(n)^{-1} y\right)
$$

Since $\sum_{z \in \Gamma} H_{r}(e, z) \leqslant C / \sqrt{A+R-r}$ by Corollary 3.3 and $e^{-\rho n}$ is summable, this proves the lemma.
3.2. Symbolic dynamics. For a nice introduction to the topics of this paragraph and the next one, see CF10.

Let $S$ be a finite symmetric generating set of the group $\Gamma$. A rooted $S$-labeled automaton (or simply automaton) is a finite directed graph $\mathcal{A}=\left(V, E, s_{*}\right)$ with distinguished vertex $s_{*}$ ("start"), and a labeling $\alpha: E \rightarrow S$ of edges by generators of the group.

A path in the graph is a sequence of edges $e_{0}, \ldots, e_{m-1}$ such that the endpoint of $e_{i}$ is the starting point of $e_{i+1}$. To such a path $\gamma$, one can associate a path $\alpha(\gamma)$ in the Cayley graph of $\Gamma$ by multiplying successively the generators read along the edges of the path. Let $\alpha_{*}(\gamma)$ be the endpoint of $\alpha(\gamma)$.

Definition 3.5. An automaton is a strongly Markov automatic structure for $\Gamma$ if
(1) Every vertex $v \in V$ is accessible from the start state $s_{*}$.
(2) For every path $\gamma$, the path $\alpha(\gamma)$ is a geodesic path in $\Gamma$.
(3) The endpoint mapping $\alpha_{*}$ induced by $\alpha$ is a bijection of the set of paths starting at $s_{*}$ onto $\Gamma$.

In particular, the sphere $\mathbb{S}_{k}$ of $\Gamma$ is in bijection with the set of paths of length $k$ starting from $s_{*}$.

Every Gromov-hyperbolic group admits such a strongly Markov automatic structure, by a theorem of Cannon Can84. Let us fix once and for all such an automaton. An infinite path in the graph $\mathcal{A}$ determines a semi-infinite geodesic in the group starting from $e$, and therefore a point on the boundary at infinity. In this way, we extend $\alpha_{*}$ to a map from infinite paths to $\partial \Gamma$.

We will denote by $\Sigma^{*}$ the set of finite paths in the graph, by $\Sigma$ the set of semiinfinite paths, and $\bar{\Sigma}=\Sigma^{*} \cup \Sigma$. These sets are endowed with a metric $d\left(\omega, \omega^{\prime}\right)=2^{-n}$ where $n$ is the first time the paths $\omega$ and $\omega^{\prime}$ differ. With this metric, $\Sigma^{*}$ is a dense open subset of the compact space $\bar{\Sigma}$. The map $\alpha_{*}$ is continuous from $\bar{\Sigma}$ to $\Gamma \cup \partial \Gamma$.

Denote by $\mathcal{H}^{\beta}$ the space of $\beta$-Hölder continuous functions on $\bar{\Sigma}$. For $0<\beta<\beta^{\prime}$, one has the following basic inequality (which is true in any metric space):

$$
\begin{equation*}
\|f\|_{\mathcal{H}^{\beta}} \leqslant 2\|f\|_{C^{0}}^{1-\beta / \beta^{\prime}}\|f\|_{\mathcal{H}^{\beta^{\prime}}}^{\beta / \beta^{\prime}} \tag{3.1}
\end{equation*}
$$

In particular, if a sequence of functions $f_{n}$ converges in $C^{0}$ and remains bounded in $\mathcal{H}^{\beta^{\prime}}$, then it converges in $\mathcal{H}^{\beta}$.

Note that a Hölder-continuous function on $\Sigma^{*}$ uniquely extends to a Höldercontinuous function on $\bar{\Sigma}$. Finally, let $\sigma: \bar{\Sigma} \rightarrow \bar{\Sigma}$ be the left shift, forgetting the first edge of a path.

When a statement is true in all Hölder spaces $\mathcal{H}^{\beta}$ for small enough $\beta$, we will simply say that it is true in Hölder spaces, without specifying $\beta$. We will also write $\|f\|$ for $\|f\|_{\mathcal{H}^{\beta}}$, unless the specific value of $\beta$ is important.

The connectedness properties of the graph $\mathcal{A}$ will play an important role below. A component of the automaton is a maximal subset of the set $V$ of vertices in which any vertex can be reached from any other vertex by a finite-length path. The simplest case is when there is a single component of the graph with cardinality bigger than 1: all nontrivial behavior takes place in this component. This is the case for well-chosen automata for subgroups of PSL $(2, \mathbb{R})$, but not for general hyperbolic groups. Identifying points belonging to the same component, one obtains a new directed graph, the components graph, in which there is no loop. This graph encodes how different components interact.
3.3. Peripheral spectrum of transfer operators. Since the spectral description of transfer operators is very classical, we will only sketch the proofs in this section, referring to [PP90] for more details.

Consider a finite directed graph $\mathcal{A}$, let $\bar{\Sigma}$ be the set of finite or infinite paths in $\mathcal{A}$, and let $\sigma$ be the left shift. (If one is uncomfortable with the idea of considering finite paths in the graph, one can equivalently add a cemetery to the graph, which can be reached from any vertex, and extend a finite path by infinitely many steps in the cemetery.) To any real-valued Hölder-continuous function $\varphi: \bar{\Sigma} \rightarrow \mathbb{R}$ (called a potential), one associates the so-called transfer operator $\mathcal{L}_{\varphi}$, defined on the set of Hölder-continuous functions by

$$
\mathcal{L}_{\varphi} f(\omega)=\sum_{\sigma\left(\omega^{\prime}\right)=\omega} e^{\varphi\left(\omega^{\prime}\right)} f\left(\omega^{\prime}\right)
$$

where for $\omega=\emptyset$ the empty path we consider only the nonempty preimages of $\omega$. The iterates of this operator encode a lot of information on the Birkhoff sums
$S_{n} \varphi(\omega)=\sum_{j=0}^{n-1} \varphi\left(\sigma^{j} \omega\right)$ of the potential $\varphi$. For instance, one has

$$
\mathcal{L}_{\varphi}^{n} 1(\emptyset)=\sum e^{S_{n} \varphi(\omega)},
$$

where the sum is over all paths of length $n$.
In the case of hyperbolic groups, we will be interested in the asymptotics of such sums, since for suitable potentials $\varphi_{r}$ they correspond to the sum of $H_{r}$ over the sphere of radius $n$ in $\Gamma$ (this is one of the quantities we want to estimate precisely to improve on Proposition 3.2). Such asymptotics can be read from the spectrum of $\mathcal{L}_{\varphi}$ that we now describe.

The simplest situation is when the graph is topologically mixing-i.e., one can go from any vertex to any other vertex (one says that the graph is recurrent) and for any $a, b \in \mathcal{A}$, for any large enough $n$, there is a path of length exactly $n$ from $a$ to $b$. In this case, the spectral description of $\mathcal{L}_{\varphi}$ is very simple and is given by the following theorem (called the Ruelle-Perron-Frobenius theorem).
Theorem 3.6. Assume $\mathcal{A}$ is topologically mixing. The operator $\mathcal{L}_{\varphi}$ acting on the space of Hölder-continuous functions has a unique eigenvalue of maximal modulus denoted by $e^{\operatorname{Pr}(\varphi)}$, and the rest of its spectrum is contained in a disk of strictly smaller radius. Moreover, the corresponding eigenfunction $h$ (suitably normalized) is strictly positive everywhere, and the eigenprojector is given by $\Pi f=\left(\int f \mathrm{~d} \lambda\right) h$ for some probability measure $\lambda$ whose support is the set $\Sigma$ of infinite paths. Finally, the probability measure $h \mathrm{~d} \lambda$ is invariant under $\sigma$ and ergodic.

In other words, one has

$$
\left\|\mathcal{L}_{\varphi}^{n} f-e^{n \operatorname{Pr}(\varphi)}\left(\int f \mathrm{~d} \lambda\right) h\right\| \leqslant C\|f\| e^{-n \varepsilon} e^{n \operatorname{Pr}(\varphi)}
$$

for some $C>0$ and $\varepsilon>0$. This is Theorem 2.2 in [PP90] (the statement there is only given on $\Sigma$, but the proofs readily adapt to $\bar{\Sigma}$ ). The real number $\operatorname{Pr}(\varphi)$ is called the pressure of the potential $\varphi$.

Assume now that $\mathcal{A}$ is recurrent, but not mixing: there is a minimal period $p>1$ such that any path from a vertex to itself has length $n p$ for some integer $n$. In this case, the set $V$ of vertices of $\mathcal{A}$ is a disjoint union $\bigsqcup_{j=0}^{p-1} V_{j}$, where for any $j \in \mathbb{Z} / p \mathbb{Z}$ an outgoing edge of $V_{j}$ is an ingoing edge of $V_{j+1}$ (we call this decomposition a cyclic decomposition of $V$ ). Denoting by $\bar{\Sigma}_{j}$ the set of paths beginning from a vertex in $V_{j}$ and the empty path, then $\sigma$ maps $\bar{\Sigma}_{j}$ to $\bar{\Sigma}_{j+1}$. Moreover, the restriction of $\sigma^{p}$ to any $\bar{\Sigma}_{j}$ is a topologically mixing subshift of finite type, to which Theorem 3.6 applies. This readily implies that the eigenvalues of maximal modulus of $\mathcal{L}_{\varphi}$ are of the form $e^{2 \mathrm{i} k \pi / p} e^{\operatorname{Pr}(\varphi)}$ for some real number $\operatorname{Pr}(\varphi)$, they are all simple, and the rest of the spectrum of $\mathcal{L}_{\varphi}$ is contained in a disk of strictly smaller radius. More specifically, there exist positive functions $h_{j}$ on $\bar{\Sigma}_{j}$ and probability measures $\lambda_{j}$ with support equal to $\Sigma_{j}$ such that

$$
\left\|\mathcal{L}_{\varphi}^{n} f-e^{n \operatorname{Pr}(\varphi)} \sum_{j=0}^{p-1}\left(\int f \mathrm{~d} \lambda_{(j-n \bmod p)}\right) h_{j}\right\| \leqslant C\|f\| e^{-n \varepsilon} e^{n \operatorname{Pr}(\varphi)}
$$

Assume finally that $\mathcal{A}$ is not even recurrent. In this case, one can associate to any component $\mathcal{C}$ the restriction of $\varphi$ to paths staying in $\mathcal{C}$ and the corresponding transfer operator $\mathcal{L}_{\mathcal{C}}$. The previous description applies to $\mathcal{L}_{\mathcal{C}}$ : it has finitely many eigenvalues of maximal modulus $e^{\operatorname{Pr}(\varphi)}$, they are of the form $e^{2 \mathrm{i} k \pi / p_{c}} e^{\operatorname{Pr} \mathcal{C}(\varphi)}$ for
some $k \in \mathbb{Z} / p_{\mathcal{C}} \mathbb{Z}$, and $\mathcal{L}_{\mathcal{C}}$ has a spectral gap. Let $\operatorname{Pr}(\varphi)$ be the maximum of $\operatorname{Pr}_{\mathcal{C}}(\varphi)$ over all components. We call a component maximal if $\operatorname{Pr}_{\mathcal{C}}(\varphi)=\operatorname{Pr}(\varphi)$. The dominating terms in $\mathcal{L}_{\varphi}^{n}$ come from the maximal components. We will say that $\varphi$ is semisimple if there is no directed path from a maximal component to a different maximal component. Otherwise, the eigenvalue $e^{\operatorname{Pr}(\varphi)}$ has nontrivial Jordan blocks, which makes the precise spectral description more cumbersome.

Lemma 3.7. Consider some edge $e_{0}$, and let $k>0$ be such that there is a path from $e_{0}$ to successively $k$ different maximal components. For any nonnegative function $f$ with $f \geqslant 1$ on the set of paths starting with $e_{0}$, one has $\mathcal{L}_{\varphi}^{n} f(\emptyset) \geqslant C n^{k-1} e^{n \operatorname{Pr}(\varphi)}$.

In the semisimple case, the asymptotics of $\mathcal{L}_{\varphi}^{n}$ can be described as follows.
Theorem 3.8. Assume that $\varphi$ is semisimple. Denote by $\mathcal{C}_{1}, \ldots, \mathcal{C}_{I}$ the maximal components, with corresponding period $p_{i}$, and consider for each $i$ a cyclic decomposition $\mathcal{C}_{i}=\bigsqcup_{j \in \mathbb{Z} / p_{i} \mathbb{Z}} \mathcal{C}_{i, j}$. There exist functions $h_{i, j}$ and measures $\lambda_{i, j}$ with $\int h_{i, j} \mathrm{~d} \lambda_{i, j}=1$ such that

$$
\begin{equation*}
\left\|\mathcal{L}_{\varphi}^{n} f-e^{n \operatorname{Pr}(\varphi)} \sum_{i=1}^{I} \sum_{j=0}^{p_{i}-1}\left(\int f \mathrm{~d} \lambda_{i,\left(j-n \bmod p_{i}\right)}\right) h_{i, j}\right\| \leqslant C\|f\| e^{-n \varepsilon} e^{n \operatorname{Pr}(\varphi)} \tag{3.2}
\end{equation*}
$$

The probability measures $\mathrm{d} \mu_{i}=\frac{1}{p_{i}} \sum_{j=0}^{p_{i}-1} h_{i, j} \mathrm{~d} \lambda_{i, j}$ are invariant under $\sigma$ and ergodic.

Denote by $\mathcal{C}_{\rightarrow, i, j}$ the set of edges from which one can reach $\mathcal{C}_{i, j}$ with a path of length in $p_{i} \mathbb{N}$, and by $\mathcal{C}_{i, j, \rightarrow}$ the set of edges that can be reached from $\mathcal{C}_{i, j}$ by a path of length in $p_{i} \mathbb{N}$. The function $h_{i, j}$ is bounded from below on paths beginning by an edge in $\mathcal{C}_{i, j, \rightarrow}$ (and the empty path) and vanishes elsewhere. The support of the measure $\lambda_{i, j}$ is the set of infinite paths beginning in $\mathcal{C}_{\rightarrow, i, j}$ with infinitely many coordinates in $\mathcal{C}_{i}$.

Proof of Lemma 3.7 and Theorem [3.8. The lemma and the theorem will follow from the decomposition of $\mathcal{L}$ into parts coming from the various components, since this decomposition is upper-triangular. Adding a constant to $\varphi$ if necessary, we can without loss of generality assume that $\operatorname{Pr}(\varphi)=0$.

It will be convenient to add one fake component to the components graph, corresponding to the empty word, that can be reached from any true component. For any Hölder function $f$ and any true component $\mathcal{C}$, let $f_{\mathcal{C}}(\omega)=f(\omega)$ if the beginning of $\omega$ is in $\mathcal{C}$, and $f_{\mathcal{C}}(\omega)=0$ otherwise. Let also $f_{\emptyset}(\omega)=f(\emptyset)$ if $\omega=\emptyset$, and 0 otherwise. We get a decomposition $f=\sum_{\mathcal{C}} f_{\mathcal{C}}$ (where $\mathcal{C}$ ranges among all components), corresponding to a decomposition of the space $\mathcal{H}$ of Hölder functions as $\bigoplus_{\mathcal{C}} \mathcal{H}_{\mathcal{C}}$. Denoting by $\Pi_{\mathcal{C}}$ the projection $f \mapsto f_{\mathcal{C}}$, we can split the operator $\mathcal{L}$ into pieces $\mathcal{L}_{\mathcal{C}^{\prime} \rightarrow \mathcal{C}}=\Pi_{\mathcal{C}} \mathcal{L} \Pi_{\mathcal{C}^{\prime}}$, mapping functions supported on paths starting in $\mathcal{C}^{\prime}$ to functions supported on paths starting in $\mathcal{C}$. This gives a block decomposition of $\mathcal{L}$ into positive operators, which is upper-triangular since the components graph has no loop. It follows that one can read the spectrum of $\mathcal{L}$ from the spectrum of the diagonal components $\mathcal{L}_{\mathcal{C} \rightarrow \mathcal{C}}$ (which we simply denote by $\mathcal{L}_{\mathcal{C}}$ ).

Note that $\mathcal{H}_{\mathcal{C}}$ is not the space of Hölder functions defined on paths staying in $\mathcal{C}$ (denoted by $\tilde{\mathcal{H}}_{\mathcal{C}}$ ): a function in $\mathcal{H}_{\mathcal{C}}$ can also be nonzero on paths starting in $\mathcal{C}$, but exiting after some time. Nevertheless, one easily checks that the dominating eigenvalue of $\mathcal{L}_{\mathcal{C} \rightarrow \mathcal{C}}$ coincides on $\mathcal{H}_{\mathcal{C}}$ and $\tilde{\mathcal{H}}_{\mathcal{C}}$.

Let us now give the details of the proof of Lemma3.7. It can be written using the above formalism, but it turns out to be simpler with respect to periodicity problems to give a direct elementary proof (although the intuition to keep in mind is the above one, of an upper-triangular matrix). We fix some paths $u_{0}, \ldots, u_{k-1}$, connecting successive maximal components $\mathcal{C}_{1}, \ldots, \mathcal{C}_{k}$, with $u_{0}$ starting with the edge $e_{0}$. We can then form paths in the automaton of the form $\omega=\left(u_{0}, v_{1}, u_{1}, \ldots, u_{k-1}, v_{k}\right)$, where the path $v_{i}$ is in $\mathcal{C}_{i}$, starts at the end of $u_{i-1}$, and ends at the beginning of $u_{i}$ (except for $v_{k}$, for which there is no condition on its end). Let $f$ be a nonnegative function with $f \geqslant 1$ on the set of paths starting with $e_{0}$. We get

$$
\mathcal{L}_{\varphi}^{n} f(\emptyset) \geqslant \sum_{\omega} e^{S_{n} \varphi(\omega)},
$$

where the sum is restricted to those $\omega$ as above of length $n$. For such an $\omega$, we have since $\varphi$ is Hölder-continuous

$$
S_{n} \varphi(\omega) \geqslant \sum_{i} S_{\left|v_{i}\right|} \varphi\left(v_{i}, u_{i}, v_{i+1}, \ldots\right)-C \geqslant \sum_{i} S_{\left|v_{i}\right|} \varphi\left(v_{i}\right)-C^{\prime} .
$$

Hence, writing $a_{i}$ for $\left|v_{i}\right|$ and $b$ for $\sum\left|u_{i}\right|$, we get

$$
\mathcal{L}_{\varphi}^{n} f(\emptyset) \geqslant e^{-C^{\prime}} \sum_{a_{1}+\cdots+a_{k}=n-b}\left(\sum_{\left|v_{1}\right|=a_{1}} e^{S_{a_{1}} \varphi\left(v_{1}\right)}\right) \cdots\left(\sum_{\left|v_{k}\right|=a_{k}} e^{S_{a_{k}} \varphi\left(v_{k}\right)}\right) .
$$

For $i<k$, denoting by $p_{i}$ the period of $\mathcal{\mathcal { C } _ { i }}$, we restrict now to values $q_{i}$ of $a_{i} \bmod p_{i}$ so that there are trajectories of length $a_{i}$ from the end of $u_{i-1}$ to the start of $u_{i}$. For such an $a_{i}$, the spectral properties of the transfer operator restricted to $\mathcal{C}_{i}$ give $\sum e^{S a_{i} \varphi\left(v_{i}\right)} \geqslant C e^{a_{i} \operatorname{Pr}(\varphi)}=C$ since $\operatorname{Pr}(\varphi)=0$. Summing over $i$, we obtain

$$
\mathcal{L}_{\varphi}^{n} f(\emptyset) \geqslant C \sum_{\substack{a_{1}+\cdots+k_{k}=n-b \\ \forall i<k, a_{i}=q_{i} \bmod p_{i}}} 1 .
$$

This grows at least like $C n^{k-1}$, proving the lemma.
Let us now prove Theorem 3.8] For $f \in \mathcal{H}$, we decompose $\mathcal{L}^{n} f$ as

$$
\begin{equation*}
\sum_{\tau=\left(\mathcal{C}_{1} \rightarrow \cdots \rightarrow \mathcal{C}_{k}\right)} \sum_{a_{1}+\cdots+a_{k}=n-k+1} \mathcal{L}_{\mathcal{C}_{k}}^{a_{k}} \mathcal{L}_{\mathcal{C}_{k-1} \rightarrow \mathcal{C}_{k}} \mathcal{L}_{\mathcal{C}_{k-1}}^{a_{k-1}} \cdots \mathcal{L}_{\mathcal{C}_{1}}^{a_{1}} f \tag{3.3}
\end{equation*}
$$

where the first sum is over all paths $\tau$ in the components graph. (Note that there are finitely many such paths since there is no loop in the components graph.) If such a path does not involve any maximal component, then the norm of each factor $\mathcal{L}_{\mathcal{C}_{\ell}}^{a_{\ell}}$ is bounded by $C e^{-a_{\ell} \varepsilon}$ for some $\varepsilon>0$, yielding a bound $C e^{-n \varepsilon}\|f\|$ for this contribution. On the other hand, if there is a maximal component, it has to be unique by semisimplicity, say at the index $i$. Denote by $p_{i}$ the period of $\mathcal{C}_{i}$ and by $\mathcal{C}_{i}=\bigsqcup_{j=0}^{p_{i}-1} \mathcal{C}_{i, j}$ the cyclic decomposition of $\mathcal{C}_{i}$. We have $\mathcal{L}_{\mathcal{C}_{i}}^{a_{i}} g=$ $\sum_{j=0}^{p_{i}-1}\left(\int g \mathrm{~d} \lambda_{j-a_{i} \operatorname{cod} p_{i}}^{\mathcal{C}_{i}}\right) h_{j}^{\mathcal{C}_{i}}+O\left(e^{-a_{i} \varepsilon}\|g\|\right)$ for any $g \in \mathcal{H}_{\mathcal{C}_{i}}$, where the nonnegative function $h_{j}^{\mathcal{C}_{i}}$ is supported on paths starting with a point in $\mathcal{C}_{i, j}$, and the measure $\lambda_{j}^{\mathcal{C}_{i}}$ is supported on infinite paths staying in $\mathcal{C}_{i}$ and starting with a point in $\mathcal{C}_{i, j}$. It readily follows that the term in (3.3) corresponding to the path $\tau=\left(\mathcal{C}_{1} \rightarrow \cdots \rightarrow \mathcal{C}_{k}\right)$
gives a contribution

$$
\sum_{j=0}^{p_{i}-1}\left(\int f \mathrm{~d} \lambda_{j-n \bmod p_{i}}^{\tau}\right) h_{j}^{\tau}+O\left(e^{-n \varepsilon}\|f\|\right)
$$

where

$$
h_{j}^{\tau}=\sum_{a_{i+1}+\cdots+a_{k}+k-i=0 \bmod p_{i}} \mathcal{L}_{\mathcal{C}_{k}}^{a_{k}} \mathcal{L}_{\mathcal{C}_{k-1} \rightarrow \mathcal{C}_{k}} \mathcal{L}_{\mathcal{C}_{k-1}}^{a_{k-1}} \cdots \mathcal{L}_{\mathcal{C}_{i} \rightarrow \mathcal{C}_{i+1}} h_{j}^{\mathcal{C}_{i}}
$$

[note that this series is well-defined since $\left.\left\|\mathcal{L}_{\mathcal{C}_{j}}^{a_{j}}\right\|=O\left(e^{-a_{j} \varepsilon}\right)\right]$ and

$$
\lambda_{j}^{\tau}=\sum_{a_{1}+\cdots+a_{i-1}+i=0 \bmod p_{i}}\left(\mathcal{L}_{\mathcal{C}_{i-1} \rightarrow \mathcal{C}_{i}} \mathcal{L}_{\mathcal{C}_{i-1}}^{a_{i-1}} \cdots \mathcal{L}_{\mathcal{C}_{1}}^{a_{1}}\right)^{*} \lambda_{j}^{\mathcal{C}_{i}},
$$

this series being again finite. The result follows by adding the contributions of all such paths $\tau$.

We will need the following simple lemma later on.
Lemma 3.9. Under the assumptions of the above theorem, let $\beta_{i}=\sum_{j} \lambda_{i, j}$. Then $\sigma_{*} \beta_{i}$ is absolutely continuous with respect to $\beta_{i}$.

Proof. The measures $\lambda_{i, j}$ are constructed as eigenmeasures of the operator $\left(\mathcal{L}_{\varphi}^{*}\right)^{p_{i}}$. More precisely, they satisfy $\mathcal{L}_{\varphi}^{*} \lambda_{i, j}=e^{\operatorname{Pr}(\varphi)} \lambda_{i,\left(j-1 \bmod p_{i}\right)}$. In particular, $\mathcal{L}_{\varphi}^{*} \beta_{i}=$ $e^{\operatorname{Pr}(\varphi)} \beta_{i}$.

Consider a cylinder $\left[\omega_{0}, \ldots, \omega_{n}\right]$, i.e., the set of paths that start with those symbols. The function $\mathcal{L}_{\varphi} 1_{\left[\omega_{0}, \ldots, \omega_{n}\right]}$ is uniformly bounded on the image of this cylinder under $\sigma$, i.e., $\left[\omega_{1}, \ldots, \omega_{n}\right]$, and it vanishes elsewhere. Hence,

$$
\begin{aligned}
\beta_{i}\left(\left[\omega_{0}, \ldots, \omega_{n}\right]\right) & =e^{-\operatorname{Pr}(\varphi)} \mathcal{L}_{\varphi}^{*} \beta_{i}\left(1_{\left[\omega_{0}, \ldots, \omega_{n}\right]}\right)=e^{-\operatorname{Pr}(\varphi)} \beta_{i}\left(\mathcal{L}_{\varphi} 1_{\left[\omega_{0}, \ldots, \omega_{n}\right]}\right) \\
& \leqslant C \beta_{i}\left(\left[\omega_{1}, \ldots, \omega_{n}\right]\right) .
\end{aligned}
$$

Since $\sigma^{-1}\left(\left[\omega_{1}, \ldots, \omega_{n}\right]\right)$ is a finite union of cylinders of the form $\left[\omega_{0}, \ldots, \omega_{n}\right]$, we obtain $\beta_{i}\left(\sigma^{-1}\left[\omega_{1}, \ldots, \omega_{n}\right]\right) \leqslant C \beta_{i}\left(\left[\omega_{1}, \ldots, \omega_{n}\right]\right)$. As cylinders generate the topology, it follows that $\beta_{i}\left(\sigma^{-1} A\right) \leqslant C \beta_{i}(A)$ for any measurable set $A$.

Finally, we will need to describe what happens under perturbations of the potential.

Proposition 3.10. Let $\varphi \in \mathcal{H}^{\beta}$ be a semisimple Hölder potential, with maximal components $\mathcal{C}_{1}, \ldots, \mathcal{C}_{I}$ and a spectral description as in Theorem 3.8. There exist $\varepsilon>0$ and $C>0$ such that, for any $\psi$ that is small enough in $\mathcal{H}^{\beta}$, there exist functions $h_{i, j}^{\psi}$ and measures $\lambda_{i, j}^{\psi}$ (with the same support as, respectively, $h_{i, j}$ and $\left.\lambda_{i, j}\right)$ and numbers $\operatorname{Pr}_{i}(\varphi+\psi)$ with

$$
\left\|\mathcal{L}_{\varphi+\psi}^{n} f-\sum_{i=1}^{I} e^{n \operatorname{Pr}_{i}(\varphi+\psi)} \sum_{j=0}^{p_{i}-1}\left(\int f \mathrm{~d} \lambda_{i,\left(j-n \bmod p_{i}\right)}^{\psi}\right) h_{i, j}^{\psi}\right\| \leqslant C\|f\| e^{-n \varepsilon} e^{n \operatorname{Pr}(\varphi)}
$$

The maps $\psi \mapsto \operatorname{Pr}_{i}(\varphi+\psi), \psi \mapsto h_{i, j}^{\psi}$, and $\psi \mapsto \lambda_{i, j}^{\psi}$ are real analytic (in the norm sense) from a small ball around 0 in $\mathcal{H}^{\beta}$ to, respectively, $\mathbb{R}, \mathcal{H}^{\beta}$ and the dual of $\mathcal{H}^{\beta}$.

Finally,

$$
\operatorname{Pr}_{i}(\varphi+\psi)=\operatorname{Pr}(\varphi)+\int \psi \mathrm{d} \mu_{i}+O\left(\|\psi\|_{\mathcal{H}^{\beta}}^{2}\right),
$$

where $\mathrm{d} \mu_{i}=\frac{1}{p_{i}} \sum_{j=0}^{p_{i}-1} h_{i, j} \mathrm{~d} \lambda_{i, j}$.
Proof. Let us first assume that the shift is topologically mixing. In this case, the dominating eigenvalue $e^{\operatorname{Pr}(\varphi)}$ of $\mathcal{L}_{\varphi}$ is simple. Simple isolated eigenvalues and the corresponding eigenprojectors and eigenfunctions depend in an analytic way on the operator, by classical perturbation theory [Kat66, Theorem VII.1.8]. Moreover, by semicontinuity of the spectrum, the perturbed operators $\mathcal{L}_{\varphi+\psi}$ also have a spectral gap, uniformly in $\psi$ close enough to 0 . One gets

$$
\mathcal{L}_{\varphi+\psi}^{n} f=e^{n \operatorname{Pr}(\varphi+\psi)}\left(\int f \mathrm{~d} \lambda^{\psi}\right) h^{\psi}+O\left(e^{-n \varepsilon} e^{n \operatorname{Pr}(\varphi)}\right)
$$

for some $\operatorname{Pr}(\varphi+\psi), \lambda^{\psi}$, and $h^{\psi}$ that depend analytically on $\psi$. This almost completes the proof of the theorem in this case; it only remains to show that $\operatorname{Pr}(\varphi+\psi)=\operatorname{Pr}(\varphi)+\int \psi h \mathrm{~d} \lambda+O\left(\|\psi\|^{2}\right)$. By analyticity, it is sufficient to show that the derivative of the pressure at 0 is given by the integral with respect to the measure $h \mathrm{~d} \lambda$. This is [PP90, Proposition 4.10].

The topologically transitive case readily reduces to the mixing case by consider$\operatorname{ing} \sigma^{p}$ where $p$ is the period.

In the general case, one obtains different pressures $\operatorname{Pr}_{i}(\varphi+\psi)$ on each component $\mathcal{C}_{i}$. On other components that were not maximal for $\varphi$, the pressure of $\varphi+\psi$ remains bounded away from $\operatorname{Pr}(\varphi)$. It follows that the maximal components of $\varphi+\psi$ are contained in those of $\varphi$ if $\psi$ is small enough. In particular, $\varphi+\psi$ is semisimple, and $\operatorname{Pr}(\varphi+\psi)=\max \operatorname{Pr}_{i}(\varphi+\psi)$. Finally, the spectral description of $\mathcal{L}_{\varphi+\psi}$ follows from the description on each component $\mathcal{C}_{i}$ separately.
3.4. Transfer operators in hyperbolic groups. Let $\Gamma$ be a nonelementary Gromov-hyperbolic group, and $\mu$ a probability measure satisfying strong Ancona inequalities. Consider a strongly Markov automatic structure for $\Gamma$, given by a directed graph $\mathcal{A}=\left(V, E, s_{*}\right)$ and a labeling $\alpha: E \rightarrow S$ of edges by generators of the group. We will use freely the notations of Paragraph 3.2.

For $r \in[1, R]$, let us define a potential $\varphi_{r}$ on the set $\Sigma^{*}$ of finite paths in the automaton by

$$
\varphi_{r}(\omega)=\log \left(\frac{H_{r}\left(e, \alpha_{*}(\omega)\right)}{H_{r}\left(e, \alpha_{*}(\sigma \omega)\right)}\right)
$$

Consider a path $\omega=\omega_{0} \cdots \omega_{n-1}$ of length $n$. Then

$$
e^{S_{n} \varphi_{r}(\omega)}=\frac{H_{r}\left(e, \alpha_{*}\left(\omega_{0} \cdots \omega_{n-1}\right)\right)}{H_{r}(e, e)} .
$$

Let $E_{*}$ be the set of edges starting from the vertex $s_{*}$ of the graph $\mathcal{A}$, and let $1_{\left[E_{*}\right]}$ be the function equal to 1 on paths starting with an edge in $E_{*}$, and 0 elsewhere. Using the language of transfer operators, we have

$$
\begin{aligned}
H_{r}(e, e) \cdot \mathcal{L}_{\varphi_{r}}^{n} 1_{\left[E_{*}\right]}(\emptyset) & =H_{r}(e, e) \sum_{\omega=\omega_{0} \cdots \omega_{n-1}} e^{S_{n} \varphi_{r}(\omega)} 1\left(\omega_{0} \in E_{*}\right) \\
& =\sum H_{r}\left(e, \alpha_{*}\left(\omega_{0} \cdots \omega_{n-1}\right)\right) 1\left(\omega_{0} \in E_{*}\right) .
\end{aligned}
$$

Since $\alpha_{*}$ induces a bijection between the paths of length $n$ starting from $s_{*}$ and the sphere $\mathbb{S}_{n}$ of radius $n$ in $\Gamma$, we obtain

$$
\sum_{x \in \mathbb{S}_{n}} H_{r}(e, x)=H_{r}(e, e) \mathcal{L}_{\varphi_{r}}^{n} 1_{\left[E_{*}\right]}(\emptyset) .
$$

Therefore, the spectrum of $\mathcal{L}_{\varphi_{r}}$ will give asymptotics for $\sum_{x \in \mathbb{S}_{n}} H_{r}(e, x)$. To be able to use the results of the previous paragraph, one should check that $\varphi_{r}$ is Hölder-continuous.

Lemma 3.11. There exists $\beta>0$ such that, for any $r \in[1, R]$, the function $\varphi_{r}$ is Hölder-continuous of exponent $\beta$ on $\Sigma^{*}$. Therefore, it extends to a Höldercontinuous function on $\bar{\Sigma}$ that we still denote by $\varphi_{r}$. It satisfies $\left\|\varphi_{r}\right\|_{\mathcal{H}^{\beta}} \leqslant C$, uniformly in $r \in[1, R]$. Moreover,

$$
\begin{equation*}
\left\|\varphi_{r}-\varphi_{R}\right\|_{\mathcal{H}^{\beta}} \leqslant C(R-r)^{1 / 3} . \tag{3.4}
\end{equation*}
$$

Proof. Consider two finite paths $\omega$ and $\omega^{\prime}$ with $d\left(\omega, \omega^{\prime}\right)=2^{-n}<1$; by definition they match up to length $n \geqslant 1$. In particular, $\omega_{0}=\omega_{0}^{\prime}$. Let $x=\alpha_{*}(\omega)$, $x^{\prime}=\alpha_{*}\left(\omega^{\prime}\right)$, and $a=\alpha_{*}\left(\omega_{0}\right)$, so that $\varphi_{r}(\omega)=\log \left(H_{r}(e, x) / H_{r}(a, x)\right)$ and $\varphi_{r}\left(\omega^{\prime}\right)=$ $\log \left(H_{r}\left(e, x^{\prime}\right) / H_{r}\left(a, x^{\prime}\right)\right)$. The points $e, a$ and $x, x^{\prime}$ are in the situation of strong Ancona inequalities (Definition 2.8) with a separating distance $n-1$. Since $\mu$ satisfies strong uniform Ancona inequalities, it follows that $\left|\varphi_{r}(\omega)-\varphi_{r}\left(\omega^{\prime}\right)\right| \leqslant C e^{-\rho n}$ for some $\rho>0$. Hence, for some $\beta^{\prime}>0, \varphi_{r}$ belongs to $\mathcal{H}^{\beta^{\prime}}$ and is uniformly bounded in this space.

Lemma 3.4implies that $\left\|\varphi_{r}-\varphi_{R}\right\|_{C^{0}} \leqslant C(R-r)^{1 / 2}$. Together with the uniform boundedness of $\varphi_{r}$ in $\mathcal{H}^{\beta^{\prime}}$, this shows that $\left\|\varphi_{r}-\varphi_{R}\right\|_{\mathcal{H}^{\beta}} \leqslant C(R-r)^{1 / 3}$ if $\beta$ is small enough, by (3.1).

Finally, we have proved all those inequalities on the space $\Sigma^{*}$ of finite paths. Since Hölder-continuous functions on $\Sigma^{*}$ extend to Hölder-continuous functions on $\bar{\Sigma}$, the result follows.

Remark 3.12. One could in fact show that $\left\|\varphi_{r}-\varphi_{R}\right\|_{\mathcal{H}^{\beta}} \leqslant C(R-r)^{1 / 2}$ by mimicking the proof of Lemma 3.4 at the level of Hölder exponents. Since this computation is lengthy and (3.4) will be sufficient for our purposes, we omit it.

Lemma 3.13. We have $\operatorname{Pr}\left(\varphi_{R}\right)=0$. Moreover, $\varphi_{R}$ is semisimple.
Proof. Suppose $\operatorname{Pr}\left(\varphi_{R}\right)<0$. Then $\mathcal{L}_{\varphi_{R}}^{n} 1_{\left[E_{*}\right]}$ goes to zero exponentially fast in the space of Hölder functions. In particular, $\sum_{x \in \mathbb{S}_{n}} H_{R}(e, x)=H_{R}(e, e) \mathcal{L}_{\varphi_{R}}^{n} 1_{\left[E_{*}\right]}(\emptyset)$ is exponentially small. One can use this estimate to prove that the series $G_{R+\varepsilon}(e, e)$ converges for some $\varepsilon>0$ : this is the content of the proof of Proposition 7.1 in GL13] (the proof is written for symmetric measures, but it applies equally well in nonsymmetric situations). This is a contradiction since, by definition, $R$ is the radius of convergence of the series $G_{r}(e, e)$. Hence, $\operatorname{Pr}\left(\varphi_{R}\right) \geqslant 0$.

If $\operatorname{Pr}\left(\varphi_{R}\right)$ were strictly positive, or $\operatorname{Pr}\left(\varphi_{R}\right)=0$ but $\varphi_{R}$ were not semisimple, then Lemma 3.7 would imply that $\mathcal{L}_{\varphi_{R}}^{n} 1_{\left[E_{*}\right]}(\emptyset)$ would tend to infinity. This quantity is equal to $H_{R}(e, e)^{-1} \sum_{x \in \mathbb{S}_{n}} H_{R}(e, x)$. Since it remains bounded by Lemma 2.5, we obtain a contradiction.

One can now come back to Corollary 3.3. Since $\operatorname{Pr}\left(\varphi_{R}\right)=0$ and $\varphi_{R}$ is semisimple, Theorem 3.8 implies in particular that $\mathcal{L}_{\varphi_{R}}^{n} 1_{\left[E_{*}\right]}(\emptyset)$ is bounded from below. Since
it coincides with $H_{R}(e, e)^{-1} \sum_{x \in \mathbb{S}_{n}} H_{R}(e, x)$, we get $\sum_{x \in \Gamma} H_{R}(e, x)=+\infty$. This shows that the constant $A$ in Corollary 3.3 vanishes, and therefore

$$
\begin{equation*}
\frac{C^{-1}}{\sqrt{R-r}} \leqslant \sum_{x \in \Gamma} H_{r}(e, x) \leqslant \frac{C}{\sqrt{R-r}} \tag{3.5}
\end{equation*}
$$

Let us introduce a convenient notation: we will reserve the notation $\tau(r)$ (possibly with some indices) for continuous functions of $r$ taking values in $(0,+\infty)$ that extend continuously up to $r=R$ and are bounded away from zero.

We will now use the spectral perturbation given by Proposition 3.10 to study $\mathcal{L}_{\varphi_{r}}$. If $r$ is close to $R$, then $\varphi_{r}-\varphi_{R}$ is small in $\mathcal{H}^{\beta}$ by Lemma 3.11, Applying the proposition on spectral perturbation to the function $f=1_{\left[E_{*}\right]}$, we get the following. Let $p$ be the least common multiple of the periods of the maximal components of $\varphi_{R}$. For any $q \in[0, p)$, one has [since $\left.\operatorname{Pr}\left(\varphi_{R}\right)=0\right]$

$$
\mathcal{L}_{\varphi_{r}}^{n p+q} 1_{\left[E_{*}\right]}(\emptyset)=\sum_{i=1}^{I} e^{(n p+q) \operatorname{Pr}_{i}\left(\varphi_{r}\right)} \tau_{0}(q, i, r)+O\left(e^{-n \varepsilon}\right),
$$

for some functions $\tau_{0}(q, i, r)$ (as in the notation we introduced in the previous paragraph). This is equal to $H_{r}(e, e)^{-1} \sum_{x \in \mathbb{S}_{n p+q}} H_{r}(e, x)$. Since $\sum_{x \in \Gamma} H_{r}(e, x)<$ $\infty$, it follows in particular that $\operatorname{Pr}_{i}\left(\varphi_{r}\right)$ is strictly negative for all $i$.

Summing over $n$ and $q$, we get

$$
\begin{align*}
\sum_{x \in \Gamma} H_{r}(e, x) & =H_{r}(e, e) \sum_{n, q} \mathcal{L}_{\varphi_{r}}^{n p+q} 1_{\left[E_{*}\right]}(\emptyset) \\
& =H_{r}(e, e) \sum_{q=0}^{p-1} \sum_{i=1}^{I} \frac{e^{q \operatorname{Pr}_{i}\left(\varphi_{r}\right)}}{1-e^{p \operatorname{Pr}_{i}\left(\varphi_{r}\right)}} \tau_{0}(q, i, r)+O(1) \\
& =\sum_{i=1}^{I} \frac{\tau_{1}(i, r)}{\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|}+O(1), \tag{3.6}
\end{align*}
$$

for some functions $\tau_{1}(i, r)$.
By (3.5), $\left|\operatorname{Pr}\left(\varphi_{r}\right)\right|=\inf _{i}\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|$ is comparable to $\sqrt{R-r}$. It will be important to show that all the $\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|$ are of the same order of magnitude: otherwise, some components would not play a significant role for $r<R$ while they would become important at $r=R$, ruining the continuity properties we are seeking. This is the main difference with the transitive situation, where there is only one eigenvalue to consider.

Theorem 3.14. For any $i \in[1, I]$, the ratio $\operatorname{Pr}_{i}\left(\varphi_{r}\right) / \operatorname{Pr}\left(\varphi_{r}\right)$ tends to 1 when $r \rightarrow R$.

We will prove this theorem in the next subsection. It follows from this result that

$$
\begin{equation*}
\sum_{x \in \Gamma} H_{r}(e, x)=\frac{\tau_{2}(r)}{\left|\operatorname{Pr}\left(\varphi_{r}\right)\right|}+O(1) \tag{3.7}
\end{equation*}
$$

Hence, the spectral data of $\mathcal{L}_{\varphi_{r}}$ are related to the function $\eta(r)=\sum_{x \in \Gamma} H_{r}(e, x)$.
3.5. Pressure does not depend on the component. In this subsection, we prove Theorem 3.14. We will in particular rely on the estimate $\left\|\varphi_{r}-\varphi_{R}\right\|_{\mathcal{H}^{\beta}} \leqslant$ $C(R-r)^{1 / 3}$ from (3.4) that in turn was proved using the a priori estimates from Lemma 3.4

By Proposition 3.10, the variation of the pressure mainly depends on the integral $\int\left(\varphi_{r}-\varphi_{R}\right) \mathrm{d} \mu_{i}$. We will show that this integral does not depend on $i$, using a geometric argument in the group due to CF10.

Fix some $r \in[1, R]$. For $c \in \mathbb{R}$, we define a set $U(c) \subset \partial \Gamma$ as the set of points $\xi$ such that, along some geodesic from $e$ to $\xi, \log H_{r}(e, x) / d(e, x) \rightarrow c$. Equivalently, this convergence holds along any geodesic tending to $\xi$, and one can replace $H_{r}(e, x)$ with $H_{r}(a, x)$ and $d(e, x)$ with $d(b, x)$ for any $a, b \in \Gamma$. Indeed, geodesics tending to $\xi$ remain within a bounded distance from each other, by [GdIH90, Proposition 7.2] (therefore, by Harnack inequalities $H_{r}$ varies by at most a multiplicative constant when one changes geodesics), and the ratio $H_{r}(e, x) / H_{r}(a, x)$ also remains bounded from above and from below again by Harnack inequalities. In particular, $U(c)$ is invariant under the action of $\Gamma$ : for any $g \in \Gamma, g \cdot U(c)=U(c)$.

Let $c_{i}=\int \varphi_{r} \mathrm{~d} \mu_{i}$, we will show that, for all $i \neq i^{\prime}, g \cdot U\left(c_{i}\right)$ intersects $U\left(c_{i^{\prime}}\right)$ for some $g \in \Gamma$. This will give $U\left(c_{i}\right)=U\left(c_{i^{\prime}}\right)$, and hence $c_{i}=c_{i^{\prime}}$ as desired. To prove this, we will show that the sets $U\left(c_{i}\right)$ all have positive measures for some measure on $\partial \Gamma$ that is ergodic under the action of $\Gamma$.

Let us first construct the measure. Let $p$ be the least common multiple of the periods $p_{i}$, and fix $q \in[0, p)$. It follows from the spectral description of $\mathcal{L}_{\varphi_{R}}$ (Theorem (3.8) that, for any Hölder-continuous function $f$ on $\bar{\Sigma}, \mathcal{L}_{\varphi_{R}}^{n p+q} f(\emptyset)$ converges when $n \rightarrow \infty$. In turn, this convergence follows for any continuous function, by approximation (since the iterates of $\mathcal{L}_{\varphi_{R}}$ on $C^{0}$ remain bounded, since $\mathcal{L}_{\varphi_{R}}^{n} 1$ itself remains bounded). If $f$ is a continuous function on $\Gamma \cup \partial \Gamma$, then $\sum_{x \in \mathbb{S}_{n}} H_{R}(e, x) f(x)=H_{R}(e, e) \mathcal{L}_{\varphi_{R}}^{n}\left(1_{\left[E_{*}\right]} \cdot f \circ \alpha_{*}\right)(\emptyset)$, and $f \circ \alpha_{*}$ is continuous. Let us define a measure $m_{n}$ supported on $\mathbb{S}_{n}$ by $m_{n}=\sum_{x \in \mathbb{S}_{n}} H_{R}(e, x) \delta_{x}$; this shows that the sequence of measures $m_{n p+q}$ converges to a limiting measure (which is supported on $\partial \Gamma$ and has mass bounded from above and from below). We deduce that the measures $\left(\sum_{1}^{N} m_{n}\right) /\left(\sum_{1}^{N} m_{n}(\Gamma)\right)$ converge to a probability measure on $\partial \Gamma$, which we denote by $\nu_{R}$. It also follows that this measure can be constructed using the Patterson-Sullivan technique: the measures

$$
\theta_{s}=\sum_{x \in \Gamma} H_{R}(e, x) e^{-s|x|} \delta_{x} / \sum_{x \in \Gamma} H_{R}(e, x) e^{-s|x|}
$$

are well-defined for $s>0$, and they converge when $s$ tends to 0 toward $\nu_{R}$.
For $g \in \Gamma$, let us denote by $L_{g}$ the left multiplication by $\Gamma$. Then, for any $x \in \Gamma$,

$$
\begin{aligned}
\left(L_{g}\right)_{*} \theta_{s}(x) & =\theta_{s}\left(g^{-1} x\right)=\frac{H_{R}\left(e, g^{-1} x\right)}{H_{R}(e, x)} e^{-s\left(\left|g^{-1} x\right|-|x|\right)} \theta_{s}(x) \\
& =\tilde{K}_{x}(g) e^{-s\left(\left|g^{-1} x\right|-|x|\right)} \theta_{s}(x)
\end{aligned}
$$

where $\tilde{K}_{x}(g)=H_{R}(g, x) / H_{R}(e, x)$ is the Martin Kernel associated to $H_{R}$. When $x$ tends to a point $\xi \in \partial \Gamma$, this quantity converges to a limit denoted by $\tilde{K}_{\xi}(g)$. Since $\left|g^{-1} x\right|-|x|$ is uniformly bounded when $x$ varies in $\Gamma$, we deduce letting $s$ tend to 0 that

$$
\begin{equation*}
\frac{\mathrm{d}\left(L_{g}\right)_{*} \nu_{R}}{\mathrm{~d} \nu_{R}}(\xi)=\tilde{K}_{\xi}(g) \tag{3.8}
\end{equation*}
$$

Following the classical arguments of Patterson-Sullivan (due in this context to Coo93 and BHM11), we deduce the following.

Proposition 3.15. The measure $\nu_{R}$ is ergodic for the action of $\Gamma$.
Proof. We want to apply the results of Coo93 and BHM11 saying that a PattersonSullivan measure is ergodic. Thus, we should interpret the function $\tilde{K}_{\xi}(g)$ in (3.8) as the exponential of a Busemann cocycle. As $\tilde{K}_{\xi}(g)$ is the limit of $H_{R}(g, x) / H_{R}(e, x)$, the function $\log \tilde{K}_{\xi}$ would be the Busemann cocycle associated to a distance $\tilde{d}$ if $H_{R}(x, y)=C e^{-\tilde{d}(x, y)}$, for some constant $C$. Let us therefore set $\tilde{d}(x, y)=$ $-\log \left(F_{R}(x, y) F_{R}(y, x)\right)$, where $F_{R}$ is the first visit Green function. We should show that $\tilde{d}$ is a distance, that it is equivalent to $d$, and that it is hyperbolic, to be able to apply the results of Coo93 and BHM11.

The subadditivity (2.3) of $F_{R}$ shows that $\tilde{d}$ satisfies the triangular inequality. For $x \neq y$, considering $n$ concatenations of paths from $x$ to $y$ then to $x$, one gets

$$
G_{R}(x, x) \geqslant \sum_{n=0}^{\infty}\left(F_{R}(x, y) F_{R}(y, x)\right)^{n}
$$

Since $G_{R}(x, x)$ is finite, this shows that $F_{R}(x, y) F_{R}(y, x)<1$. Hence, $\tilde{d}$ is a distance. It is a variant of the Green distance studied in [BHM11].

The quantity $G_{R}(e, e)$, which is finite, equals $\sum_{\gamma} w_{R}(\gamma)$ [where the sum is over all paths from $e$ to itself, and the notation $w_{R}(\gamma)$ for the $R$-weight of a path $\gamma$ has been introduced in Subsection 2.1. Excluding finitely many paths, one can make the remaining sum arbitrarily small. If $x$ is not on one of those finitely many paths, then $F_{R}(e, x) F_{R}(x, e)$ is bounded by the remaining sum and is therefore arbitrarily small. This shows that $\tilde{d}(e, x)$ tends to infinity when $x \rightarrow \infty$ in $\Gamma$.

Since $G_{R}$ (or, equivalently, $F_{R}$ ) satisfies Ancona inequalities, there exists $D>0$ such that $\tilde{d}(x, z) \geqslant \tilde{d}(x, y)+\tilde{d}(y, z)-D$ whenever $x, y, z$ are on a geodesic segment in this order. Let $L$ be such that $\tilde{d}(e, x) \geqslant 2 D$ for $|x| \geqslant L$. By induction, this implies that $\tilde{d}(e, x) \geqslant(n+1) D$ for $|x| \geqslant n L$. In particular, there exists a constant $C>0$ such that, for all $x, \tilde{d}(e, x) \geqslant C^{-1}|x|$. By Harnack inequalities (2.2), we also have $\tilde{d}(e, x) \leqslant C|x|$. This shows that the distance $\tilde{d}$ is equivalent to the word distance $d$.

The word distance is hyperbolic. It does not immediately follow that $\tilde{d}$ is hyperbolic, since the metric space $(\Gamma, \tilde{d})$ is usually not geodesic (while the invariance of hyperbolicity under quasi-isometries requires such an assumption). However, BHM11 proves that if Ancona inequalities hold, then $\tilde{d}$ is hyperbolic (the proof given in their Theorem 1.1 is for the usual Green metric, but it applies verbatim in our setting).

Finally, we can apply the results of Paragraphs 2.2 and 2.3 in BHM11. Equation (3.8) shows that $\nu_{R}$ is quasiconformal for a distance at infinity coming from the hyperbolic distance $\tilde{d}$ on $\Gamma$. Therefore, [BHM11, Theorem 2.7] implies that $\nu_{R}$ is ergodic.

Proposition 3.16. For $i \neq i^{\prime}$, one has $\int \varphi_{r} \mathrm{~d} \mu_{i}=\int \varphi_{r} \mathrm{~d} \mu_{i^{\prime}}$.
Proof. The limit of $\mathcal{L}_{\varphi_{R}}^{n p+q} f(\emptyset)$ is given by $\sum_{i=1}^{I} \sum_{j=0}^{p_{i}-1}\left(\int f \mathrm{~d} \lambda_{i,\left(j-q \bmod p_{i}\right)}\right) h_{i, j}(\emptyset)$. Since $h_{i, j}(\emptyset)$ is bounded from above and from below, we deduce that $\nu_{R}$ is equivalent
to the push-forward under $\alpha_{*}$ of the measure $\sum_{i, j} \lambda_{i, j}$ restricted to the set of paths beginning with an edge in $E_{*}$.

The probability measure $\mathrm{d} \mu_{i}=\frac{1}{p_{i}} \sum_{j=0}^{p_{i}-1} h_{i, j} \mathrm{~d} \lambda_{i, j}$ is invariant and ergodic for the left shift $\sigma$. Let $O_{i} \subset \Sigma$ denote the set of points such that the normalized Birkhoff sums $S_{n} f / n$ converge to $\int f \mathrm{~d} \mu_{i}$ for any continuous function $f$. By the Birkhoff ergodic theorem, $\mu_{i}\left(O_{i}\right)=1$. Since $\mu_{i}$ is equivalent to $\beta_{i}=\sum_{j} \lambda_{i, j}$ restricted to the set $\bar{\Sigma}_{i}$ of paths staying in the component $\mathcal{C}_{i}$, we get $\beta_{i}\left(O_{i}^{c} \cap \bar{\Sigma}_{i}\right)=0$ (where $O_{i}^{c}$ denotes the complement of $O_{i}$ ). We deduce that

$$
\begin{equation*}
\beta_{i}\left(O_{i}^{c}\right)=0 . \tag{3.9}
\end{equation*}
$$

Otherwise, since $\beta_{i}$ almost every point ends up in $\bar{\Sigma}_{i}$ after finitely many iterations, we would have $\beta_{i}\left(O_{i}^{c} \cap \sigma^{-k} \bar{\Sigma}_{i}\right)>0$ for some $k \geqslant 0$; hence $\beta_{i}\left(\sigma^{-k}\left(\sigma^{k} O_{i}^{c} \cap \bar{\Sigma}_{i}\right)\right)>0$. Since $\sigma_{*}^{k} \beta_{i}$ is absolutely continuous with respect to $\beta_{i}$ by Lemma 3.9, and $\sigma^{k} O_{i}^{c} \subset$ $O_{i}^{c}$, this gives $\beta_{i}\left(O_{i}^{c} \cap \bar{\Sigma}_{i}\right)>0$, a contradiction.

Let us now show that

$$
\begin{equation*}
\nu_{R}\left(\alpha_{*}\left(O_{i} \cap\left[E_{*}\right]\right)\right)>0, \tag{3.10}
\end{equation*}
$$

where $O_{i} \cap\left[E_{*}\right]$ denotes the set of paths in $O_{i}$ beginning with an edge in $E_{*}$. Otherwise, since the image of $\beta_{i}\left(\cdot \cap\left[E_{*}\right]\right)$ is absolutely continuous with respect to $\nu_{R}$, we would get $\left(\alpha_{*} \beta_{i}\right)\left(\alpha_{*}\left(O_{i} \cap\left[E_{*}\right]\right)\right)=0$; hence $\beta_{i}\left(O_{i} \cap\left[E_{*}\right]\right)=0$. Since $\beta_{i}\left(O_{i}^{c}\right)=0$ by (3.9), we get $\beta_{i}\left(\left[E_{*}\right]\right)=0$. This is a contradiction since Theorem 3.8 shows that $\beta_{i}$ gives positive weight to $\left[E_{*}\right]$.

Consider now $\omega \in O_{i} \cap\left[E_{*}\right]$, and let $\xi=\alpha_{*}(\omega) \in \partial \Gamma$. The path $\alpha(\omega)$ is a geodesic converging to $\xi$. In particular, denoting by $\bar{\omega}_{n}$ the beginning of $\omega$ of length $n, x_{n}=\alpha_{*}\left(\bar{\omega}_{n}\right)$ is a sequence of points converging to $\xi$ along a geodesic ray. Moreover,

$$
\log H_{r}\left(e, x_{n}\right)=S_{n} \varphi_{r}\left(\bar{\omega}_{n}\right)+\log H_{r}(e, e) .
$$

Since $\varphi_{r}$ is Hölder-continuous, $S_{n} \varphi_{r}\left(\bar{\omega}_{n}\right)-S_{n} \varphi_{r}(\omega)$ remains uniformly bounded. Hence, $\log H_{r}\left(e, x_{n}\right) / n=S_{n} \varphi_{r}(\omega) / n+o(1)$ tends to $c_{i}=\int \varphi_{r} \mathrm{~d} \mu_{i}$ by definition of $O_{i}$. This shows that $\xi \in U\left(c_{i}\right)$. Therefore, $\alpha_{*}\left(O_{i} \cap\left[E_{*}\right]\right) \subset U\left(c_{i}\right)$. With (3.10), this gives $\nu_{R}\left(U\left(c_{i}\right)\right)>0$.

Since $\nu_{R}$ is ergodic for the action of $\Gamma$ by Proposition 3.15, and the sets $U(c)$ are $\Gamma$-invariant, we deduce that $U\left(c_{i}\right)$ has full measure. Therefore, all those sets have to coincide.

Proof of Theorem 3.14. By Proposition 3.10, the pressure $\operatorname{Pr}_{i}\left(\varphi_{r}\right)$ on the component $\mathcal{C}_{i}$ is equal to $\int\left(\varphi_{r}-\varphi_{R}\right) \mathrm{d} \mu_{i}+O\left(\left\|\varphi_{r}-\varphi_{R}\right\|\right)^{2}$. The integral does not depend on $i$, by Proposition 3.16. Considering $i^{\prime}$ such that the pressure is maximal, we obtain

$$
\operatorname{Pr}_{i}\left(\varphi_{r}\right)=\operatorname{Pr}\left(\varphi_{r}\right)+O\left(\left\|\varphi_{r}-\varphi_{R}\right\|\right)^{2} .
$$

Moreover, by (3.5) and (3.6), the ratio between $\operatorname{Pr}\left(\varphi_{r}\right)$ and $-\sqrt{R-r}$ is bounded from above and below. Since $\left\|\varphi_{r}-\varphi_{R}\right\|^{2}=O(R-r)^{2 / 3}$ by Lemma 3.11, we obtain $\operatorname{Pr}_{i}\left(\varphi_{r}\right)=\operatorname{Pr}\left(\varphi_{r}\right)+o\left(\operatorname{Pr}\left(\varphi_{r}\right)\right)$. This concludes the proof.
3.6. Estimating the second derivative of the Green function. To improve on Proposition 3.2, one should get asymptotics for $\sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)$ in terms of $\eta(r)=\sum G_{r}(e, x) G_{r}(x, e)=\sum H_{r}(e, x)$ or, equivalently, in terms of $\operatorname{Pr}\left(\varphi_{r}\right)$.

Proposition 3.17. One has when $r \rightarrow R$

$$
\begin{equation*}
\sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)=c(r) \eta(r)^{3}+O\left(\eta(r)^{2}\right) \tag{3.11}
\end{equation*}
$$

for some nonnegative function $c(r)$ that extends continuously to $r=R$.
This subsection is devoted to the proof of this proposition. We will need to express things in terms of transfer operators on the symbolic space. Let

$$
\begin{equation*}
\Phi_{r}(x)=\frac{\sum_{y} G_{r}(e, y) G_{r}(y, x)}{G_{r}(e, x)} \tag{3.12}
\end{equation*}
$$

and define for $r<R$ a probability measure $\nu_{r}$ on $\Gamma$ by

$$
\begin{equation*}
\int f \mathrm{~d} \nu_{r}=\frac{\sum_{x \in \Gamma} H_{r}(e, x) f(x)}{\sum_{x \in \Gamma} H_{r}(e, x)} . \tag{3.13}
\end{equation*}
$$

The sum in (3.11) is equal to

$$
\eta(r) \int \Phi_{r} \mathrm{~d} \nu_{r}
$$

To estimate it, we should understand $\Phi_{r}$ and $\nu_{r}$.
Proposition 3.18. When $r \rightarrow R$, the sequence of probability measures $\nu_{r}$ on the compact space $\Gamma \cup \partial \Gamma$ converges weakly to a probability measure $\nu_{R}$, which is supported on $\partial \Gamma$.

Proof. If $\nu_{r}$ converges weakly, then the limiting measure can give no weight to $\Gamma$, since $\nu_{r}(x)=H_{r}(e, x) / \sum_{y \in \Gamma} H_{r}(e, y)$ tends to 0 by (3.5).

Therefore, we just have to prove the convergence of $\nu_{r}(f)$ for any continuous function, or even for $f$ in a dense set of functions. We will consider those $f$ such that the function $\tilde{f}$ defined on $\Sigma^{*}$ by $\tilde{f}(\omega)=f\left(\alpha_{*}(\omega)\right)$ belongs to $\mathcal{H}^{\beta}$. For such a function, we have

$$
\sum_{x \in \Gamma} H_{r}(e, x) f(x)=H_{r}(e, e) \sum_{n \in \mathbb{N}} \mathcal{L}_{\varphi_{r}}^{n}\left(1_{\left[E_{*}\right]} \tilde{f}\right)(\emptyset) .
$$

Using the spectral description of Proposition 3.10, we deduce that this can be written as

$$
\sum_{i=1}^{I} c_{\tilde{f}}(i, r) /\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|+O(1)
$$

as in (3.6), for some functions $c_{\tilde{f}}(i, r)$ that extend continuously up to $r=R$. Therefore, by (3.7),

$$
\nu_{r}(f)=\frac{\sum_{i=1}^{I} c_{\tilde{f}}(i, r) /\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|+O(1)}{\tau_{2}(r) /\left|\operatorname{Pr}\left(\varphi_{r}\right)\right|+O(1)}
$$

Since all the quantities $\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|$ are asymptotic to $\left|\operatorname{Pr}\left(\varphi_{r}\right)\right|$ by Theorem 3.14 and tend to 0 , this converges when $r$ tends to $R$, to $\sum_{i=1}^{I} c_{\tilde{f}}(i, R) / \tau_{2}(R)$.
Remark 3.19. One can easily check that the measure $\nu_{R}$ in Proposition 3.18 is the same as the measure we constructed in Subsection 3.5 and was already denoted by $\nu_{R}$. This will have no importance for our purposes.

To estimate $\Phi_{r}$ [defined in (3.12)], let us first note the following estimate.
Lemma 3.20. We have

$$
\Phi_{r}(x) \leqslant C(1+|x|) \eta(r) .
$$

Proof. The proof relies on the same argument as Lemma 3.4. Denote by $\gamma$ a geodesic segment from $e$ to $x$, and by $\Gamma_{n}$ (for $\left.0 \leqslant n \leqslant|x|\right)$ the set of points whose first projection on $\gamma$ is the point $\gamma(n)$, at distance $n$ of $e$. For $y \in \Gamma_{n}$, one has by Ancona inequalities

$$
\begin{aligned}
G_{r}(e, y) G_{r}(y, x) & \leqslant C G_{r}(e, \gamma(n)) G_{r}(\gamma(n), y) G_{r}(y, \gamma(n)) G_{r}(\gamma(n), x) \\
& \leqslant C H_{r}(\gamma(n), y) G_{r}(e, x) .
\end{aligned}
$$

Therefore,

$$
\Phi_{r}(x)=\sum_{n=0}^{|x|} \sum_{y \in \Gamma_{n}} \frac{G_{r}(e, y) G_{r}(y, x)}{G_{r}(e, x)} \leqslant C \sum_{n=0}^{|x|} \sum_{y \in \Gamma_{n}} H_{r}(\gamma(n), y) \leqslant C(|x|+1) \eta(r)
$$

To obtain a convergence instead of bounds, we will use a similar argument, but we will need to replace the wild sets $\Gamma_{n}$ by a nicer version given by partitions of unity, as in Lemma 8.5 of GL13] (that we recall for the convenience of the reader).
Lemma 3.21. For $K$ large enough, we can associate to any geodesic segment $\gamma$ in the Cayley graph of length $2 K+1$ centered around e a function $\kappa_{\gamma}: \Gamma \rightarrow[0,1]$ with the following properties:
(1) The function $\kappa_{\gamma}$ extends continuously to $\Gamma \cup \partial \Gamma$.
(2) Let $\pi_{\gamma}(y)$ be the set of points on $\gamma$ that are closest to $y \in \Gamma$. Then $\kappa_{\gamma}(y)=0$ if $\pi_{\gamma}(y)$ contains a point at distance $\geqslant K / 4$ of $e$.
(3) Let $\gamma^{\prime}$ be any bi-infinite geodesic passing through e. Adding the functions $\kappa_{\gamma}$ along the subsegments of $\gamma^{\prime}$ of length $2 K+1$ one gets the function identically equal to 1. More formally, for all $y \in \Gamma$,

$$
\begin{equation*}
\sum_{n \in \mathbb{Z}} \kappa_{\gamma^{\prime}(n)^{-1} \gamma^{\prime}[n-K, n+K]}\left(\gamma^{\prime}(n)^{-1} y\right)=1 \tag{3.14}
\end{equation*}
$$

Let us now define for $r \in[1, R)$ a function $\Psi_{r}$ on geodesic segments $\gamma$ through $e$, as follows. Let $a$ and $b$ be the endpoints of $\gamma$. If $d(e, a) \leqslant K$ or $d(e, b) \leqslant K$, let $\Psi_{r}(\gamma)=0$. Otherwise, let

$$
\Psi_{r}(\gamma)=\eta(r)^{-1} \sum_{y \in \Gamma} \kappa_{\gamma[-K, K]}(y) G_{r}(a, y) G_{r}(y, b) / G_{r}(a, b) .
$$

Consider a geodesic segment $\gamma$ from $e$ to a point $x$, and denote by $\sigma^{n} \gamma$ the shifted segment, i.e., $\gamma(n)^{-1} \gamma$. Then we have

$$
\begin{equation*}
\Phi_{r}(x)=\eta(r) \sum_{n=0}^{|x|} \Psi_{r}\left(\sigma^{n} \gamma\right)+O(\eta(r)) \tag{3.15}
\end{equation*}
$$

Indeed, by (3.14), when one adds all the quantities $\Psi_{r}\left(\sigma^{n} \gamma\right)$, one counts every point in the group with a coefficient 1 , except those whose projection on $\gamma$ is close to $e$ or $x$. They contribute to the sum by an amount at most $C \eta(r)$, as explained in the proof of Lemma 3.20

Lemma 3.22. The functions $\Psi_{r}$ are uniformly bounded and Hölder-continuous for $r \in[1, R)$. They converge uniformly when $r$ tends to $R$.

By Hölder-continuous, we mean that, if two geodesics $\gamma$ and $\gamma^{\prime}$ coincide on a ball of size $n$ around $e$, then $\left|\Psi_{r}(\gamma)-\Psi_{r}\left(\gamma^{\prime}\right)\right| \leqslant C e^{-\rho n}$ for some $\rho>0$.

Proof. This is essentially Lemma 8.6 in [GL13]. The uniform Hölder continuity is proved there and relies uniquely on strong Ancona inequalities. On the other hand, the proof of the convergence when $r \rightarrow R$ has to be modified slightly due to the presence of several maximal components.

Since the functions $\Psi_{r}$ are uniformly Hölder-continuous, it is sufficient to show that they converge simply to get uniform convergence. Fix some geodesic segment $\gamma$ through $e$, with endpoints $a$ and $b$ at distance at least $K$ of $e$. We have

$$
G_{r}(a, b) \Psi_{r}(\gamma)=\frac{1}{\eta(r)} \sum_{y \in \Gamma} \kappa_{\gamma[-K, K]}(y) \frac{G_{r}(a, y) G_{r}(y, b)}{G_{r}(e, y) G_{r}(y, e)} H_{r}(e, y)=\int f_{r}(y) \mathrm{d} \nu_{r},
$$

where $f_{r}(y)=\kappa_{\gamma[-K, K]}(y) \frac{G_{r}(a, y) G_{r}(y, b)}{G_{r}(e, y) G_{r}(y, e)}$. This is a function on $\Gamma$ that extends continuously to $\partial \Gamma$ by the strong Ancona inequalities (and since $\kappa_{\gamma[-K, K]}$ is continuous). Moreover, $f_{r}$ converges uniformly to a function $f_{R}$ when $r$ tends to $R$, by Lemma 3.4 Since $\nu_{r}$ converges weakly by Proposition 3.18, it follows that $\int f_{r} \mathrm{~d} \nu_{r}$ converges.

Lemma 3.23. There exists a family of functions $h_{r}$ on $\bar{\Sigma}$ for $r \in[1, R)$ with the following properties:
(1) The functions $h_{r}$ are Hölder-continuous, and they converge in the Hölder topology to a function $h_{R}$ when $r \rightarrow R$.
(2) For any $\omega \in \Sigma^{*}$ of length $n$,

$$
\begin{equation*}
\Phi_{r}\left(\alpha_{*}(\omega)\right)=\eta(r) S_{n} h_{r}(\omega)+O(\eta(r)) . \tag{3.16}
\end{equation*}
$$

We recall that $S_{n} h_{r}$ is the Birkhoff sum $\sum_{k=0}^{n-1} h_{r} \circ \sigma^{k}$.
Proof. We will need to work with the bilateral shift $\sigma_{\mathbb{Z}}$ on the space $\bar{\Sigma}_{\mathbb{Z}}$ of bilateral paths in the automaton (that may be infinite in zero, one, or both directions). We define a function $g_{r}$ on the set $\Sigma_{\mathbb{Z}}^{*}$ of finite paths by $g_{r}(\omega)=0$ if the length of $\omega$ in the future or in the past is less than $K$, and $g_{r}(\omega)=\Psi_{r}(\alpha(\omega))$ otherwise, where $\alpha(\omega)$ is the geodesic segment going through

$$
\ldots, \alpha\left(\omega_{-1}\right)^{-1} \alpha\left(\omega_{-2}\right)^{-1}, \alpha\left(\omega_{-1}\right)^{-1}, e, \alpha\left(\omega_{0}\right), \alpha\left(\omega_{0}\right) \alpha\left(\omega_{1}\right), \ldots
$$

Lemma 3.22 ensures that the functions $g_{r}$ are Hölder-continuous and that they converge uniformly when $r$ tends to $R$. By (3.1), they also converge in some Hölder topology. Moreover, they extend to Hölder-continuous functions on $\bar{\Sigma}_{\mathbb{Z}}$.

Consider now a finite path $\omega$ in $\Sigma^{*}$, of length $n$. One may consider it as a path in $\Sigma_{\mathbb{Z}}^{*}$ with empty coordinates for negative time. Equation (3.15) reads

$$
\Phi_{r}\left(\alpha_{*}(\omega)\right)=\eta(r) \sum_{k=0}^{n} g_{r}\left(\sigma_{\mathbb{Z}}^{k} \omega\right)+O(\eta(r)) .
$$

This is almost the required property, but the function $g_{r}$ is defined on the bilateral shift instead of the unilateral shift as desired. This problem is solved using a classical coboundary trick: for any Hölder-continuous function $g$ on $\bar{\Sigma}_{\mathbb{Z}}$, there exist two Hölder-continuous functions $h$ and $u$ on $\bar{\Sigma}_{\mathbb{Z}}$ (for a smaller Hölder exponent) such that $g=h+u-u \circ \sigma_{\mathbb{Z}}$, and $h$ depends only on positive coordinates. Moreover, $h$ and $u$ depend linearly (and continuously) on $g$. This is Proposition 1.2 in [PP90. The proof is given there for subshifts where one only allows infinite paths, but it readily adapts to the situation where finite paths are allowed (or one can reduce to
the infinite paths situation by adding two cemeteries, one for the past and one for the future).

Writing $g_{r}=h_{r}+u_{r}-u_{r} \circ \sigma_{\mathbb{Z}}$ as above, we obtain

$$
\Phi_{r}\left(\alpha_{*}(\omega)\right)=\eta(r) \sum_{k=0}^{n} h_{r}\left(\sigma_{\mathbb{Z}}^{k} \omega\right)+\eta(r)\left(u_{r}(\omega)-u_{r}\left(\sigma_{\mathbb{Z}}^{n+1} \omega\right)\right)+O(\eta(r)) .
$$

Since $u_{r}$ is uniformly bounded, this is the desired decomposition.
Proof of Proposition 3.17. The sum in (3.11) can be written as $\sum_{x \in \Gamma} H_{r}(e, x) \Phi_{r}(x)$. Since $\alpha_{*}$ induces a bijection between the finite paths in the automaton starting from $s_{*}$ and the group, this is equal to $\sum H_{r}\left(e, \alpha_{*}(\omega)\right) \Phi_{r}\left(\alpha_{*}(\omega)\right)$, where the sum is restricted to those paths with $\omega_{0} \in E_{*}$.

Consider now, in this sum, the contribution of paths of length $n$. By definition of the transfer operator $\mathcal{L}_{\varphi_{r}}$, it is equal to $H_{r}(e, e) \mathcal{L}_{\varphi_{r}}^{n}\left(1_{\left[E_{*}\right]} \cdot \Phi_{r} \circ \alpha_{*}\right)(\emptyset)$. Using the decomposition (3.16) for $\Phi_{r} \circ \alpha_{*}$, we get

$$
\sum_{x, y \in \Gamma} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)=\eta(r) \sum_{n \in \mathbb{N}} \mathcal{L}_{\varphi_{r}}^{n}\left(1_{\left[E_{*}\right]} S_{n} h_{r}+O(1)\right)(\emptyset) .
$$

The contribution of the error term $O(1)$ in this equation is bounded by

$$
\eta(r) \sum\left\|\mathcal{L}_{\varphi_{r}}^{n} 1\right\| \leqslant C \eta(r) /\left|\operatorname{Pr}\left(\varphi_{r}\right)\right| \leqslant C \eta(r)^{2} .
$$

It is therefore compatible with the error term in the statement of Proposition 3.17
Since $\mathcal{L}_{\varphi_{r}}(u \cdot v \circ \sigma)=v \mathcal{L}_{\varphi_{r}} u$, we have $\mathcal{L}_{\varphi_{r}}^{n}\left(1_{\left[E_{*}\right]} S_{n} h_{r}\right)=\sum_{k=1}^{n} \mathcal{L}_{\varphi_{r}}^{k}\left(h_{r} \mathcal{L}_{\varphi_{r}}^{n-k} 1_{\left[E_{*}\right]}\right)$. Therefore, the previous equation becomes

$$
\begin{aligned}
& \eta(r) \sum_{n=0}^{\infty} \sum_{k=1}^{n} \mathcal{L}_{\varphi_{r}}^{k}\left(h_{r} \mathcal{L}_{\varphi_{r}}^{n-k} 1_{\left[E_{*}\right]}\right)(\emptyset)+O\left(\eta(r)^{2}\right) \\
& \quad=\eta(r) \sum_{k=1}^{\infty} \mathcal{L}_{\varphi_{r}}^{k}\left(h_{r} \sum_{\ell=0}^{\infty} \mathcal{L}_{\varphi_{r}}^{\ell} 1_{\left[E_{*}\right]}\right)(\emptyset)+O\left(\eta(r)^{2}\right)
\end{aligned}
$$

Using the spectral description of Proposition 3.10, one can write

$$
\sum_{\ell=0}^{\infty} \mathcal{L}_{\varphi_{r}}^{\ell} 1_{\left[E_{*}\right]}=\sum_{i=1}^{I} f_{i, r} /\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|+O(1)
$$

for some Hölder-continuous functions $f_{i, r}$ that converge when $r$ tends to $R$. Again, the $O(1)$ results in an error $O\left(\eta(r)^{2}\right)$ in the final formula. It remains to understand $\sum_{k=1}^{\infty} \mathcal{L}_{\varphi_{r}}^{k}\left(h_{r} f_{i, r}\right)$. Using again the spectral description of $\mathcal{L}_{\varphi_{r}}$, one may write it as $\sum_{j=1}^{I} g_{i, j, r} /\left|\operatorname{Pr}_{j}\left(\varphi_{r}\right)\right|$ for some Hölder-continuous functions $g_{i, j, r}$ that depend continuously on $r$. Finally, we have obtained

$$
\sum_{x, y \in \Gamma} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)=\eta(r) \sum_{i=1}^{I} \sum_{j=1}^{I} \frac{g_{i, j, r}(\emptyset)}{\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|\left|\operatorname{Pr}_{j}\left(\varphi_{r}\right)\right|}+O\left(\eta(r)^{2}\right)
$$

By Theorem 3.14 and (3.7), $\left|\operatorname{Pr}_{i}\left(\varphi_{r}\right)\right|$ is asymptotic to $\tau_{2}(r) / \eta(r)$ for some continuous function $\tau_{2}$ that admits a positive limit at $r=R$. Since all the quantities $g_{i, j, r}(\emptyset)$ converge when $r$ tends to $R$, the proposition follows.
3.7. Asymptotics of the Green function. In this subsection, we prove Theorem 3.1. We rely on the asymptotics for $\sum G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)$ that were obtained in Proposition 3.17 and the differential equation for $G_{r}(e, e)$.

More precisely, define as in the proof of Corollary 3.3a function $F(r)=r^{2} \eta(r)$, so that $F^{\prime}(r)=2 r \sum_{x, y} G_{r}(e, y) G_{r}(y, x) G_{r}(x, e)$. By Proposition 3.17, $2 F^{\prime}(r) / F(r)^{3}$ converges to a constant $c$ when $r$ tends to $R$. By Proposition 3.2, $c$ is nonzero. By integration, it follows that $1 / F(r)^{2}-1 / F(R)^{2} \sim c(R-r)$. Since $F(R)=+\infty$, we get $F(r) \sim c^{-1 / 2}(R-r)^{-1 / 2}$. This proves the desired asymptotics of $\partial G_{r}(e, e) / \partial r$.

Fix now a point $a \in \Gamma$, and let us compute the asymptotics of $\partial G_{r}(e, a) / \partial r$ or, equivalently, of $\sum_{x} G_{r}(e, x) G_{r}(x, a)$. We write this sum as

$$
\sum_{x} H_{r}(e, x) \frac{G_{r}(x, a)}{G_{r}(x, e)}=\eta(r) \int_{\Gamma} f_{r} \mathrm{~d} \nu_{r}
$$

where $f_{r}(x)=G_{r}(x, a) / G_{r}(x, e)$ and the measure $\nu_{r}$ has been defined in (3.13). The functions $f_{r}$ extend continuously to $\Gamma \cup \partial \Gamma$ by the strong Ancona inequalities and converge uniformly when $r$ tends to $R$, by Lemma 3.4. Since $\nu_{r}$ converges weakly when $r \rightarrow R$ by Proposition 3.18, we deduce that $\int f_{r} \mathrm{~d} \nu_{r}$ converges. Therefore, the asymptotics of $\partial G_{r}(e, a) / \partial r$ follow from those of $\eta(r)$.

## 4. Asymptotics of transition probabilities

Theorem 1.1 follows directly from the asymptotics of the Green function proved in Theorem 3.1 and from Theorem 9.1 in (GL13): this theorem shows that, for symmetric measures, one can read the behavior of transition probabilities from the behavior of the Green function.

Let us explain quickly why symmetry matters. The Green function is $G_{r}(x, y)=$ $\sum r^{n} p_{n}(x, y)$, and its derivative is $\sum n r^{n-1} p_{n}(x, y)$. If $\partial G_{r}(x, y) / \partial r \sim C(R-$ $r)^{-1 / 2}$, it follows from Karamata's tauberian theorem that

$$
\sum_{k=1}^{n} k R^{k} p_{k}(x, y) \sim C^{\prime} n^{1 / 2}
$$

This is a local limit theorem in Cesaro average. If $R^{n} p_{n}(x, y)$ were monotone, the desired asymptotics of $p_{n}(x, y)$ would follow readily. Symmetry is used to obtain almost monotonicity: up to an exponentially small error (that does not matter in the estimates), $R^{n} p_{n}(x, y)$ is indeed decreasing when the random walk is aperiodic and the measure is symmetric. This is a consequence of spectral results for the (self-adjoint) Markov operator associated to the random walk.

From Theorem 1.1, one can also derive asymptotics for the first return probabilities. We describe the result in the aperiodic case; the periodic one is handled similarly by looking at $\mu^{2}$.

Proposition 4.1. Consider a probability measure $\mu$ on a countable group $\Gamma$ such that the associated transition probabilities satisfy $p_{n}(x, y) \sim C(x, y) R^{-n} n^{-\beta}$ for some $R \geqslant 1$ and $\beta>1$. Let $f_{n}(x, y)$ be the first visit probabilities from $x$ to $y$ at time $n$, i.e.,

$$
f_{n}(x, y)=\mathbb{P}_{x}\left(X_{1}, \ldots, X_{n-1} \neq y, X_{n}=y\right) .
$$

Then $f_{n}(x, y) \sim C^{\prime}(x, y) R^{-n} n^{-\beta}$ for some constants $C^{\prime}(x, y)$.
For the proof, we will mainly rely on the following theorem [CNW73, Theorem 1].

Theorem 4.2. Consider a function $A(z)=\sum_{n=0}^{\infty} a_{n} z^{n}$ with $a_{n} \geqslant 0$ and $a_{n} \sim$ $c R^{-n} n^{-\beta}$ with $\beta>1$. Consider also a function $\Phi$ that is analytic on a neighborhood of $\{A(z):|z| \leqslant R\}$. Then the coefficients $b_{n}$ of the series expansion $\Phi(A(z))=$ $\sum b_{n} z^{n}$ satisfy

$$
b_{n} \sim a_{n} \Phi^{\prime}\left(\sum a_{n} R^{n}\right) .
$$

Proof of Proposition 4.1. Decomposing a path from $e$ to itself into successive excursions, one gets the renewal equation

$$
\begin{equation*}
\sum_{n=0}^{\infty} p_{n}(e, e) z^{n}=\frac{1}{1-\sum_{n=1}^{\infty} f_{n}(e, e) z^{n}} \tag{4.1}
\end{equation*}
$$

Since $\sum p_{n}(e, e) R^{n}=G_{R}(e, e)<\infty$, one deduces $\sum f_{n}(e, e) R^{n}<1$. Therefore, $1-\sum f_{n}(e, e) z^{n}$ does not vanish for $|z| \leqslant R$, and $\sum p_{n}(e, e) z^{n}$ is well-defined and nonzero for any such $z$. From (4.1), one gets

$$
\sum f_{n}(e, e) z^{n}=1-1 /\left(\sum p_{n}(e, e) z^{n}\right)=\Phi\left(\sum p_{n}(e, e) z^{n}\right)
$$

where we set $\Phi(t)=1-1 / t$. Since $p_{n}(e, e) \sim C R^{-n} n^{-\beta}$, we may apply Theorem4.2 to obtain $f_{n}(e, e) \sim C^{\prime} R^{-n} n^{-\beta}$.

For $x \neq y$, one has $\sum p_{n}(x, y) z^{n}=\left(\sum f_{n}(x, y) z^{n}\right) \cdot\left(\sum p_{n}(e, e) z^{n}\right)$. The functions $\sum p_{n}(x, y) z^{n}$ and $1 / \sum p_{n}(e, e) z^{n}$ both have coefficients that are asymptotic to a constant times $R^{-n} n^{-\beta}$. Since the set of all functions with this property is closed under multiplication (see CNW73, Lemma 1]), we get $f_{n}(x, y) \sim C^{\prime}(x, y) R^{-n} n^{-\beta}$ as desired.

## Appendix A. Ancona inequalities for surface groups

In this appendix, we prove Ancona inequalities for surface groups without any symmetry assumption on the measure.

Theorem A.1. Let $\Gamma$ be a cocompact Fuchsian group, i.e., a cocompact discrete subgroup of $\operatorname{PSL}(2, \mathbb{R})$. Let $\mu$ be an admissible finitely supported probability measure on $\Gamma$. Then it satisfies strong uniform Ancona inequalities.

This theorem has several corollaries.
Corollary A.2. Under the assumptions of the theorem, the Martin boundary for $R$-harmonic functions coincides with the geometric boundary of the group, i.e., the unit circle $S^{1}$.

Corollary A.3. Under the assumptions of the theorem, for any $x, y \in \Gamma$, there exists $C(x, y)>0$ such that $\sum_{k=1}^{n} k R^{k} p_{k}(x, y) \sim C(x, y) n^{1 / 2}$.

The first corollary is a classical consequence of Ancona inequalities. For the second corollary, we rely on Section 3 (or on GL13, since the Cannon automaton is transitive) to deduce that the Green function satisfies $\partial G_{r}(x, y) / \partial r \sim C(x, y) /$ $\sqrt{R-r}$ when $r$ tends to $R$. Using Karamata's tauberian theorem, this readily implies the statement of the corollary (see the arguments in Section (4). Note that we are unable to deduce the true local limit theorem from this estimate since we do not know if $R^{n} p_{n}(x, y)$ is decreasing, or sufficiently well approximated by a decreasing sequence, as in the symmetric situation.

The strong uniform Ancona inequalities of Theorem A. 1 are a consequence of estimates for the weight of paths avoiding a ball, given in the following proposition, and of Lemma 2.7,

Proposition A.4. Let $\Gamma$ be a cocompact Fuchsian group in $\operatorname{PSL}(2, \mathbb{R})$. Let $\mu$ be an admissible finitely supported probability measure on $\Gamma$. For any $K>0$, there exists $n_{0}>0$ such that, for any $n \geqslant n_{0}$, for any points $x, y, z$ on a geodesic segment (in this order) with $d(x, y) \in[n, 100 n]$ and $d(y, z) \in[n, 100 n]$,

$$
G_{R}\left(x, z ; B(y, n)^{c}\right) \leqslant K^{-n} .
$$

The rest of this section is devoted to the proof of the proposition. The main tool in this proof is superadditivity (as in Lemma [2.6 which relies on Lemma 2.5). This time, it is in the form of Kingman's subadditive ergodic theorem, or rather a bilateral version of this theorem that we now give.

Theorem A.5. Let $T: \Omega \rightarrow \Omega$ be an ergodic probability preserving invertible automorphism of a probability space $(\Omega, \mathbb{P})$. Consider for each bounded interval $I \subset \mathbb{Z}$ an integrable function $\Phi_{I}: \Omega \rightarrow \mathbb{R}$ with the following properties:
(1) If I is the disjoint union of two intervals $I_{1}$ and $I_{2}$, then $\Phi_{I}(\omega) \leqslant \Phi_{I_{1}}(\omega)+$ $\Phi_{I_{2}}(\omega)$.
(2) One has $\Phi_{[m, n]}(\omega)=\Phi_{[m-1, n-1]}(T \omega)$.
(3) The quantity $n^{-1} \int \Phi_{[0, n)}(\omega) \mathrm{d} \mathbb{P}(\omega)$ is bounded from below.

Then, for almost every $\omega$, the quantity $(m+n)^{-1} \Phi_{[-m, n)}(\omega)$ converges when $m+n \rightarrow$ $\infty$ (and $m, n \geqslant 0$ ) toward $\inf n^{-1} \int \Phi_{[0, n)}(\omega) \mathrm{d} \mathbb{P}(\omega)=\lim n^{-1} \int \Phi_{[0, n)}(\omega) \mathrm{d} \mathbb{P}(\omega)$.
Proof. Let $\Psi_{n}(\omega)=\Phi_{[0, n)}(\omega)$. The assumptions give, for any $m, n \geqslant 0$,

$$
\begin{aligned}
\Psi_{m+n}(\omega) & =\Phi_{[0, m+n)}(\omega) \leqslant \Phi_{[0, n)}(\omega)+\Phi_{[n, n+m)}(\omega)=\Phi_{[0, n)}(\omega)+\Phi_{[0, m)}\left(T^{n} \omega\right) \\
& =\Psi_{n}(\omega)+\Psi_{m}\left(T^{n} \omega\right)
\end{aligned}
$$

This shows that $\Psi_{n}$ is a subadditive cocycle in the usual sense of Kingman's ergodic theorem (see for instance Kre85, Theorem I.5.3]). Therefore, $n^{-1} \Psi_{n}$ converges almost surely and in $L^{1}$ to the limit $c=\inf n^{-1} \int \Phi_{[0, n)}(\omega) \mathrm{dP}(\omega)=$ $\lim n^{-1} \int \Phi_{[0, n)}(\omega) d \mathbb{P}(\omega)$.

In the same way, $\Phi_{[-n,-1]}(\omega)$ is a subadditive cocycle for the transformation $T^{-1}$. Hence, $n^{-1} \Phi_{[-n,-1]}(\omega)$ converges almost surely to $\lim n^{-1} \int \Phi_{[-n,-1]}(\omega) d \mathbb{P}(\omega)$. Changing variables by $\omega^{\prime}=T^{-n} \omega$, this integral is equal to $\int \Phi_{[0, n)}\left(\omega^{\prime}\right) d \mathbb{P}\left(\omega^{\prime}\right)$. Therefore, the limit is again $c$.

Consider now a generic point $\omega$, and $m, n \geqslant 0$. We want to show that if $m+n$ is large, then $(m+n)^{-1} \Phi_{[-m, n)}(\omega)$ is close to $c$. We will do so if $m$ is large. The case $n$ large is handled similarly. Let $\varepsilon>0$ be small. Since $j^{-1} \Phi_{[0, j)}\left(\omega^{\prime}\right)$ converges almost everywhere to $c$, it converges uniformly on a set $A$ of a measure arbitrarily close to 1 . In particular, there exists $N$ such that, for all $j \geqslant N$ and for all $\omega^{\prime} \in A$, one has $j^{-1} \Phi_{[0, j)}\left(\omega^{\prime}\right) \in[c-\varepsilon, c+\varepsilon]$. Since $\omega$ is generic and $\mathbb{P}(A)$ is very close to 1 , the orbit of $\omega$ spends a very large proportion of its time in $A$. In particular, one may find for every large enough $m$ an integer $k \in[m+\varepsilon m, m+2 \varepsilon m]$ such that $T^{-k} \omega \in A$. We get for this $k$ (and for any $n \geqslant 0$ )

$$
\begin{aligned}
\Phi_{[0, k+n)}\left(T^{-k} \omega\right) & \leqslant \Phi_{[0, k-m)}\left(T^{-k} \omega\right)+\Phi_{[k-m, k+n)}\left(T^{-k} \omega\right) \\
& =\Phi_{[0, k-m)}\left(T^{-k} \omega\right)+\Phi_{[-m, n)}(\omega) .
\end{aligned}
$$

If $m$ is large enough, then $k-m \geqslant \varepsilon m$ is larger than $N$. Since $T^{-k} \omega$ belongs to $A$, we obtain

$$
\Phi_{[-m, n)}(\omega) \geqslant(k+n)(c-\varepsilon)-(k-m)(c+\varepsilon)=(m+n)(c+O(\varepsilon)) .
$$

This shows that $\liminf (m+n)^{-1} \Phi_{[-m, n)}(\omega) \geqslant c$. On the other hand,

$$
\Phi_{[-m, n)}(\omega) \leqslant \Phi_{[-m,-1]}(\omega)+\Phi_{[0, n)}(\omega)=m(c+o(1))+n(c+o(1))
$$

Therefore, $\lim \sup (m+n)^{-1} \Phi_{[-m, n)}(\omega) \leqslant c$. This concludes the proof.
Let us now start the proof of Proposition A.4 We can assume without loss of generality that $y=e$. Let $\ell>0$ be large (it will not depend on $n$ ). We will construct $\ell$ suitable barriers $A_{1}, \ldots, A_{\ell}$ between $x$ and $z$ such that

$$
\begin{equation*}
\sum_{a \in A_{i}, b \in A_{i+1}} G_{R}(a, b)^{2} \leqslant e^{-\rho n} \tag{A.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\sum_{a \in A_{1}} G_{R}(e, a)^{2} \leqslant 1, \quad \sum_{a \in A_{\ell}} G_{R}(a, e)^{2} \leqslant 1 \tag{A.2}
\end{equation*}
$$

for some $\rho>0$ that does not depend on $\ell, x$, or $z$. From the last equation, we obtain $\sum_{a \in A_{1}} G_{R}(x, a)^{2} \leqslant C^{n}$ thanks to Harnack inequalities [and since $d(x, e) \leqslant$ $100 n]$, and $\sum_{a \in A_{\ell}} G_{R}(a, z)^{2} \leqslant C^{n}$. Arguing as in the beginning of the proof of Lemma [2.6] we define operators $L_{i}$ from $\ell^{2}\left(A_{i+1}\right)$ to $\ell^{2}\left(A_{i}\right)$. They satisfy $\left\|L_{0}\right\| \leqslant$ $C^{n / 2},\left\|L_{\ell+1}\right\| \leqslant C^{n / 2}$, and $\left\|L_{i}\right\| \leqslant e^{-\rho n / 2}$ for $1 \leqslant i \leqslant \ell$. Therefore,

$$
G_{R}\left(x, z ; B(e, n)^{c}\right) \leqslant \prod\left\|L_{i}\right\| \leqslant C^{n} e^{-(\ell-1) \rho n / 2}
$$

Taking $\ell$ large, we can ensure that this is bounded by $K^{-n}$ as desired, for any $K>0$.

The key point of the argument is the construction of the barriers. The problem with the argument in Lemma 2.6 is that we only have a control on $G_{R}(a, b) G_{R}(b, a)$ coming from Lemma [2.5, not $G_{R}(a, b)^{2}$. The idea is that those controls would be equivalent if $G_{R}(a, b)$ and $G_{R}(b, a)$ were of the same order of magnitude. For symmetric measures, this is always the case. For nonsymmetric measures, we will be able to enforce it by constructing the barriers using another, symmetric, random walk. We will use the Kingman subadditive ergodic theorem to show that for typical points both $G_{R}(a, b)$ and $G_{R}(b, a)$ grow at the same speed. It will then follow from Lemma 2.5 that they are both exponentially small.

Let us stress that this kind of argument cannot work for all points. For instance, consider in the free group on two generators $a$ and $b$ a random walk that goes toward $a$ with probability $1-3 \varepsilon$, and toward $a^{-1}, b$, and $b^{-1}$ with probability $\varepsilon$, for some small enough $\varepsilon$. It is easy to check that $G_{R}\left(e, a^{n}\right)$ is exponentially large [while $G_{R}\left(a^{n}, e\right)$ is exponentially small]. In particular, $\sum_{x \in \mathbb{S}_{n}} G_{R}(e, x)^{2}$ grows exponentially fast, but this growth is due to a rather small number of points. The barriers we construct have to avoid those points.

We turn to details. The barriers we will construct will not depend on the points $x$ and $z$ (but the order in which they will be encountered will depend on those points, of course). Since $\Gamma$ is a cocompact discrete subgroup of $\operatorname{PSL}(2, \mathbb{R})$, it acts on the hyperbolic disk $\mathbb{H}^{2}$. If $O$ is a suitably chosen reference point in this disk, the points $\gamma O$ (for $\gamma \in \Gamma$ ) are pairwise disjoint; hence $\Gamma$ can be identified with $\Gamma O$. Moreover, this identification is a quasi-isometry between $\Gamma$ (with the word distance
coming from its Cayley graph) and $\mathbb{H}^{2}$. In particular, the geometric boundary of $\Gamma$ is identified with $S^{1}=\partial \mathbb{H}^{2}$. We can assume that $O$ is the center of the hyperbolic disk.

Let us fix an admissible symmetric measure $\nu$ on $\Gamma$, supported on the set of generators, and let us consider the corresponding random walk. We claim that the following lemma holds. Here and henceforth, $G_{R}$ always denotes the Green function associated to the original measure $\mu$.

Lemma A.6. There exist $\rho>0$ and $v>0$ with the following properties.
(1) For almost every trajectory $X_{k}$ of the random walk given by $\nu, d\left(X_{k}, e\right) \sim$ $k v$.
(2) For almost every trajectory $X_{k}$, for all large enough $k, G_{R}\left(e, X_{k}\right) \leqslant e^{-\rho k}$ and $G_{R}\left(X_{k}, e\right) \leqslant e^{-\rho k}$.
(3) For almost every pair of independent trajectories $X_{k}$ and $Y_{k}$, for all large enough $k$ and $k^{\prime}, G_{R}\left(X_{k}, Y_{k^{\prime}}\right) \leqslant e^{-\rho\left(k+k^{\prime}\right)}$.
Let us admit the lemma for the moment. We choose $M \ell$ points in $S^{1}$ that are evenly spaced (for some large enough $M$ ), and $M \ell$ small intervals $I_{i}$ around those points. The Poisson boundary of the random walk given by $\nu$ is $S^{1}$, and the hitting measure has full support. Therefore, there is positive probability to hit the boundary in any of the intervals $I_{i}$. Let us choose for each $i$ a trajectory $X_{k}^{(i)}$ of the random walk that ends up in $I_{i}$. We will also require each of those trajectories to be typical, so that they satisfy the conclusions of Lemma A. 6 ,

Since the trajectories $X_{k}^{(i)}$ converge to different points on the boundaries, they are disjoint outside of a large enough compact set. Let $\Gamma_{i}(n)$ be the set of points $X_{k}^{(i)}$ that are at a distance at least $n$ of $e$, and let $B_{i}(n)$ be a thickening of $\Gamma_{i}(n)$, i.e., $B_{i}(n)=\bigcup_{a \in \Gamma_{i}(n)} B\left(a, C_{0}\right)$ for some large constant $C_{0}$. If $n$ is large enough, the sets $\left(B_{i}(n)\right)_{i \leqslant 2 \ell+2}$ are mutually disjoint. Since $d\left(X_{k}^{(i)}, e\right) \sim v k$, the set $\Gamma_{i}(n)$ only contains points among the $X_{k}^{(i)}$ with $k \geqslant n /(2 v)$. Therefore,

$$
\sum_{a \in \Gamma_{i}(n)} G_{R}(e, a)^{2} \leqslant \sum_{k=n /(2 v)}^{\infty} G_{R}\left(e, X_{k}^{(i)}\right)^{2} \leqslant \sum_{k=n /(2 v)}^{\infty} e^{-2 \rho k} \leqslant C e^{-\rho n / v}
$$

Since any point in the thickening $B_{i}(n)$ is a bounded distance away from a point in $\Gamma_{i}(n)$, a similar estimate holds with $B_{i}(n)$ instead of $\Gamma_{i}(n)$ thanks to Harnack inequalities. Arguing in the same way for the other inequalities, we obtain

$$
\begin{align*}
\sum_{a \in B_{i}(n)} G_{R}(e, a)^{2} & \leqslant C e^{-\rho n / v}, \quad \sum_{a \in B_{i}(n)} G_{R}(a, e)^{2} \leqslant C e^{-\rho n / v}, \\
\sum_{a \in B_{i}(n), b \in B_{j}(n)} G_{R}(a, b)^{2} & \leqslant C e^{-2 \rho n / v} . \tag{A.3}
\end{align*}
$$

Consider now two points $x$ and $z$ at a distance at least $n$ of $e$ such that $e$ is on a geodesic segment from $x$ to $z$. The Cayley geodesic from $x$ to $z$ is a quasigeodesic in hyperbolic space, which remains in a bounded size neighborhood of a true hyperbolic geodesic (and this geodesic passes close to $O$ ). In particular, the visual angle from $O$ between $x$ and $z$ is bounded from below. Avoiding a ball around $O$ in $\mathbb{H}^{2}$, one can go from $x$ to $z$ in two directions around this ball, clockwise or counterclockwise. Since the limit intervals $I_{i}$ are evenly spaced, it follows that, in any of those directions, one meets successively at least $\ell$ sets $B_{i}(n)$ if $M$ is large
enough. Denote by $A_{1}$ the union of the two sets $B_{i}(n)$ that are closest to $x$ (ignoring the single set that might contain $x$ ), then by $A_{2}$ the union of the next two ones, and so on. We get barriers $A_{1}, \ldots, A_{\ell}$ between $x$ and $z$ as desired. Moreover, the estimates (A.3) show that those barriers satisfy (A.1) and (A.2) if $n$ is large enough (for a different value of $\rho$ ). This concludes the proof of Proposition A.4 modulo Lemma A. 6

Proof of Lemma A.6. We will use Kingman's theorem on the space $\Omega=\Gamma^{\mathbb{Z}}$ with the product measure $\nu^{\otimes \mathbb{Z}}$. In other words, an element $\omega \in \Omega$ is a sequence of elements $\omega_{i}$ of $\Gamma$ that are drawn independently according to $\nu$. Such an $\omega$ can be viewed as the increments of a random walk distributed according to $\nu$ : let $X_{n}(\omega)=\omega_{0} \cdots \omega_{n-1}$ for $n \geqslant 0$ and $X_{n}(\omega)=\omega_{-1}^{-1} \cdots \omega_{n}^{-1}$ for $n<0$, so that $X_{0}=e$ and $X_{n+1}=X_{n} \omega_{n}$. Let $T$ be the left shift on $\Omega$; it is ergodic, it preserves the measure, and $X_{n}(T \omega)=\omega_{0}^{-1} X_{n+1}(\omega)$.

We define a subadditive cocycle $\Phi_{[m, n)}(\omega)=-\log F_{R}\left(X_{m}(\omega), X_{n}(\omega)\right)$, where $F_{R}(x, y)=G_{R}(x, y) / G_{R}(e, e)$ is the first entrance Green function for the measure $\mu$, defined in Subsection 2.1. This function satisfies $F_{R}(x, y) F_{R}(y, z) \leqslant F_{R}(x, z)$ by (2.3); hence $\Phi_{[m, n)}$ is subadditive. By Harnack inequalities, $\left|\Phi_{[m, n)}(\omega)\right| \leqslant$ $C|n-m|$. Therefore, the integrability assumptions of Theorem A. 5 are satisfied. We deduce that $(m+n)^{-1} \log F_{R}\left(X_{-m}, X_{n}\right)$ converges almost surely to $c=$ $\lim n^{-1} \int \log F_{R}\left(e, X_{n}(\omega)\right) \mathrm{d} \mathbb{P}(\omega)$. We can also apply Kingman's theorem to the cocycle $\tilde{\Phi}_{[m, n)}(\omega)=-\log F_{R}\left(X_{n}(\omega), X_{m}(\omega)\right)$, to get that $(m+n)^{-1} \log F_{R}\left(X_{n}, X_{-m}\right)$ almost surely converges, to $c^{\prime}=\lim n^{-1} \int \log F_{R}\left(X_{n}(\omega), e\right) d \mathbb{P}(\omega)$. By left-invariance of the Green function, $F_{R}\left(X_{n}, e\right)=F_{R}\left(e, X_{n}^{-1}\right)$. Since $X_{n}^{-1}$ is distributed like $X_{n}$ by symmetry of the random walk given by $\nu$, this yields $\int \log F_{R}\left(X_{n}(\omega), e\right) \mathrm{d} \mathbb{P}(\omega)=$ $\int \log F_{R}\left(e, X_{n}(\omega)\right) \mathrm{d} \mathbb{P}(\omega)$. Dividing by $n$ and letting $n$ tend to infinity gives $c=c^{\prime}$.

The quantities $(m+n)^{-1} \log F_{R}\left(X_{-m}, X_{n}\right)$ and $(m+n)^{-1} \log F_{R}\left(X_{n}, X_{-m}\right)$ both converge to $c$. Taking $m=0$, we get in particular that $n^{-1} \log F_{R}\left(e, X_{n}\right)$ and $n^{-1} \log F_{R}\left(X_{n}, e\right)$ almost surely converge to $c$.

We will now prove that $c$ is strictly negative. There exists a real number $h \geqslant$ 0 (the entropy of the $\nu$-random walk) such that the random walk at time $n$ is essentially supported by $e^{h n}$ points, with a probability $e^{-h n}$ to reach each of those points. More precisely (see for instance [Fur02, Theorem 2.28]), for any $\varepsilon>0$, if $n$ is large enough, there exists a subset $E_{n}$ of $\Omega$, with probability at least $3 / 4$, such that for any $\omega \in E_{n}$ one has

$$
\mathbb{P}\left\{\omega^{\prime}: X_{n}\left(\omega^{\prime}\right)=X_{n}(\omega)\right\} \in\left[e^{(-h-\varepsilon) n}, e^{(-h+\varepsilon) n}\right]
$$

Since $\log F_{R}\left(e, X_{n}\right)$ and $\log F_{R}\left(X_{n}, e\right)$ almost surely converge to $c$, we can also assume (shrinking $E_{n}$ a little bit) that for any $\omega \in E_{n}$ one has $F_{R}\left(e, X_{n}\right) \geqslant e^{(c-\varepsilon) n}$ and $F_{R}\left(X_{n}, e\right) \geqslant e^{(c-\varepsilon) n}$. Let $\tilde{E}_{n} \subset \Gamma$ be the set of points $X_{n}(\omega)$ for $\omega \in E_{n}$. It has cardinality at least $C e^{(h-\varepsilon) n}$, it is contained in $B(e, n)$ (since the steps of the $\nu$-random walk have length at most 1 by definition), and for any $x \in \tilde{E}_{n}$ one has $F_{R}(e, x) F_{R}(x, e) \geqslant e^{2(c-\varepsilon) n}$. Therefore,

$$
\begin{aligned}
\sum_{x \in B(e, n)} F_{R}(e, x) F_{R}(x, e) & \geqslant \sum_{x \in \tilde{E}_{n}} F_{R}(e, x) F_{R}(x, e) \geqslant \operatorname{Card} \tilde{E}_{n} \cdot e^{2(c-\varepsilon) n} \\
& \geqslant C e^{(h-\varepsilon) n} e^{2(c-\varepsilon) n}
\end{aligned}
$$

By Lemma 2.5, the sum $\sum_{x \in \mathbb{S}_{k}} F_{R}(e, x) F_{R}(x, e)$ is uniformly bounded [as $F_{R}(x, y)=$ $\left.G_{R}(x, y) / G_{R}(e, e)\right]$. Therefore, $\sum_{x \in B(e, n)} F_{R}(e, x) F_{R}(x, e) \leqslant C(n+1)$. We deduce that $(h-\varepsilon)+2(c-\varepsilon)$ is nonpositive. Finally, letting $\varepsilon$ tend to 0 , we get $c \leqslant-h / 2$. Since entropy is nonzero in nonamenable groups (see for instance Fur02, Proposition $2.35]$ ), we get $c<0$ as desired.

Let us now prove the estimates of the lemma. The first item (positive escape rate) is classical and follows from Kingman's theorem for the existence of the escape rate, and from the inequality $h \leqslant v \zeta$ for its positivity (where $\zeta$ is the exponential growth rate of the cardinality of balls); see Fur02, Proposition 2.32]. For the second item, consider a typical trajectory $X_{k}$ of the random walk. Since $\log F_{R}\left(e, X_{k}\right) \sim c k$ with $c<0$, we deduce that for large enough $k$ one has $F_{R}\left(e, X_{k}\right) \leqslant e^{c k / 2}$. Since $G_{R}(x, y)=F_{R}(x, y) G_{R}(e, e)$, the exponential decay of $G_{R}\left(e, X_{k}\right)$ follows. The decay of $G_{R}\left(X_{k}, e\right)$ is handled in the same way. Finally, consider two independent trajectories $X_{k}$ and $Y_{k}$ of the random walk. Define (for $k>0$ ) $X_{-k}=Y_{k}$. By symmetry of $\nu,\left(X_{k}\right)_{k \in \mathbb{Z}}$ is a typical trajectory for the bilateral random walk. Applying Theorem A.5, we deduce that $\log F_{R}\left(X_{-m}, X_{n}\right) \sim(m+n) c$, i.e., $\log F_{R}\left(Y_{m}, X_{n}\right) \sim$ $(m+n) c$ when $m+n$ tends to infinity. This is the desired exponential decay.

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