An infinite-dimensional Cartan development

> Pierre Perruchaud University of Notre Dame

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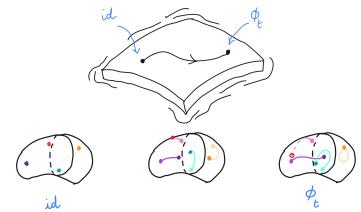
- The Lagrangian approach to fluid mechanics An illustrated guide to manifolds of maps
- II Stochastic processes on manifolds Construction of the Brownian motion The Cartan development
- III Infinite-dimensional geometry A roadmap to rigour in Euler geometry
- IV The infinite-dimensional Cartan development *Developing fluids?*

Part I

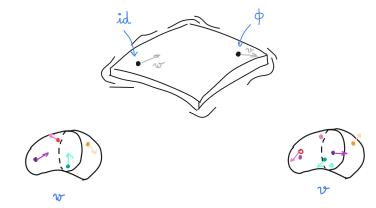
The Lagrangian approach to fluid mechanics

Fix a (finite dimensional) closed Riemannian manifold M of dimension d, and call μ its measure. Consider \mathcal{D} some space of maps $M \to M$, and $\mathcal{D}_{\mu} \subset \mathcal{D}$ some space of volume-preserving diffeomorphisms.

We think of points in \mathcal{D} as fluid configurations, of curves in \mathcal{D} as fluid motions. The objective is to define Brownian fluids (and more general stochastic fluids).

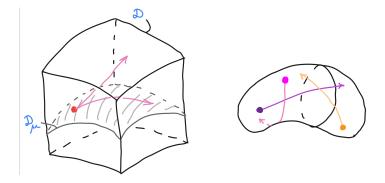


- $\blacktriangleright \text{ Point in } \mathcal{D} \leftrightarrow \mathsf{map}$
- $\blacktriangleright \ \mathsf{Curve in} \ \mathcal{D} \leftrightarrow \mathsf{flow}$
- Tangent vector at id $\in \mathcal{D} \leftrightarrow$ vector field
- Tangent vector at $\phi \in \mathcal{D} \leftrightarrow$ vector field rooted at ϕ .



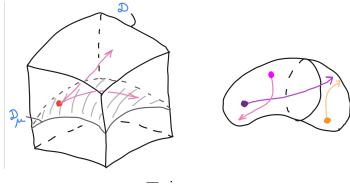
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For *some* Riemannian metric on \mathcal{D} , the geodesics follow the inviscid Burgers equation: every fluid particle goes straight ahead (we may have immediate collision).



 $\overline{
abla}_{\dot{\phi}}\dot{\phi}=0, \qquad
abla_{\dot{\phi}(x)}(\dot{\phi}(x))=0 \qquad ext{ for all } x\in M.$

For the same Riemannian metric restricted to \mathcal{D}_{μ} , the geodesics follow the incompressible inviscid Euler equations: the fluid particles cannot collide (they have an associated volume) but try to go straight ahead as much as possible.



$$\overline{\nabla}_{\dot{\phi}}\phi = -\nabla p.$$

Objective: define a common framework for classical fluids, Brownian fluids, stochastic perturbations of classical fluids, and possibly more. Solution: the Cartan development Part II

Stochastic processes on manifolds

Brownian motion is the process with generator $\frac{1}{2}\Delta_M$.

Locally, up to order 1, Δ_M is a sum of squares: if $(X_1(x), \ldots, X_d(x))$ is an orthonormal basis at each point,

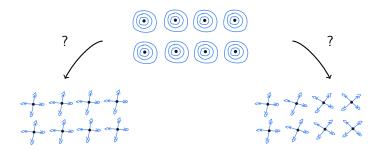
$$\Delta_M f = X_1 \cdot (X_1 \cdot f) + \cdots + X_d \cdot (X_d \cdot f) + X_0 \cdot f,$$

for some vector field X_0 . It means the Brownian motion x satisfies

$$dx_t = X_1(x_t) \circ dW_t^1 + \cdots + X_d(x_t) \circ dW_t^d + X_0(x_t)dt.$$

Somewhat unsatisfactory for a few reasons.

- Arbitrary choice for the basis; we break the symmetry.
- X_0 depends on the choice of basis.
- > Patching the local SDEs requires a higher-dimensional noise.



Idea of Eells, Elworthy and Malliavin: enlarge the space to the orthonormal frame bundle *OM*.

u = (x, u) is a point x on the manifold together with some orthonormal basis $(u(\epsilon_i))_i$ at $x (u : \mathbb{R}^d \to T_x M$ is an isometry). Informal equation:

$$dx_t = u_t(\epsilon_1) \circ dW_t^1 + \dots + u_t(\epsilon_d) \circ dW_t^d, \quad \nabla_{dx_t} u_t = 0.$$

= $u_t(\circ dW_t)$

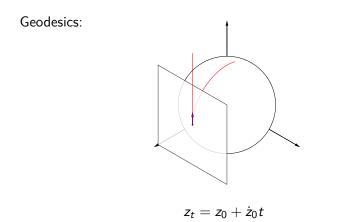
Actual equation:

$$d(x_t, u_t) = \underbrace{H(x_t, u_t)}_{\text{horizontal lift}} \circ dW_t.$$

II. Stochastic processes on manifolds

In fact, we can do much more than just Brownian motion: Cartan development.

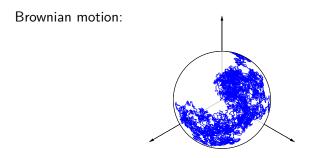
 $\mathrm{d} u_t = H(u_t) \mathrm{d} z_t$



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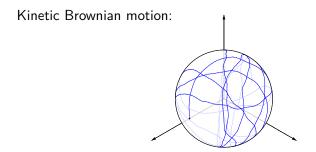


$$z_t = W_t$$

II. Stochastic processes on manifolds

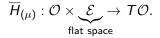
In fact, we can do much more than just Brownian motion: Cartan development.

 $\mathrm{d} u_t = H(u_t) \mathrm{d} z_t$



 $\dot{z} =$ Brownian motion on the sphere

The plan: Define infinite-dimensional versions $\mathcal{O}_{(\mu)}$ of the orthonormal frame bundle, with smooth horizontal lifts



Then, for any flat motion z of your choosing, we can try to solve the smooth controlled equation

$$\mathsf{d} u_t = \overline{H}_{(\mu)}(u_t) \circ \mathsf{d} z_t.$$

Part III

The infinite-dimensional orthonormal bundle

1. Topology

Recall *M* is a (finite-dimensional) closed Riemannian manifold of dimension *d*, and fix s > d/2.

There is a way to define sections of regularity H^s for every given bundle F over M, and they comes with a natural manifold structure. For instance:

- the H^s vector fields over M are the H^s sections of TM;
- the H^s maps $\phi: M \to M$ are the H^s sections of $M \times M$;
- ► the H^s "unrooted" vector fields are the H^s sections of TM × M.

We write $H^{s}(F)$ for the manifold of H^{s} -regular sections of F.

1. Topology

Lemma (Omega lemma)

Pointwise operations are smooth.

Example:

For M an oriented surface, the operation R that sends

- ▶ a collection of vectors v(x) (a vector field $v \in H^s(TM)$)
- ▶ and a collection of angles $\theta(x)$ (a function $\theta \in H^{s}(\mathbb{R} \times M)$)
- ► to the collection of vectors R_{θ(x)}(v(x)) rotated clockwise (a vector field R(θ, v) ∈ H^s(TM))

is pointwise. By the Omega lemma, it is smooth.

1. Topology

Recall s > d/2. Define

$$\mathcal{D} := \{ \phi : M \xrightarrow{H^s} M \},\$$
$$\mathcal{D}_{\mu} := \{ \phi : M \xrightarrow{H^{s+1}} M \text{ volume-preserving diffeomorphism} \}.$$

(Not obvious but true: \mathcal{D}_{μ} is a submanifold of \mathcal{D}^{s+1} .)

The tangent space at $\phi \in \mathcal{D}$ is the space of H^s vector fields rooted at ϕ , so the total tangent space is $H^s(TM)$. The tangent space at $\phi \in \mathcal{D}_{\mu}$ is the space of divergence-free H^{s+1} vector fields rooted at ϕ .

2. Geometry

Metric = total kinetic energy. For $v \in T_{\phi}\mathcal{D}$ (vector field rooted at ϕ),

$$\frac{1}{2}|v|_{L^2}^2 = \frac{1}{2}\int_M |v(x)|^2 dx.$$

Weak Riemannian structure: need not admit geodesics, a connection, etc. However, in $\mathcal{D}_{(\mu)}$, it does! In \mathcal{D} , everything is easy because it works pointwise:

- $t \mapsto \phi_t$ geodesic means $t \mapsto \phi_t(x)$ geodesic for all x;
- for $t \mapsto v_t$ a vector field along $t \mapsto \phi_t$,

$$(\overline{\nabla}_{\dot{\phi}}v)(x) = \nabla_{\dot{\phi}(x)}v(x).$$

2. Geometry

In \mathcal{D}_{μ} , everything is also nice because the orthogonal projections $P_{\phi}: T_{\phi}\mathcal{D} \to T_{\phi}\mathcal{D}$ on the space of divergence-free vector fields rooted at $\phi \in \mathcal{D}_{\mu}$ patch into a smooth $P: T\mathcal{D}_{|\mathcal{D}_{\mu}} \to T\mathcal{D}$.

Important difference: the geodesics are only defined locally.

Part IV

The infinite-dimensional Cartan development

Pairs (ϕ, u) , where $\phi : T_{id}\mathcal{D}_{(\mu)} \to T_{\phi}\mathcal{D}_{(\mu)}$ is an isometry.

- Pointwise (easy) isometries: choose an isometry e(x) : T_xM → T_{φ(x)}M for each x, and compose at the target. We write u = [e].
- Non-local (subtle) isometries: anything else.
 For instance, move the energy across the Fourier modes.

There is an easy way to construct the pointwise orthonormal frame bundle and the pointwise Cartan development. It turns out that they are enough to construct the Cartan development in \mathcal{D} , but not in \mathcal{D}_{μ} .

IV. Infinite-dimensional Cartan development

Because of the non-local character, it is not obvious that the set of pairs (ϕ, u) should admit a manifold structure.

Trick: pull back the non-local issues to a fixed Banach space. For any isometry $u: T_{id}\mathcal{D} \to T_{\phi}\mathcal{D}$ and collection e of isometries $e(x): T_xM \to T_{\phi(x)}M$,

$$u = [e] \circ \underbrace{([e]^{-1} \circ u)}_{=: u^0}.$$

The space of triples (ϕ, e, u^0) is a product

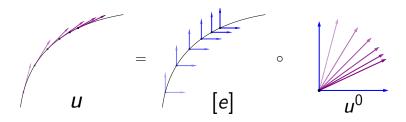
 $H^{s}(F) \times \operatorname{Isom}(T_{\operatorname{id}}\mathcal{D})$

for *F* well-chosen, and Isom($T_{id}D$) inside to a fixed Banach space.

This is big enough to be used as a $\mathcal{O}_{(\mu)}$, and the horizontal lift H over \mathcal{O} is easily defined.

What about \mathcal{D}_{μ} ? Rough idea: $u = [e] \circ u^0$ with

- e parallel transported with no constraint,
- u^0 continuously adjusting to preserve the volume.



Actual construction technical but not clever; same as parallel transport in finite-dimensional submanifolds.

Thank you for your attention.

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