

# The Galois groupoid of Picard-Painlevé VI equation

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The sixth Painlevé equation for special values of classical parameters ( $\alpha = \beta = \gamma = 0, \delta = 1/2$ ) was discovered by E.Picard in [10] as an exemple of order two non-linear equation without movable singularities. The usual form of this equation is called in this paper  $PP_6$ :

$$y'' = \frac{1}{2} \left( \frac{1}{y} + \frac{1}{y-1} + \frac{1}{y-x} \right) (y')^2 + \left( \frac{1}{x-y} + \frac{1}{1-x} - \frac{1}{x} \right) y' + \frac{y(y-1)}{2x(x-1)(y-x)}.$$

Among all the  $P_6$  equations this one has the property to be solved by a formula :

$$(1) \quad y(x) = \wp(a\omega_1(x) + b\omega_2(x); \omega_1(x), \omega_2(x))$$

with  $a$  and  $b$  two constants,  $\omega_{1,2}$  a basis of periods of  $t^2 = y(y-1)(y-x)$  and  $\wp(\cdot; \omega_1, \omega_2)$  the corresponding Weierstass function. For rational  $a$  and  $b$ , the solution is algebraic but for other values the solution is not even a classical function in the sense of H.Umemura [12] despite the formula to express it as it is proved by H. Watanabe in [15]. For a complete study of this equation, see the article of M.Mazzocco [8].

This is a common belief that this kind of property must be explained by a non-linear Galois theory. Two essentially equivalent very general differential Galoisian theories have been proposed in the last ten years by H.Umemura [13, 14] and B.Malgrange [6]. Because of its geometric flavor, we will focus on Malgrange's Galois groupoid to explain the existence of a formula to solve  $PP_6$ . The computation of the Galois groupoid of  $PP_6$  can be reduced to computation already done by P.Painlevé in ([9] pp 501–517). Painlevé remarked that  $PP_6$  is irreducible in his sense (which is very close to Nishioka-Umemura definition [11]) but it admits a system of first integrals in a Picard-Vessiot extension of the partial differential field  $(\mathbb{C}(x, y, y'); \partial/\partial x, \partial/\partial y, \partial/\partial y')$ . The aim of this article is to integrate this remark from Painlevé in the framework of Malgrange's Galois groupoid. Solutions of Painlevé equations describe isomonodromy deformations of rank two Fuchsian systems with four singularities on  $\mathbb{CP}^1$ . Algebraic properties of the monodromy data are related to the transcendence nature of the corresponding Painlevé solution [4]. In the Picard-Painlevé case, the monodromy data are not special but the space of monodromy data and the non-linear monodromy are special [8]

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# 1 The Galois groupoid of a vector field in $\mathbb{C}^3$

Let  $X$  be a rational vector field in  $\mathbb{C}^3$ . In general it is not complete and its flows are only defined on open set small enough. All the dynamic of this vector field is contained in the pseudogroup of transformations of  $\mathbb{C}^3$  generated by these local flows. By keeping only the germs of diffeomorphisms from this pseudogroup one gets a groupoid,  $TanX$ , acting on  $\mathbb{C}^3$ .

The Galois groupoid of  $X$  is the Zariski closure of  $TanX$  for a (nearly) obvious embedding of  $TanX$  in an infinite dimensional algebraic variety.

This variety is the space  $J^*$  of formal diffeomorphisms of  $\mathbb{C}^3$ , *i.e.* the set of formal invertible maps  $\varphi : \widehat{\mathbb{C}^3, a} \rightarrow \widehat{\mathbb{C}^3, b}$  and the embedding is the Taylor expansion of elements of  $TanX$ . The space  $J^*$  can be presented as the projective limit of the spaces  $J_q^*$  of order  $q$  jets of diffeomorphisms. These ones are isomorphic to

$$\mathbb{C}^3 \times \mathbb{C}^3 \times \text{Gl}(\mathbb{C}^3) \times_{\substack{\alpha \in \mathbb{N}^3 \\ |\alpha| \leq q}} \mathbb{C}^{3|\alpha|}$$

and their coordinates rings

$$O(J_q^*) = \mathbb{C} \left[ x, y, y', \bar{x}^\alpha, \bar{y}^\alpha, \bar{y}'^\alpha, \frac{1}{\det(\bar{x}^{\epsilon_i}, \bar{y}^{\epsilon_i}, \bar{y}'^{\epsilon_i})} \mid \alpha \in \mathbb{N}^3; |\alpha| \leq q \right]$$

are the rings of partial differential equations of order  $q$  in three functions of three arguments with non vanishing jacobian. Futhermore these varieties get natural groupoid structures given by the computation rules for the Taylor expansion of the composition of formal diffeomorphisms. An algebraic subgroupoid of  $J_q^*$  is an algebraic subvariety whose ideal satisfies some stability conditions under inversion and composition.

The space  $J^*$  is  $\varprojlim J_q^*$  and its ring  $O(J^*) = \varinjlim O(J_q^*)$  is the commutative differential ring of non-linear partial differential equations on germs of diffeomorphisms. The derivations are given by the natural actions of  $\partial/\partial x$ ,  $\partial/\partial y$ ,  $\partial/\partial y'$  on partial differential equations.

**Definition 1.1 (Malgrange [6])** *Let  $G$  be a subvariety of  $J^*$  described by a differential ideal. It is an algebraic  $\mathcal{D}$ -groupoid on  $\mathbb{C}^3$  if there is a subvariety  $Z \subset \mathbb{C}^3$  such that the projections of  $G$  on the finite order jet spaces defined algebraic subgroupoids on  $\mathbb{C}^3 - Z$ .*

**Definition 1.2 (Malgrange [6])** *The Galois groupoid of  $X$  is the smallest algebraic  $\mathcal{D}$ -groupoid on  $\mathbb{C}^3$  containing  $TanX$ .*

Roughly speaking, the Galois groupoid of  $X$  is the set of all the germs of diffeomorphisms of  $\mathbb{C}^3$  solutions of all the PDE's vanishing on  $TanX$ . Because  $L_X X = 0$ , the Galois groupoid is a subgroupoid of the groupoid of transformations preserving  $X$  *i.e.* germs  $\varphi$  such that  $\varphi_* X = X$ . When the vector field is divergence free, one can add the equations given by the coordinates of  $\varphi^*(dx \wedge dy \wedge dy') = dx \wedge dy \wedge dy'$ .

Near a regular point of  $X$  one can choose a flowbox given by two transversal coordinates  $t_1, t_2$  and a tangent one  $z$ . A germ of diffeomorphism preserving  $X$  can be written

$$\varphi : \begin{cases} \bar{t}_1(t_1, t_2) \\ \bar{t}_2(t_1, t_2) \\ \bar{z}(z, t_1, t_2). \end{cases}$$

By looking at the PDE's of the Galois groupoid in these coordinates, two type of PDE can be distinguished. The PDE's vanishing on the germs such that  $\bar{z} = z$  are called tangential equations. The others are the transversal ones. For example, the transversal equations of the set of PDE :  $L_X X = 0$  are given by  $L_X X \wedge X = 0$ .

Using Lie-Cartan local classification of pseudogroups acting on  $\mathbb{C}^2$  established in [1], one has the following proposition

**Proposition 1.3 ([3])** *If  $X$  is divergence free and  $\gamma$  is the closed 2-form vanishing on  $X$ , i.e.  $\iota_X \gamma = 0$ , then one of the following situations occurs:*

- *Gal( $X$ ) is imprimitive in codimension one: there exists an algebraic 1-form  $\theta$  such that  $\theta \wedge d\theta = 0$  and  $\theta(X) = 0$ ,*
- *Gal( $X$ ) is transversally affine: there exists two algebraic independent 1-forms  $\theta_{1,2}$  vanishing on  $X$  and a traceless matrix of 1-form  $(\theta_{1,2}^{1,2})$  such that  $d\theta_i = \sum_j \theta_i^j \wedge \theta_j$  and  $d\theta_i^j = \sum_k \theta_i^k \wedge \theta_k^j$ ,*
- *the only transversal equations of Gal( $X$ ) are those of invariance of  $\gamma$ .*

The result presented in this paper is the following.

**Theorem 1.4 (Painlevé [9])** *The Galois groupoid of  $PP_6$  is transversally affine.*

Transversally affine vector fields admit very special first integrals given by the affine structure out of a codimension one subvariety. These are described in the next section. The vector field over  $\mathbb{C}(x, y, y')$  of  $PP_6$  is not divergence free for the usual 'volume' form  $dx \wedge dy \wedge dy'$  but for the canonical one. All the Painlevé equation can be express as time dependent Hamiltonians [7]. For  $PP_6$ , the Hamiltonian is :

$$\begin{cases} \frac{dq}{dx} = \frac{\partial K}{\partial p} \\ \frac{dp}{dx} = -\frac{\partial K}{\partial q} \\ K = \frac{1}{x(x-1)}[q(q-1)(q-x)p^2 - pq(q-1) + \frac{1}{4}(q-x)] \\ q = y \text{ and } p = \left(\frac{x(x-1)}{2y(y-1)(y-x)}\right) y' + \frac{1}{2(y-x)}. \end{cases}$$

Let  $\theta_1$  and  $\theta_2$  be the forms  $dq - (\partial K / \partial p)dx$  and  $dp + (\partial K / \partial q)dx$ ,  $\theta_1 \wedge \theta_2 = dq \wedge dp + dH \wedge dx$  is a closed 2-form vanishing on  $PP_6$ .

## 2 Classical first integrals and Galois groupoid

The classical functions over  $\mathbb{C}$  were introduced in [12] by H.Umemura. A function of one variable is said to be classical if one can find it in an ordinary differential field extension of the rational functions field  $\mathbb{C}(x)$  built by successive strongly normal extensions [5] or algebraic extensions. In general two successive strongly normal extensions fail to be a strongly normal one. Umemura's definition of classical functions of  $n$  variables is the following.

**Definition 2.1 (Umemura [12])** *Let  $\mathbb{C}(x_1, \dots, x_n)$  be the partial differential field of rational functions of  $n$  arguments with derivations  $\partial_{x_1} \dots \partial_{x_n}$  and  $\mathcal{K}$  be a differential extension of  $\mathbb{C}(x_1, \dots, x_n)$ . It is said to be classical if one can find a tower of differential extensions*

$$\mathbb{C}(x_1, \dots, x_n) = \mathcal{K}_0 \subset \mathcal{K}_1 \dots \subset \mathcal{K}_p = \mathcal{K}$$

*such that  $K_i \subset K_{i+1}$  is one of the following*

- algebraic,
- Picard-Vessiot:  $K_{i+1} = K_i(H_p^q)$  where the  $H$ 's are entries of a fundamental solution of linear equations  $\partial_{x_\ell} H_p^q = A_{\ell,p}^m H_m^q$  with  $A$ 's matrices with entries in  $K_i$ ,
- 'Abelian': There exists an Abelian function  $A$  of  $m$  arguments and  $f_1, \dots, f_m \in K_i$ ,  $K_{i+1} = K_i \langle A(f_1, \dots, f_m) \rangle$ .

Most of the studies of Painlevé equations focus on their classical solutions ( $n = 1$  in the definition) but it fails to explain the 'solvability' of the Picard-Painlevé equations. Nevertheless it is well known that this equation has classical first integrals ( $n = 3$  in the definition) ([9]).

One of the basics properties of this type of extension is to have a finite transcendence degree. From the lemmas 4.4.5 and 4.4.6 of [2] and proposition 1.3 of this article, one deduces the particular form of the Galois groupoid of a vector field with classical first integrals.

**Proposition 2.2** *Let  $X$  be a divergence free vector field on  $\mathbb{C}^3$  with two independent classical first integrals. Then its Galois groupoid is transversally affine or imprimitive in codimension one.*

In fact all the imprimitive in codimension one cases cannot occur. It is possible to give a more precise statement for this proposition but it is not needed in this paper.

Conversely any transversally affine vector field has classical first integrals. They are built by solving the following linear system:

$$dL_i^j = \sum_k L_i^k \theta_k^j \quad dH_i = \sum_j L_i^j \theta_j.$$

These functions are more than classical, they are in a Picard-Vessot extension of the field of rational functions of  $\mathbb{C}^3$ .

### 3 The first integrals of $PP_6$

In this section, computation of special first integrals of  $PP_6$  is done to prove theorem 1.4. This computation follows P.Painlevé [9].

The  $PP_6$  equations was discovered by E.Picard as the pull-back of a linear order two equation by a transcendental function. For this reason it is solvable by the formula 1. This formula and the pull-back are given by a integral with  $x$  as parameter

$$\int_0^{y(x)} \frac{d\xi}{\sqrt{\xi(\xi-1)(\xi-x)}} = a\omega_1(x) + b\omega_2(x),$$

$\omega_{1,2}$  form a basis of periods of  $t^2 = y(y-1)(y-x)$  i.e. the right hand side of the equality is the general solution of the Picard-Fuchs ( $PF$ ) equation  $4x(x-1)w'' - 4(2x-1)w' - w = 0$ . Let  $X_{PP}$  and  $X_{PF}$  be the vector fields corresponding to the  $PP_6$  and  $PF$  equations on their phase spaces.  $X_{PP}$  is the pull-back of  $X_{PF}$  by the following map:

$$\begin{cases} x = x \\ w = \int_0^y \frac{d\xi}{\sqrt{\xi(\xi-1)(\xi-x)}} \\ w' = \frac{y'}{\sqrt{y(y-1)(y-x)}} + \int_0^y \frac{d\xi}{2(\xi-x)\sqrt{\xi(\xi-1)(\xi-x)}}. \end{cases}$$

To get two first integrals for  $PP_6$ , one pull-backs two first integrals of  $PF$ . Because it is a linear equation, it has first integrals linear on the ‘fibers’ *i.e.*  $H = \alpha(x)w + \beta(x)w'$  where  $(\alpha, \beta)$  is a solution of the  $PF$ 's adjoint equation :

$$\begin{cases} \alpha' = \beta \frac{-1}{4x(x-1)} \\ \beta' = \beta \frac{1-2x}{x(x-1)} - \alpha. \end{cases}$$

Let  $H_{\alpha, \beta}$  be the function

$$\frac{y'\beta}{\sqrt{y(y-1)(y-x)}} + \int_0^y \left( \alpha + \frac{\beta}{2(\xi-x)} \right) \frac{d\xi}{\sqrt{\xi(\xi-1)(\xi-x)}}.$$

It turns out that this function is a primitive of a closed 1-form with coefficients in the field  $\mathbb{C}(x, \alpha, \beta, y, y' \sqrt{y(y-1)(y-x)})$  :

$$H_{\alpha, \beta} = \frac{y'\beta}{\sqrt{y(y-1)(y-x)}} + \int \left( \alpha + \frac{\beta}{2(y-x)} \right) \frac{dy}{\sqrt{y(y-1)(y-x)}} - \frac{\beta y(y-1)}{2x(1-x)(y-x)} \frac{dx}{\sqrt{y(y-1)(y-x)}}.$$

The transversally affine structure is given by a sequence of 1-forms. It is derived from two first integrals  $H_1 = H_{\alpha_1, \beta_1}$  and  $H_2 = H_{\alpha_2, \beta_2}$  with  $\alpha_{1,2}$  and  $\beta_{1,2}$  the entries of a fundamental solution  $F$  of the  $PF$ 's adjoint equation. The derivatives give

$$dH_i = L_i^1 \theta_1 + L_i^2 \theta_2$$

where the  $\theta$ 's are the 1-forms given by the Hamiltonian form of  $PP_6$ . By construction,  $H_1$  and  $H_2$  are linear in  $\alpha_{1,2}$  and  $\beta_{1,2}$  and the matrix  $L$  equals  $(\sqrt{y(y-1)(y-x)})^{-1}MF$  for a matrix  $M$  with entries in  $\mathbb{C}(x, y, y')$ .

The matrix  $F$  satisfies a linear equation  $dF = F\Omega$ ,  $\Omega$  a matrix of 1-forms with coefficients in  $\mathbb{C}(x, y, y')$ . This implies that  $dL = L\underline{\Omega}$  with  $\underline{\Omega} = -(1/2)(1/y + 1/(y-1) + 1/(y-x))dy \text{Id} + dMM^{-1} + M\underline{\Omega}M^{-1}$ ,  $\text{Id}$  is the identity matrix. Because  $d\Omega = 0$ , one gets  $d\underline{\Omega} = \underline{\Omega} \wedge \underline{\Omega}$ . The 1-forms  $\theta_{1,2}$  and the coefficients  $\theta_{1,2}^{1,2}$  of  $\underline{\Omega}$  give the transversally affine structure.

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