

# The circular telegraph process

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## The persistent walk on $\mathbb{Z}/N\mathbb{Z}$

$(V_n)_{n \geq 0}$  markov chain on  $\{\pm 1\}$  with transition matrix

$$\begin{pmatrix} \frac{1+\alpha}{2} & \frac{1-\alpha}{2} \\ \frac{1-\alpha}{2} & \frac{1+\alpha}{2} \end{pmatrix}$$

$(X_n)_{n \geq 0}$  second order markov chain on  $\mathbb{Z}/N\mathbb{Z}$  with

$$X_{n+1} = X_n + V_n$$

### Proposition

The maximal rate of convergence to equilibrium is obtain for

$$\alpha_{opt}^{(N)} = \frac{1 - \sin(\pi/N)}{1 + \sin(\pi/N)}$$

$$\mathbb{P}_{opt}(V_{n+1} = -V_n) \simeq \frac{\pi}{N}$$

## The circular telegraph process

If  $N_t \sim \text{Poisson}(a)$ , let  $(X_t, V_t)$  be markov process on  $\mathbb{T} \times \{\pm 1\}$  defined by

$$\begin{aligned}dX_t &= V_t dt \\V_t &= (-1)^{N_t}\end{aligned}$$

or by its generator on  $L^2(\mu)$  ( $\mu$  being the uniform measure on  $\mathbb{T} \times \{\pm 1\}$ )

$$\begin{aligned}L_a f(x, v) &= v \partial_x f(x, v) + a(f(x, -v) - f(x, v)) \\&= \overbrace{Tf(x, v)} + \overbrace{Kf(x, v)}\end{aligned}$$

and its semi-group  $P_t^a$ . We denote  $f_t = P_t f$  and

$$\begin{aligned}F_t &= \|f_t - \mu f\|_{L^2(\mu)}^2 \\ \|P_t - \mu\| &= \sup_{\|f\|=1} \sqrt{F_t}\end{aligned}$$

There's no randomness on the position variable :

$$\begin{aligned} F'_t &= 2 \langle Kf_t, f_t \rangle \\ &= 0 \quad \text{if } f_t \in \text{Ker}K \end{aligned}$$

and no coercivity from the deterministic part of the process. Nevertheless,

### Proposition

Let  $\lambda_a = -\sup\{\text{Re}(\nu), \nu \in \sigma(L) \setminus \{0\}\}$ . Then

$$\lambda_a = a - \sqrt{(a^2 - 1)_+}$$

$$\|P_t - \mu\| = g_a(t)e^{-\lambda_a t}$$

with  $g_a$  an explicit prefactor such that, in particular,

$$\|P_t - \mu\| = 1 - \frac{\min(a, 1)}{3} t^3 + o_{t \rightarrow 0}(t^3)$$

## $L^2$ norm of the semi-group

The  $W_n = \text{span}(x, v \mapsto e^{inx})$  are stable orthogonal plans, whose span space is dense in  $L^2(\mu)$ . Thus

$$\|P_t - \mu\|_{L^2(\mu)} = \sup_{n \geq 0} \|P_t - \mu\|_{W_n}.$$

In dimension 2, everything is computable :

- $\sigma(L|_{W_n}) = \{-(a \pm \sqrt{a^2 - n^2})\}$
- Let  $f_+, f_-$  be respective eigenfunctions. Then

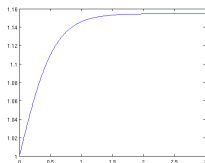
$$|\langle f_+, f_- \rangle| = \min\{|\langle g_+, g_- \rangle| \text{ with } g_{\pm} \text{ eigenfunctions of } M \\ \text{with } \|e^{tM}\| \text{ non-increasing and } \sigma(L|_{W_n}) = \sigma(M)\}$$

- $\|P_t - \mu\|_{W_n} = g_n(t) e^{-t(a - \sqrt{(a^2 - n^2)_+})}$

# Behaviour of the prefactor

$$a > n$$

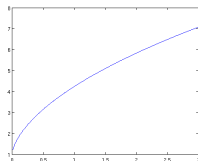
real  
spectrum



$$g_n(t) \rightarrow \frac{\left(\frac{a}{n}\right)^2}{\left(\frac{a}{n}\right)^2 - 1}$$

$$a = n$$

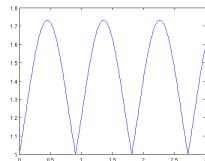
Jordan  
block



$$g_n(t) \sim 2nt$$

$$a < n$$

complex  
spectrum



$$g_n(t) \in \left[ 1, \frac{1 + \frac{a}{n}}{1 - \frac{a}{n}} \right]$$
$$\nu = |\lambda_+ - \lambda_-|$$

## Some details

- The renormalised process  $(X_n^N, V_n^N)$  converges in law to  $(X_t, V_t)$  but

$$\|P_n^N - \mu_N\| \xrightarrow{N \rightarrow \infty} \sup_{n \in \frac{1}{2}\mathbb{N}} \{g_n(t)e^{-\lambda_n t}\} \neq \sup_{n \in \mathbb{N}} \{g_n(t)e^{-\lambda_n t}\} = \|P_t - \mu\|$$

- Moreover

$$\mathcal{L}_{opt}(V_N) \xrightarrow{N \rightarrow \infty} \mathcal{E}\left(\frac{1}{2}\right) \neq \mathcal{E}(1) = \mathcal{L}_{opt}(V_t)$$

- Mostly  $\|P_t - \mu\| = \|P_t - \mu\|_{W_1}$ , but not necessarily if  $a < 1$ .
- If  $\lambda$  is the exponential rate of convergence on  $W_n$  then
  - ▶ if  $a \geq n$ ,  $F_t' + \lambda F_t$  is always non-negative.
  - ▶ if  $a < n$ ,  $F_t' + \lambda F_t = A \cos(\nu_n t + \phi)$ .



### 3-order differential (in)equation

Now  $f \in L^2(\mu)$  is fixed with  $\mu f = 0$ , and  $F_t = \|f_t\|^2$ . Then

$$\partial_t F_t = \langle (2K)f_t, f_t \rangle$$

$$\partial_t^2 F_t = \langle \left( (2K)^2 + 2[K, T] \right) f_t, f_t \rangle$$

$$\partial_t^3 F_t = \langle \left( (2K)^3 + 6[K^2, T] + 2[[K, T], T] \right) f_t, f_t \rangle .$$

In particular for  $f_t \in \text{Ker}K$  then  $F_t' = 0$  but

$$\partial_t^3 F_t = -4a \|\partial_x f_t\|^2 \leq -4a F_t$$

In fact, for  $f \in W_n$ ,

$$\left( (\partial_t + 2a)^2 + 4(n^2 - a^2) \right) (\partial_t + 2a) F_t = 0$$

Nothing like  $F_t''' + AF_t'' + BF_t' + CF_t \leq 0$  can be true for all  $f$ .

Let  $h_t = e^{2at} F_t$  then either

- $h'_t \leq 0 \quad \longrightarrow$  exponential decay.
- Else  $h_t''' + 4(1 - a^2)h'_t \leq h_t''' + 4(n^2 - a^2)h'_t = 0 \quad \longrightarrow$  uniform w.r.t.  $n$ .

The inequality can be written

$$Q_3(\partial t)F_t + Q_1(\partial t)G_t \leq 0$$

with  $Q_i$  polynomial of degree  $i$  and  $G_t = \|\partial_x f_t\|^2$ .

### Proposition

Suppose now  $a$  depends on the position  $x$ , with  $a(x) \geq a_* > 0$ . Then one can still find  $Q_1$  and  $Q_3$  such that

$$Q_3(\partial t)F_t + Q_1(\partial t)G_t \leq 0$$

for all  $f \in L^2(\mu)$ .

## From the inequality to exponential decay

$G_t$  brings the coercivity missing in  $F'_t$  :

$$\begin{aligned} G_t - F'_t &= \langle (\partial_x^* \partial_x - 2K) f_t, f_t \rangle \\ &:= k_t F_t \end{aligned}$$

with  $k_t \geq k_* > 0$ . Now we have

$$\tilde{Q}_3(\partial t) F_t + \tilde{Q}_1(\partial t) (k_t F_t) \leq 0$$

And, if  $\tilde{Q}_1 | \tilde{Q}_3$  (which can be done),

$$(\partial_t + \lambda) \left( (\partial_t + \eta)^2 + \nu_t \right) F_t \leq 0$$

With some  $\nu_t \geq \nu_*$ . By the Gronwall lemma

$$\left( (\partial_t + \eta)^2 + \nu_* \right) F_t \leq \left( (\partial_t + \eta)^2 + \nu_t \right) F_t \leq M_0 e^{-\lambda t}$$

Let  $\phi$  be the solution of

$$\begin{cases} [(\partial_t + \eta)^2 + \nu_*] \phi = M_0 e^{-\lambda t} \\ \phi(0) = F(0), \phi'(0) = F'(0). \end{cases}$$

so that  $h_t = e^{\eta t} (F_t - \phi(t))$  satisfy  $h_t'' + \nu_* h_t \leq 0$ . Then :

- Either  $\nu_* \leq 0$  and applying twice the Gronwall lemma gives  $h_t \leq 0$ .
- Either  $\nu_* > 0$ ; then while  $h_t' > 0$

$$\frac{d}{dt} \left( (h_t')^2 + \nu_* h_t^2 \right) \leq 0$$

so that  $h_t \leq -\min_{s \leq t} h_s$ . Moreover, by the Sturm theorem,  
 $|\{h_t > 0\}| \leq \frac{\pi}{\nu_*} : F$  and  $\phi$  are interlaced.

## To sum up :

The study of  $K$ ,  $T$ , etc. gives  $\lambda > 0$ ,  $\eta$  and a function  $\nu_t \geq \nu_*$  with  $\operatorname{Re}(\eta - \sqrt{\nu_*}) > 0$  s.t.

$$(\partial_t + \lambda) \left( (\partial_t + \eta)^2 + \nu_t \right) F_t \leq 0$$

and this implies

- if  $\nu_* \leq 0$ ,  $F_t \leq \phi(t)$
- if  $\nu_* > 0$ ,  $F_t \leq \phi(t) + e^{-\eta t} \sup_{s \leq t} e^{\eta s} (\phi_s - F_s)$ ; furthermore the length of a time interval where  $F > \phi$  is less than  $\frac{\pi}{\nu_*}$

with in both cases  $\phi$  solution of

$$(\partial_t + \lambda) \left[ (\partial_t + \eta)^2 + \nu_* \right] \phi(t) = 0,$$

$\phi(0) = F_0$ ,  $\phi'(0) = F'_0$  and  $\phi''(0)$  explicit.

## Some remarks :

- In particular,

$$\lim_{t \rightarrow \infty} \frac{1}{t} \ln F_t \leq -\min(\lambda, \operatorname{Re}(\eta - \sqrt{-\nu_*})) < 0$$

- As  $\phi_0''$  depends on  $F_0''$  and  $G_0$ , this does not give directly a bound for  $\|P_t - \mu\|$ .
- $F_t$  is the real  $L^2$  norm of  $f_t$ , but  $G_t$  is needed to close the inequality.
- The method works with  $V_t$  an Ornstein-Uhlenbeck process, in other words with

$$L = v\partial_x - \partial_v + \partial_v^2$$

or with the kinetic Fokker-Planck dynamics (provided  $U''$  is bounded)

$$L = v\partial_x + U'(x)\partial_v - v\partial_v + \partial_v^2$$

## Application to other models ?

- If  $a$  depends on  $x$  and  $v$  the invariant measure is  $e^{-U(x)} dx \otimes U_{\{\pm 1\}}$  where

$$a(x, 1) - a(x, -1) = U'(x).$$

Can we still apply the same method ?

- Same question for the next-to-be presented ergodic variant of the telegraph process with generator

$$Lf = v\partial_x f + \left( a\mathbf{1}_{v=\text{sign}(x)} + b\mathbf{1}_{v=-\text{sign}(x)} \right) (f(x, -v) - f(x, v))$$

with  $b > a$ . Problem : what is  $\partial_x \left( a\mathbf{1}_{v=\text{sign}(x)} + b\mathbf{1}_{v=-\text{sign}(x)} \right)$  ?

- For some switching kinetic processes we also have

$$f_t \in \text{Ker}K \Rightarrow F_t''' \leq -CF_t.$$

Thank you for your attention.

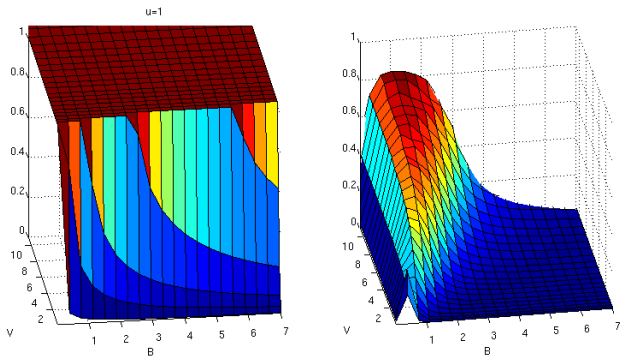


FIGURE: Theoretical and numerical rate of relaxation for the kinetic Fokker-Planck