

# Postulates of quantum mechanics viewed as a non-commutative extension of probability theory with a dynamical law

Dimitri Petritis

Institut de recherche mathématique de Rennes  
Université de Rennes 1 et CNRS (UMR 6625)

Santiago, November 2013

# Phase space and states

## Postulate

*Phase space quantum system: complex separable Hilbert space  $\mathbb{H}$ .  
Pure states are unit vectors<sup>a</sup> of  $\mathbb{H}$ .*

<sup>a</sup>Strictly speaking, equivalence classes of unit vectors differing by a global phase, called **rays**, so that  $\mathbf{S}_p \simeq \mathbb{H}_1 / \sim$  and  $h' \sim h \Leftrightarrow h' = \lambda h, \lambda \in \mathbb{C}, |\lambda| = 1$ .

- Hilbert space:  $\mathbb{C}$ -Banach space, with sesquilinear form  $\langle \cdot | \cdot \rangle : \mathbb{H} \times \mathbb{H} \rightarrow \mathbb{C}$ , s.t.
  - $\langle \alpha_1 \phi_1 + \alpha_2 \phi_2 | \psi \rangle = \overline{\alpha_1} \langle \phi_1 | \psi \rangle + \overline{\alpha_2} \langle \phi_2 | \psi \rangle$ ,
  - $\langle \phi | \psi \rangle = \overline{\langle \psi | \phi \rangle}$ ,
  - $\|\psi\|^2 = \langle \psi | \psi \rangle$ .
- Ray:  $\psi \in \mathbb{H}, \|\psi\| = 1$ .
  - But  $\mathbb{H}$  separable  $\Rightarrow \|\psi\|^2 = \sum_n |\psi_n|^2 = 1$ .
  - Hence  $(|\psi_n|^2)_n$  **probability vector** on set indexing the basis of  $\mathbb{H}$ .

# Time evolution

## Postulate

*Time evolution of **isolated** quantum system described by a **unitary** operator  $U$  acting on  $\mathbb{H}$ . Conversely, any unitary operator acting on  $\mathbb{H}$  corresponds to possible invariance<sup>a</sup> of the system.*

---

<sup>a</sup>Time evolution of isolated system leaves physical quantity “energy” invariant. Unitary operators are associated with conserved quantities.

- $U$  unitary:  $U^*U = UU^* = I \Rightarrow U^{-1} = U^*$ .
- $\phi$  ray:  $\|U\phi\|^2 = \langle U\phi | U\phi \rangle = \langle \phi | U^*U\phi \rangle = \|\phi\|^2 = 1$ .
- Hence  $U\phi$  ray.

## Sharp and unsharp observables

### Postulate

**General real sharp observable:** self-adjoint  $X$  acting on  $\mathbb{H}$ .  
**Elementary real sharp observables associated with  $X$ :**

$$\mathcal{B}(\mathbb{R}) \ni B \mapsto M(B) \in \mathfrak{P}(\mathbb{H}); M \text{ spectral projector of } X.$$

**Unsharp real observable:** resolution of identity into positive operators that are not necessarily projections.

- $X^* = X \Rightarrow \text{spec } X := \mathbb{X} \subseteq \mathbb{R}$ .
- $M$  spectral measure of  $X$ ;  $\text{supp } M = \text{spec } X$  but can trivially be extended on  $\mathbb{R}$ .
- $X = \int_{\text{spec } X} M(dx)x$ .

# Measurements (innocent-looking but against intuition)

## Postulate

Let  $M$  spectral measure of sharp observable  $X$ . **Measuring** observable  $X$  in pure state  $\psi$  is asking whether  $X$  has spectral values in  $B$  and determining probability  $\pi_M^\psi(B)$  of occurrence, by

$$\mathbf{S}_p \ni \psi \mapsto \pi_M^\psi(B) := \langle \psi | M(B)\psi \rangle \in [0, 1].$$

- $\pi_M^\psi \in \mathcal{M}_1(\mathcal{B}(\mathbb{R}))$  supported by  $\mathbb{X} := \text{spec } X$ .
- $X = \int_{\text{spec } X} M(dx)x$   
 $\Rightarrow \mathbb{E}_\psi(X) = \int_{\mathbb{X}} \pi_M^\psi(dx)x = \int_{\mathbb{X}} \langle \psi | M(dx)\psi \rangle x = \langle \psi | X\psi \rangle$ .
- $M$  in CM indicator; in QM projection operator acting on  $\mathbb{H}$ .
- Quid if 2 observables with spectral measures  $M$  and  $N$ ?

## Composite systems

### Postulate

*System composed from two subsystems — described respectively by  $\mathbb{H}_1$  and  $\mathbb{H}_2$  — is described by  $\mathbb{H}_1 \otimes \mathbb{H}_2$  (where  $\otimes$  denotes tensor product).*

Blackboard 1: Tensor product (a first look).

# A very simple illustrative example

## Phase space, states, and time evolution

- $\mathbb{H} = \mathbb{C}^2$ ;  $\forall f \in \mathbb{H} : f = f_1\epsilon_1 + f_2\epsilon_2$ ,  $\epsilon_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ ;  $\epsilon_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ .
- If  $\|f\| \neq 0$ , then  $\phi = \frac{f}{\|f\|} \in \mathbf{S}_p$ , i.e.  $|\phi_1|^2 + |\phi_2|^2 = 1$ , therefore  $\phi_1, \phi_2$  **amplitudes** of probability on  $\{1, 2\}$ .
- $\epsilon_1, \epsilon_2 \in \mathbf{S}_p$ .
- $|\phi_1|^2 = \mathbb{P}(\text{system in state } \phi \text{ is in state } \epsilon_1)$ .
- If  $\phi_1 \neq 0, \phi_2 \neq 0$ , the pure state  $\phi$  is in a **linear superposition** of other pure states.
- Time evolution of isolated system is invertible and preserves pure states.

# A very simple illustrative example

## Observables I

- Take  $X = \begin{pmatrix} 1 & 2i \\ -2i & -2 \end{pmatrix} = X^*$  as an example of sharp observable.
- $\mathbb{X} = \text{spec } X = \{-3, 2\}$ . Compute easily

Eigenvalues $x$	Eigenvectors $\mathbf{u}_x$	Orthoprojectors $M_x$
-3	$\frac{1}{\sqrt{5}} \begin{pmatrix} -i \\ 2 \end{pmatrix}$	$\frac{1}{5} \begin{pmatrix} 1 & -2i \\ 2i & 4 \end{pmatrix}$
2	$\frac{1}{\sqrt{5}} \begin{pmatrix} 2i \\ 1 \end{pmatrix}$	$\frac{1}{5} \begin{pmatrix} 4 & 2i \\ -2i & 1 \end{pmatrix}$

- $\sum_{x \in \text{spec } X} M_x = I_{\mathbb{H}}$ ,  $M_x \geq 0$  and  $M_x^2 = M_x$ . Thus  $(M_x)$  family of sharp elementary observables.
- For  $x \in \text{spec } X$ ,  $M_x \mathbb{H} = \mathbb{C} \mathbf{u}_x$ .
- $X^* = X \Rightarrow \langle \mathbf{u}_x | \mathbf{u}_y \rangle = \delta_{x,y}$ ,  $x, y \in \text{spec } X \Rightarrow M_x, M_y$  orthogonal orthoprojectors.



# A very simple illustrative example

## Observables II

- Since  $(\mathbf{u}_x)_x$  orthonormal basis,  $\mathbf{S}_p \ni \psi = \sum_x \alpha_x \mathbf{u}_x$ , with  $\sum_x |\alpha_x|^2 = 1$ , i.e.  $(|\alpha_x|^2)_x$  probability vector on spec  $X$ .
- $\mathbb{E}_\psi X = \langle \psi | X \psi \rangle = \langle \sum_x \alpha_x \mathbf{u}_x | (\sum_{x'} M_{xx'}) \sum_{x''} \alpha_{x''} \mathbf{u}_{x''} \rangle = \sum_x x |\alpha_x|^2$ .
- If  $X$  was a classical  $\mathbb{X}$ -valued r.v., with probability  $(p_x)_{x \in \mathbb{X}}$ ,

$$\mathbb{E}X = \sum_x x p_x = \sum_x \sqrt{p_x} x \sqrt{p_x} = \sum_x e^{-i\theta_x} \sqrt{p_x} x \sqrt{p_x} e^{i\theta_x} = \langle \psi | \hat{X} \psi \rangle,$$

where  $\psi = \begin{pmatrix} e^{i\theta_{-3}} \sqrt{p_{-3}} \\ e^{i\theta_2} \sqrt{p_2} \end{pmatrix}$  and  $\hat{X} = \begin{pmatrix} -3 & 0 \\ 0 & 2 \end{pmatrix}$ .

- Classical r.v. = quantum r.v. with only diagonal entries.

# A very simple illustrative example

## Measurement

- System in  $\psi \in \mathbf{S}_p$ . Ask question  $M_x$ , i.e. “does  $X$  takes value  $x$ ?” and determine its probability.
- Answer: yes, with probability  $\pi_M^\psi(x) = \langle \psi | M_x \psi \rangle = \|M_x \psi\|^2$ ,  
where  $M_x \psi = \begin{cases} \langle \mathbf{u}_x | \psi \rangle \mathbf{u}_x & \text{if } x \in \mathbb{X}, \\ 0 & \text{otherwise.} \end{cases}$
- **BUT:** Once the question has been asked, and got positive answer, the system now in new state  $\phi_x = \frac{\langle \mathbf{u}_x | \psi \rangle \mathbf{u}_x}{|\langle \mathbf{u}_x | \psi \rangle|}$ .
- Re-asking the question, gives answer “yes” with probability 1.
- Asking a question on a quantum system, **IRREVERSIBLY** alters its state.

## A very simple illustrative example

### Composite systems

- $\mathbb{H} = \mathbb{H}_1 \otimes \mathbb{H}_2$ ,  $\mathbb{H}_i \simeq \mathbb{C}^2$ . Each  $\mathbb{H}_i$  with basis  $(\epsilon_1, \epsilon_2)$ .
- Basis of  $\mathbb{H}$   $(\epsilon_1 \otimes \epsilon_1, \epsilon_1 \otimes \epsilon_2, \epsilon_2 \otimes \epsilon_1, \epsilon_2 \otimes \epsilon_2)$ .
- $\mathbb{H} \ni \psi = \psi_{11}\epsilon_1 \otimes \epsilon_1 + \psi_{12}\epsilon_1 \otimes \epsilon_2 + \psi_{21}\epsilon_2 \otimes \epsilon_1 + \psi_{22}\epsilon_2 \otimes \epsilon_2$ .
- If  $\psi_{11} = \psi_{12} = 0$ ,  $\psi_{21} \neq 0$ ;  $\psi_{22} \neq 0$ , then

$$\psi = \psi_{21}\epsilon_2 \otimes \epsilon_1 + \psi_{22}\epsilon_2 \otimes \epsilon_2 = \epsilon_2 \otimes (\psi_{21}\epsilon_1 + \psi_{22}\epsilon_2),$$

i.e.  $\psi = \psi^{(1)} \otimes \psi^{(2)}$ . **Joint probability = product measure (independence).**

- If both  $\psi_{11} \neq 0$ ;  $\psi_{22} \neq 0$  then  $\psi \neq \psi^{(1)} \otimes \psi^{(2)}$  **(non-independence)**. **Exercise: is-it possible classically?**
- “Quantenverschränkung”, “Quantum entanglement”.  
“Intrication (enchevêtrement) quantique”. Probably  
“Entrelazamiento (enredo) cuántico” in Spanish.

## An artist's view of entanglement



**Figure:** Ruth Bloch. Entanglement (bronze, 71cm, 1995). (From the web page of the artist, reproduced here with her permission; the existence of this sculpture has been brought to my attention by Dénes Petz).